


Seesaw mechanism of neutrino mass

• $G_{SM} = SU(2)_L \times U(1)_Y$

• minimal fermion

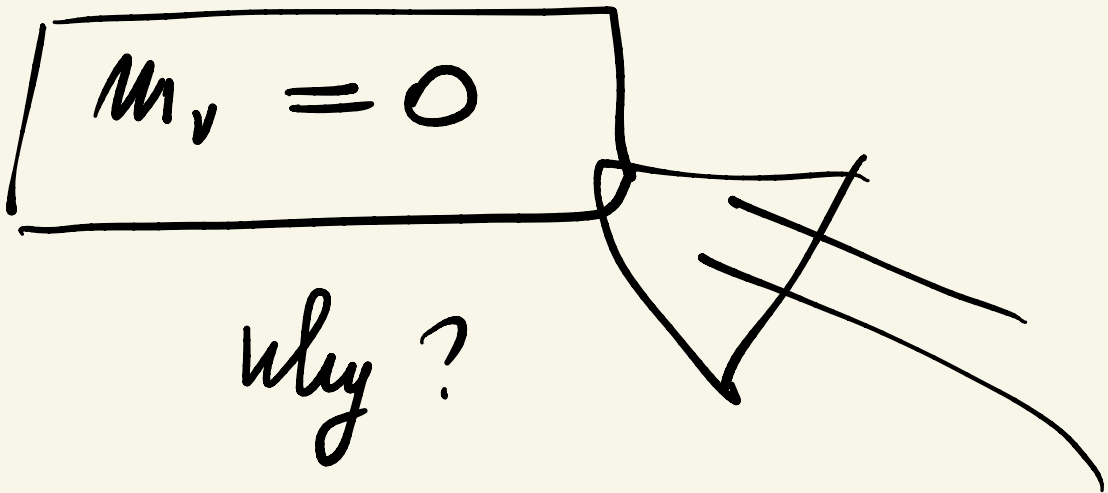
$$q = \begin{pmatrix} u \\ d \end{pmatrix}_L \quad u_R, d_R$$

$$l = \begin{pmatrix} \nu \\ e \end{pmatrix}_L \quad e_R$$

• minimal Higgs

$$\Phi$$





(a) $\exists \nu_L$ only

\Downarrow

(b) $\nu_L^T c \nu_L$ not allowed

~~$S U(2)$~~ , ~~$U(1)$~~

\Downarrow natural

$$\exists \psi_R$$

$$\gamma(\psi_R) = 0$$

$$T_a(\psi_R) = 0$$

$$a = 1, 2, 3$$

$$\left(\begin{array}{l} \cancel{E = m a^2} \\ \cancel{E = m b^2} \\ E = m c^2 ! \end{array} \right)$$

\Downarrow

$$\mathcal{L}_y(e) = \bar{l}_1 i \sigma_2 \bar{\Phi}^* \psi_R \psi_D + h.c.$$

$$+ \frac{1}{2} \psi_R^T C \psi_R \underline{M}_R + h.c.$$

$$\text{if } M_R = 0$$

$$\Rightarrow \mu_\nu = y_0 \langle \Phi \rangle$$

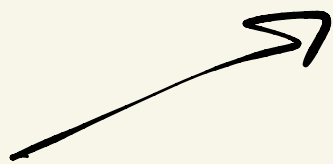
$$y_0 = \frac{\mu_\nu}{\langle \Phi \rangle} = \frac{g}{2} \frac{\mu_\nu}{M_W}$$

$$\mu_\nu \leq 1 \text{ eV}$$

$$\Rightarrow y_0 < 10^{-11}$$

$$\Gamma(h \rightarrow \nu\bar{\nu}) \propto y_0^2$$

$$\Rightarrow \mathcal{B}(h \rightarrow \nu\bar{\nu}) \leq 10^{-22} (=0)$$



true only when $M_R = 0$

- $N_L \equiv C \bar{V}_R^T = C \gamma_0 V_R^* = i \gamma_2 V_R^*$

$$M_R V_R^T C V_R + V_R^\dagger C + V_R^* M_R^*$$

$$\Rightarrow N_L^T = \bar{V}_R C^T = V_R^\dagger \gamma_0 C^T$$

$$\Rightarrow N_L^T C N_L = V_R^\dagger \gamma_0 \overbrace{C^T C}^{\equiv 1} C \gamma_0 V_R^*$$

