

GUT Course

2022 / 2023

Lecture IV

4/11 / 2022

L MU

Fall 2022



S_M - wave

$$\bar{\Phi} = \text{doublet} \quad \bar{\Phi} \rightarrow U \bar{\Phi}$$

$$Y(\bar{\Phi}) = 1$$

$$\Rightarrow \langle \bar{\Phi} \rangle = \bar{\Phi}_0 = \begin{pmatrix} 0 \\ v \end{pmatrix}$$



$$M_W = \frac{g}{2} v \quad M_t = \frac{M_W}{\cos \theta_W}$$

$$M_A = 0$$

$$\bar{\Phi}_0 \therefore T_a \bar{\Phi}_0 \neq 0$$

$$Y \bar{\Phi}_0 \neq 0$$

$$T_3 \Phi_0 = -\frac{1}{2} \Phi_0$$

$$Y \Phi_0 = \Phi_0$$

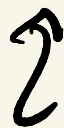


$$Q_{em} = T_3 + \frac{Y}{2} \therefore$$

$$Q_{em} \Phi_0 = 0$$

$$\Leftrightarrow M_A = 0$$

$$A = \sin \theta_w A_3 + \cos \theta_w B$$



$SU(2)$



$U(1)$

$$\sin \theta w = \frac{e}{f}, \quad \cos \theta w = \frac{e}{f'}$$

$$\left(\tan \theta w \equiv f'/f \right)$$



$$\frac{1}{e} A = \frac{1}{f} A_3 + \frac{1}{f'} B$$

$$Q = T_3 + \frac{y}{z}$$

e

f

f'

to be generalized

- $\bar{\Phi}_0^6 = \begin{pmatrix} 0 \\ e \end{pmatrix}$

- $\bar{\Phi}_0^H = \begin{pmatrix} e \\ 0 \end{pmatrix}$ what then?

$$T_3 \bar{\Phi}_0^H = \frac{1}{2} \bar{\Phi}_0^H$$

$$\frac{4}{2} \bar{\Phi}_0^H = \frac{1}{2} \bar{\Phi}_0^H$$

$$\Rightarrow Q_{em} \bar{\Phi}_0^H \neq 0$$

? ? ? ?
• • • •



$$Q_{em}^H = T_3 - \frac{4}{2}$$

$$Q_{em}^H \bar{\Phi}_0^H = 0$$



$$\boxed{Q_{em}^6 \Phi_0^6 = 0}$$

$$Q_{em}^6 = T_3 + \frac{Y}{2}$$

• $\Phi_0^3 = \begin{pmatrix} v_1 \\ v_2 \end{pmatrix} \Rightarrow$

$$Q_{em}^3 \therefore Q_{em}^3 \Phi_0^3 = 0$$

$$\boxed{Q_{em}^3 = 0 \quad ??}$$

$$\boxed{SU(N)} \quad [T_a, T_b] = i f_{abc} T_c$$

$$T_a' = U T_a U^\dagger$$

$$[T_a', T_b'] = [U T_a U^\dagger, U T_b U^\dagger]$$

$$= U [T_a, T_b] U^\dagger = U i f_{abc} T_c U^\dagger$$

$$= i f_{abc} T_c'$$

$$S.M \Rightarrow Q_{em} \bar{B}_0 \neq 0$$

$$M_A = 0$$

$$M_A^{exp} \leq 10^{-16} eV$$

$$\begin{aligned} \mathcal{L}_Y = & \gamma_e \bar{l}_L \Phi l_R + \gamma_d \bar{q}_L \Phi d_R + \\ & + \gamma_u \bar{q}_L i \sigma_2 \Phi^* u_R + h.c. \end{aligned}$$

$$\Phi_0 = \begin{pmatrix} 0 \\ v \end{pmatrix}$$

$$\Rightarrow \gamma_f \bar{f}_L f_R v + \text{h.c.} = \\ = \gamma_f v \bar{f} f$$



$u_f = \gamma_f v$	←
$M_W = \frac{g}{2} v$	←

⋈

the scale of weak int.

$$\Rightarrow \left[\gamma_f = \frac{u_f}{v} = \frac{g}{2} \frac{u_f}{M_W} \right]$$

• $\Phi \longrightarrow U(x)\Phi \quad [T_a, T_b] = i\epsilon_{abc} T_c$

$$U(x) = e^{i\theta_a(x) T_a}$$

$$\Phi \xrightarrow{U} \begin{pmatrix} 0 \\ v+h(x) \end{pmatrix} = \Phi_{un}$$

Higgs boson

$$\Phi \text{ (4 real d.o.f.)} \rightarrow \Phi_{un} \text{ (1 d.o.f.)}$$

$$\underbrace{A_i, B}_{4 \text{ massless}} \longrightarrow \underbrace{W^+, W^-, Z}_{3 \text{ massive}} ; \underbrace{A}_{M_A=0}$$



3 extra d.o.f.

