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Problem set 10

Problem 1 Less-classical equations of motion for Bloch electrons packets

In the lecture we already saw a semi-classical derivation of the equations of motion for the effective wave vector \vec{K} of wave packets constructed from Bloch electrons in the presence of electromagnetic fields

$$\hbar \frac{d\vec{K}}{dt} = -e \left[\vec{E}(\vec{R}, t) + \frac{\vec{v}(\vec{K})}{c} \times \vec{B}(\vec{R}, t) \right], \quad (1)$$

based on expanding the vector potential $\vec{A}(\vec{r})$. In this problem, we will use a less "classical" yet more restricted approach to illustrate the crucial requirements leading to eq. (1). Consider a single-particle Hamiltonian

$$\hat{h} = \varepsilon \left((\hat{\vec{p}} + (e/c)\vec{A}(\hat{R}))/\hbar \right) + U(\hat{R}), \quad (2)$$

where the function $\varepsilon(\vec{k})$ denotes the dispersion relation, $U(\hat{R})$ a potential and $\hat{\vec{p}}, \hat{R}$ are operators with canonical commutation relations $[\hat{R}_\mu, \hat{p}_\nu] = i\hbar\delta_{\mu\nu}$.

Use the Heisenberg equations of motion and the identification $\hbar\vec{K} = \hat{\vec{p}} + (e/c)\vec{A}(\hat{R})$ to show eq. (1) for this Hamiltonian.

Problem 2 Holes change signs

We consider a single-particle dispersion relation approximated near the top of the band \vec{k}_0

$$\varepsilon(\vec{k}) \approx \varepsilon(\vec{k}_0) - \frac{\hbar^2}{2m^*} (\vec{k} - \vec{k}_0)^2, \quad (3)$$

with the effective mass m^* . Show that in this case wave packets constructed from Bloch states where $\vec{K} = \vec{k}_0$ obey the equations of motion

$$\frac{d^2\vec{R}}{dt^2} = +\frac{e}{m^*} \left[\vec{E}(\vec{R}, t) + \frac{\vec{v}(\vec{K})}{c} \times \vec{B}(\vec{R}, t) \right]. \quad (4)$$

Use eq. (3) to explain why it is consistent to interpret these equations of motion as those of a hole-excitation.