$$= V_R^\dagger \gamma_0 C \gamma_0 V_R^*$$

$$= V_R^\dagger (-C) V_R^* = V_R^\dagger C^\dagger V_R^*$$



$$M_R V_R^T C V_R + M_R^* V_R^\dagger C^\dagger V_R^* =$$

$$= M_N \overline{N_L^T} C N_L + h.c.$$

$$\boxed{M_N = M_A^*}$$

$$\langle \mathcal{L}_1 \rangle = \left(\bar{v}_L Y_D v_R + \overbrace{\bar{v}_R Y_D^* v_L}^{M_D} \right) \langle \phi \rangle$$
$$+ \frac{1}{2} N_L^T C M_N N_L + \text{h.c.}$$

but: $\bar{v}_R v_L = ?$

$$N_L = C \bar{v}_R^T$$

$$\Rightarrow N_L^T = \bar{v}_R C^T$$

$$\Rightarrow N_L^T C = \bar{v}_R C^T C = \bar{v}_R$$



$$\bar{\nu}_R \nu_L = N_L^T C \nu_L$$



$$\mathcal{M}(\nu) = N_L^T C M_D \nu_L + h.c.$$

$$+ \frac{1}{2} N_L^T C M_N N_L + h.c.$$



$$M_D^T = M_N$$

$$= \frac{1}{2} N_L^T C M_D \nu_L + \frac{1}{2} N_L^T C M_D \nu_L$$

$$+ \frac{1}{2} N_L^T C M_N N_L + h.c.$$

(-c)
⇓

$$= \frac{1}{2} N_L^T C M_D \nu_L + \frac{1}{2} (-) \nu_L^T C^T M_D^T N_L$$

$$+ \frac{1}{2} N_L^T C M_N N_L + h.c.$$



$$M(v) = \frac{1}{2} \left(N_L^T C M_D v_L + v_L^T C M_D^T N_L \right. \\ \left. + M_N N_L^T C N_L \right) + h.c.$$

$$M_{(v, n)} = \begin{matrix} & \Downarrow & v & & N \\ & & & & \\ v & & 0 & & M_D^T \\ & & & & \\ N & & M_D & & M_N \end{matrix}$$

$$\Downarrow \text{1 gen}$$

$$M_{(D,N)} = \begin{pmatrix} 0 & \mu_D \\ \mu_D & \mu_N \end{pmatrix}$$

$$\boxed{\mu_N \gg \mu_D}$$

"seesaw"
approximation

$$\mu_1 + \mu_2 = \mu_N \quad (1)$$

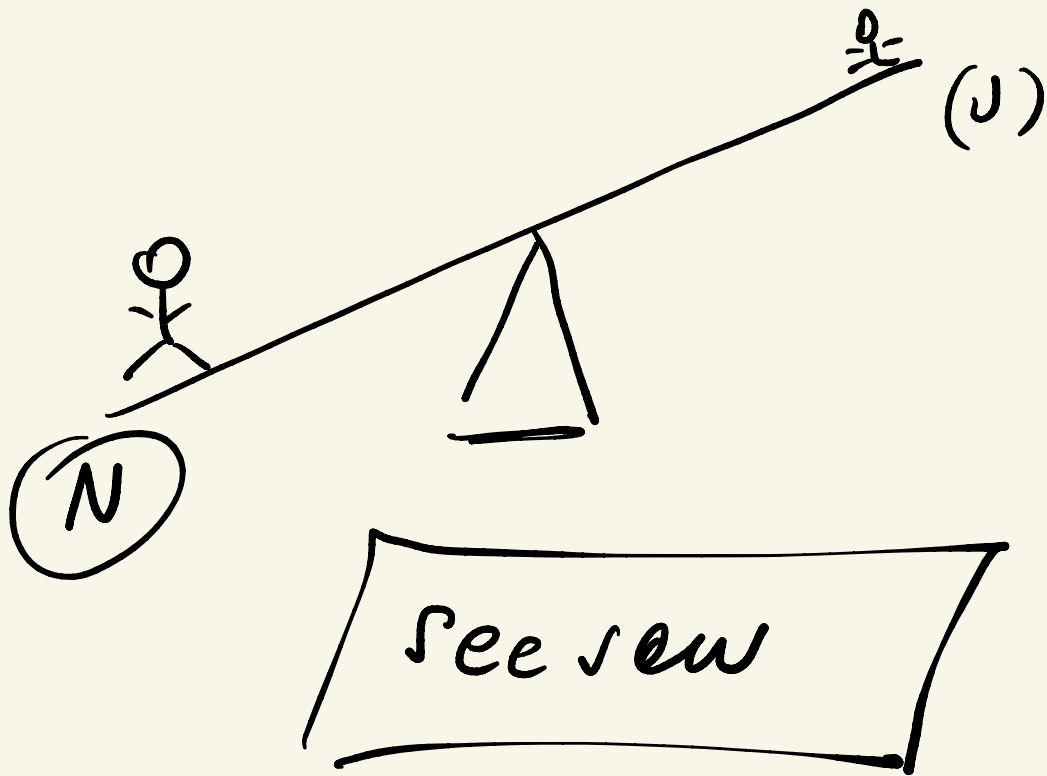
$$\mu_1 \cdot \mu_2 = -\mu_D^2 \quad (2)$$



