Quantum Field Theory (Quantum Electrodynamics)

Problem Set 12

22 & 24 January 2024

1. Quantization of Maxwell theory

1. Let's now move to the quantization of the Maxwell theory, by promoting the *a*'s into operators. Imposing canonical commutation relations between A_{μ} and \dot{A}_{ν} , i.e.

$$[A_{\mu}(t,\vec{x}), A^{\nu}(t,\vec{x'})] = 0, \quad [\dot{A}_{\mu}(t,\vec{x}), A^{\nu}(t,\vec{x'})] = i\delta^{\nu}_{\mu}\delta^{(3)}(\vec{x}-\vec{x'}) ,$$

show that the non-vanishing commutation relations between the ladder operators are

$$[a_r(\vec{k}), a_s^+(\vec{k}')] = 2\omega_{\vec{k}}(2\pi)^3 \zeta_s \delta_{rs} \delta^{(3)}(\vec{k} - \vec{k}') \; .$$

2. Argue why the Lorenz gauge condition cannot be imposed on operators, but should instead be imposed on states. Argue that fixing the gauge corresponds to

$$\left(a_0(\vec{k}) - a_3(\vec{k})\right)|\psi\rangle = 0 ,$$

(the Gupta–Bleuler condition), where $|\psi\rangle$ is a physical state.

- 3. Consider the state $a_0^+(\vec{k}) |0\rangle$ and compute its norm. Is this state physical? Modify this state accordingly, such that it satisfies the Gupta-Bleuler condition. Is it now physical?
- 4. Express the Hamiltonian in terms of the ladder operators. Check that its expectation value in the modified states which you constructed in the previous point vanishes.
- 5. What statistics do vector particles obey?
- 6. Compute the Feynman propagator for the vector field.

2. Quantum Electrodynamics (QED)

The Lagrangian of QED reads

$$\mathcal{L} = -\frac{1}{4} F_{\mu\nu} F^{\mu\nu} + i\bar{\psi}(\partial \!\!\!/ - m)\psi - e\bar{\psi} A\!\!\!/ \psi , \qquad (1)$$

with ψ a Dirac fermion with mass m.

- 1. Show that the Lagrangian is invariant under the gauge transformation $\psi \to \psi' = e^{-ie\alpha(x)}\psi$, $A_{\mu} \to A'_{\mu} = A_{\mu} + \partial_{\mu}\alpha(x)$, with a(x) an arbitrary function.
- 2. The interaction part H_{int} of the Hamiltonian reads

$$H_{int} = e \int \mathrm{d}^3 \vec{x} \, \bar{\psi} \mathcal{A} \psi$$

Compute the matrix element

$$\langle f | \int \mathrm{d}t \, H_{int} | i \rangle \;$$

where the initial state comprises a fermion-antifermion pair, while the final state a photon. Can this process take place?

3. Let's add to (1) an additional Dirac fermion Ψ with mass M, such that

$$\tilde{\mathcal{L}} = \mathcal{L} + i\bar{\Psi}(\partial \!\!\!/ - M)\Psi - e\bar{\Psi}A\!\!\!/ \Psi , \qquad (2)$$

where we may interpret ψ as an electron and Ψ as a muon. Compute, at the lowest nonvanishing order, the scattering amplitude for the process $\psi \Psi \to \psi \Psi$.

4. Square and average the above, write it in terms of the Mandelstam variables and derive the differential cross-section in the center-of-mass frame.

3. Feynman Rules

- 1. Work out the Feynman rules for the Lagrangian (2).
- 2. Using the above, compute the amplitude corresponding to the process $\psi \Psi \rightarrow \psi \Psi$ and compare your findings with Exercise 1, point 2.
- 3. Assuming that m = 0, compute the amplitude for the $\psi\psi \rightarrow \psi\psi$ scattering in terms of the Mandelstam variables. Derive the differential cross-section for this process.
- 4. Write the Feynman diagrams that contribute to the Compton scattering $\psi \gamma \rightarrow \psi \gamma$ and find the corresponding amplitude. Compute the differential cross section for this process in the laboratory frame.