$$\mu_1 \ll \mu_2$$

$$(1) \Rightarrow \boxed{\mu_2 \approx \mu_N}$$

$$(2) \quad \mu_1 = -\frac{\mu_D^2}{\mu_2} \approx \frac{\mu_D^2}{\mu_N} \approx \mu_D$$



physical

$$\begin{pmatrix} D \\ N \end{pmatrix} = U \begin{pmatrix} D \\ N \end{pmatrix}$$

$$\begin{pmatrix} \mu_{D_1} \equiv \mu_1 \\ \mu_{D_2} \equiv \mu_2 \end{pmatrix}$$

unphysical

1 rev. $V = 0 = \begin{pmatrix} \cos\theta & \sin\theta \\ -\sin\theta & \cos\theta \end{pmatrix}$

$$\begin{pmatrix} v_1 \\ v_2 \end{pmatrix} \equiv \begin{pmatrix} v'_1 \\ v'_2 \end{pmatrix} \leftarrow \text{physical}$$

$$\Rightarrow \frac{1}{2} \tan 2\theta \rightarrow 0$$

$\underbrace{\hspace{10em}}_{\mu_D \rightarrow 0}$

$$\approx \theta \quad (\theta \leq 1)$$

$$\frac{1}{2} \tan 2\theta \rightarrow 0$$

$\mu_N \rightarrow \infty$



$$\theta \approx \frac{1}{2} \tan 2\theta = \frac{u_D}{u_N} \ll 1$$

\Uparrow from

$$S = \begin{pmatrix} a & c \\ c & b \end{pmatrix}$$

$$\Rightarrow \frac{1}{2} \tan 2\theta = \frac{c}{b-a}$$

bottom line:

$$v' \approx v + \theta N$$

$$N' \approx N - \theta v$$

from now ν', N' = called
 ν, N

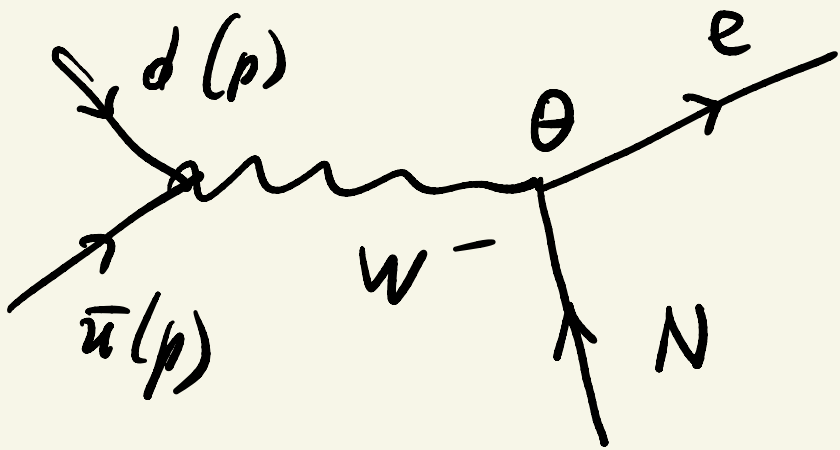
• see row ($M_N \gg M_D$)

used as a physical case!

$$M_N > M_W$$

$$\mathcal{L}_{int} = \frac{g}{\sqrt{2}} \bar{\nu}_L \gamma^\mu e_L W_\mu^+ + h.c.,$$

$$\rightarrow \frac{g}{\sqrt{2}} \bar{N}_L \gamma^\mu e_L W_\mu^+ \theta + h.c.,$$



$$\Rightarrow \sigma(N) \propto \theta^2 \approx \frac{m_D^2}{m_N^2}$$

$$\left| \frac{m_D}{m_N} \right| = \frac{m_D^2}{m_N m_N} \approx \frac{m_D^2}{m_N}$$

$$\Rightarrow \sigma(N) \propto \frac{m_D}{m_N} \rightarrow 0$$

N cannot be produced
in see saw picture

• $\gamma \perp$ generation

\Downarrow

$$\underline{M}_{(v, w)} = \begin{pmatrix} 0 & M_D^T \\ M_D & M_N \end{pmatrix}$$

$$\begin{pmatrix} v \\ N \end{pmatrix} \rightarrow U \begin{pmatrix} v \\ N \end{pmatrix}$$

$$U U^+ = 1 \quad (\text{up to } \theta^2 \text{ order})$$

$$U = \begin{pmatrix} 1 & \sim \theta^+ \\ \sim \theta & 1 \end{pmatrix}$$

$$U^T M_{(\nu N)} U = \begin{pmatrix} M_\nu & 0 \\ 0 & M_N \end{pmatrix}$$

$$\Rightarrow \frac{M_\nu}{\parallel (M_\nu \gg M_0)} = -M_D^T \frac{1}{M_N} M_D$$

symmetric ! $M_\nu^T = M_\nu$
 $M_N^T = M_N$

$$\begin{aligned} & \nu_i^T (M_\nu)_{ij} \nu_j^T = \\ & = -\nu_j^T C^T (M_\nu)_{ij} \nu_i^T \\ & = \nu_j^T C (M_\nu^T)_{ji} \nu_i^T \quad \text{Q.E.D.} \end{aligned}$$

$$(C^T = -c)$$

$$(2) \quad \theta = \frac{1}{M_N} M_D$$
$$\leq 1, \quad M_N \gg M_D$$

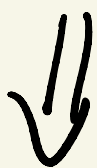
• neutrino mass \leftrightarrow neutrino oscillations

$$(a) \quad \nu_\mu \rightarrow \nu_\tau \quad (\text{atm. neutrino})$$

$$\Delta m_{ATM}^2 \approx 10^{-3} \text{ eV}^2$$

(b) $\nu_e \rightarrow \nu_\mu$ (solar neutrinos)

$$\Delta m_{21}^2 \approx 10^{-5} \text{ eV}^2$$



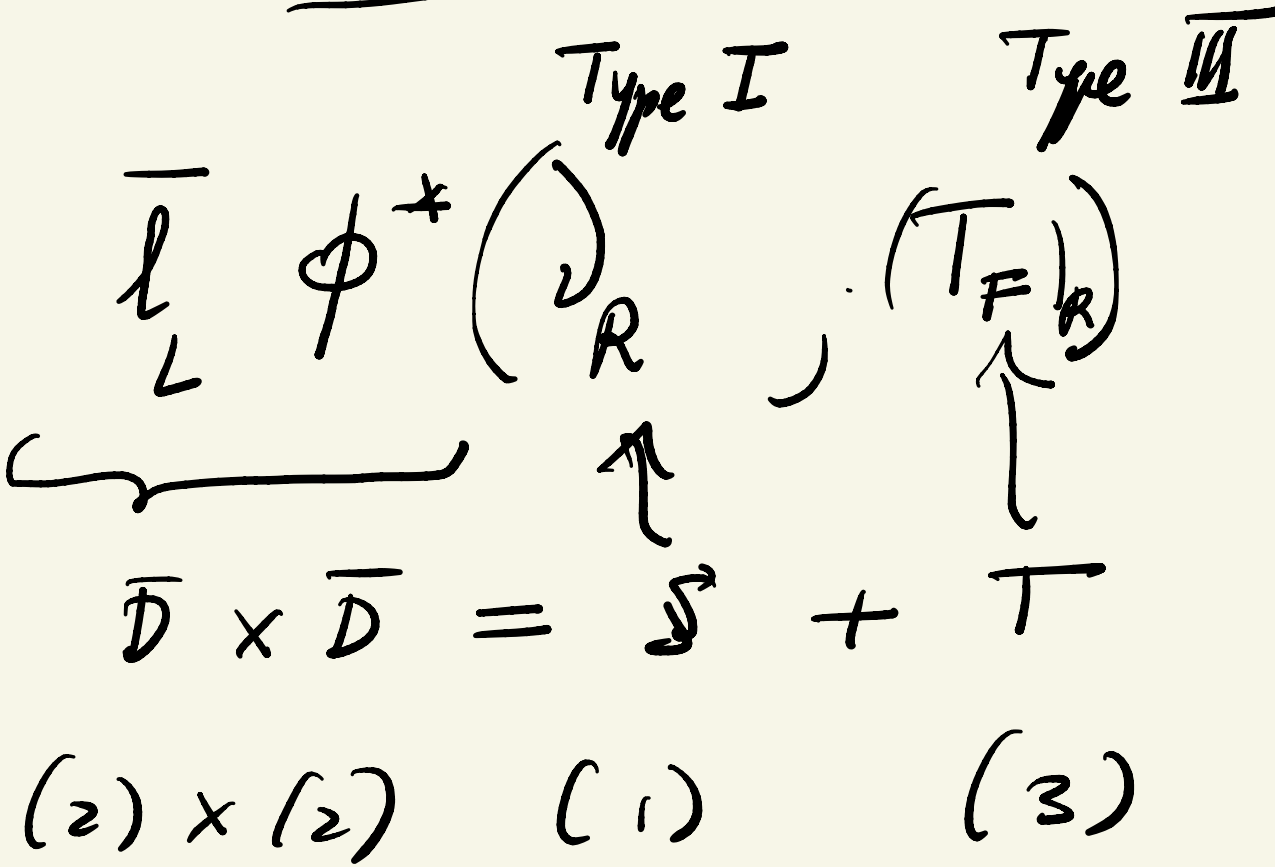
Q. how many massive neutrinos?

A. at least two!



at least two N

Other sources of
neutrino mass



$\chi(\nu_R) = 0, \chi(T_{FR}) = 0$

$$\nu_R \rightarrow N_L$$

$$(T_F)_R \rightarrow (T_F^c)_L$$

$$\therefore \frac{g}{\sqrt{2}} \overline{T_{0F}} \gamma^\mu T_{-F} W_\mu^+$$

$$(T_F)_L = \begin{pmatrix} T^+ \\ T_0 \\ T^- \end{pmatrix}_L$$

Q. If no new fermion

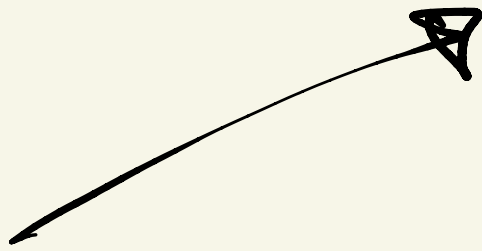
\Rightarrow what?

\Downarrow

$$\underbrace{\nu_L^T C \nu_L}$$

repr. under $SU(2)$?

$$T_3: \left(-\frac{1}{2}\right) + \left(-\frac{1}{2}\right) = -1$$



scalar triplet (Δ)!

$$l = \begin{pmatrix} \nu \\ e \end{pmatrix}_L, \quad e_R$$



$$\Delta \rightarrow U \Delta U^\dagger$$

$$L_Y(\text{new}) = l_L^T \overset{\tau}{C} \Rightarrow \boxed{Y(\Delta) = 2} l_L Y_\Delta + \text{h.c.}$$

$$\rightarrow l_L^T C U^T i\sigma_2 U \Delta U^+ U l_L Y_\Delta$$

$$\rightarrow l_L^T C i\sigma_2 \underbrace{U^+ U}_1 \Delta \underbrace{U^+ U}_1 l_L Y_\Delta$$

$\Rightarrow \text{inv.}$

Q.E.D.

$$V = V_{SM} + \frac{M_\Delta^2}{2} \overbrace{Tr \Delta^+ \Delta}^{\Delta_0 \Delta_0^* + \dots} + \dots + V_{\phi \Delta}$$

(ϕ)

$$\parallel \\ -\mu \phi^T i\sigma_2 \Delta^+ \phi$$

$$\frac{\partial V}{\partial \Delta_0^*} = m_\Delta^2 \Delta_0 - \mu \langle \phi \rangle \langle \phi \rangle + \dots = 0$$

$$\Rightarrow \langle \Delta_0 \rangle = +\mu \frac{\langle \phi \rangle^2}{m_\Delta^2}$$

induced (small) vev

$$\langle \Delta_0 \rangle \leq \text{GeV}$$

$$M_\nu = Y_\Delta \langle \Delta_0 \rangle$$

Type II seesaw

$\mu \sim m_\Delta \Rightarrow$

$$M_\nu \approx Y_\Delta \frac{M_w^2}{m_\Delta} \quad (\text{II})$$
$$M_\nu \approx \frac{Y_\Delta^2 M_w^2}{M_N} \quad (\text{I})$$