$$3 + 1 = 4$$

2 matches

⇓

$$\begin{aligned} \mathcal{L}_{kin} &= \frac{1}{2} (D_\mu \Phi)^\dagger (D^\mu \Phi) \\ &= \frac{1}{2} (\partial_\mu h)^2 + \frac{1}{2} M_W^2 W^+ W^- \left(1 + \frac{h}{v}\right)^2 \\ &\quad + \frac{1}{2} M_Z^2 z z \left(1 + \frac{h}{v}\right)^2 \end{aligned}$$

$$- \gamma_f v \bar{f} f \left(1 + \frac{h}{v}\right) +$$

$$\Phi_{kin} = \begin{pmatrix} 0 \\ v+h \end{pmatrix} = \begin{pmatrix} 0 \\ v \end{pmatrix} \left(1 + \frac{h}{v}\right)$$

$$\Downarrow \left[\begin{array}{l} M_W = \frac{g}{2} v \\ m_f = \gamma_f v \end{array} \right]$$

$$\Downarrow \rightarrow v = \frac{2 M_W}{g}$$

h : couples to your mass
(if it gives you a mass)



$$\frac{1}{2} M_W^2 w^+ w^- \left(1 + \frac{h}{v}\right)^2$$

$$= \frac{1}{2} M_W^2 w^+ w^- + M_W^2 w^+ w^- \frac{h}{v}$$

$$\downarrow$$

$$2 \frac{g}{M_W} M_W^2 w^+ w^- h$$

$$\boxed{2 h M_{Wg} W^+ W^-} \quad \checkmark$$

• fermions

$$\boxed{h y_f \bar{f} f = \frac{g}{2} \frac{m_f}{M_W} h \bar{f} f}$$

mass = dynamical

$$\Gamma(h \rightarrow f \bar{f}) \propto y_f^2 m_h$$

↓

$$\Gamma(h \rightarrow f\bar{f}) \propto g^2 \frac{m_f^2}{M_W^2} m_h$$

today

$$(m_h \approx 125 \text{ GeV})$$

w, z, t, b, τ

Higgs origin of mass

- $w^- \rightarrow e \bar{\nu}$
 - $z \rightarrow e \bar{e} (f \bar{f})$
- } ⑧

$$\frac{g}{\sqrt{2}} W_{\mu}^{-} \bar{e} \gamma^{\mu} L \nu$$

$$\frac{g}{\cos\theta_w} Z_{\mu} \bar{f} [T_3 - Q \sin^2\theta_w] f$$

$$\Gamma(W \rightarrow e + \bar{\nu}) \propto g^2 f(m_e)$$

kinematics

$$\underbrace{\Phi_1, \Phi_2}$$

$$\Phi_1^0 = \begin{pmatrix} 0 \\ v_1 \end{pmatrix}, \quad \Phi_2^0 = v_2 \begin{pmatrix} \cos\theta \\ \sin\theta \end{pmatrix}$$

Ques $\Phi_0 = 0$ NOT
a prediction

$$V = \dots + (\Phi_1 + \Phi_2) (\Phi_2 + \Phi_1) \lambda$$

$$\therefore \Phi_2^0 = \begin{pmatrix} 0 \\ v_2 \end{pmatrix}$$

$$\Phi_1 = \begin{pmatrix} 0 \\ v_1 + h_1 \end{pmatrix}$$

$$\Phi_2 = \begin{pmatrix} H^+ \\ v_2 + h_2 + i G \end{pmatrix}$$

$$h = f(h_1, h_2)$$

$$M_0 \neq h$$



limit as its mass

Higgs sector

$$V = \frac{\lambda}{4} (\Phi^\dagger \Phi - v^2)^2$$

$$\Phi_{\text{vac}} = \begin{pmatrix} 0 \\ v+h \end{pmatrix}$$



$$V = \frac{\lambda}{4} (v^2 + h^2 + 2vh - v^2)^2$$

$$= \frac{1}{2} \underbrace{2\lambda v^2 h^2} + \frac{\lambda}{4} h^4 + \lambda v h^3$$

$$\boxed{m_h^2 = 2\lambda v^2}$$

⇓

$m_h = \sqrt{2\lambda} v$ $m_W = \frac{g}{2} v$ $m_f = y_f v$	Higgs mass on same footing
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$$\alpha_w = \frac{g^2}{4\pi} \quad , \quad \alpha_f = \frac{y_f^2}{4\pi} \quad ,$$

$$\alpha_h = \frac{2\lambda}{4\pi}$$

∴

$$\alpha_i \ll L$$

perturbativity

new fermions:

$$m_f \ll \text{TeV}$$

$$(m_h \ll \text{TeV})$$

cut-off $\Lambda \approx M_{\text{pl}}$
($\approx 10^{19}$ GeV)

$$\alpha_N \approx \frac{1}{M_{pl}^2}$$

$$\alpha_N = \alpha_G \approx \frac{E^2}{M_{pl}^2} \ll 1$$
$$E \ll M_{pl}$$

- Higgs boson mass not predicted!?

NO!

$M_h =$ "predicted"

$$M_h^2 = 2 \lambda v^2$$

$$M_W = \frac{g}{2} v$$

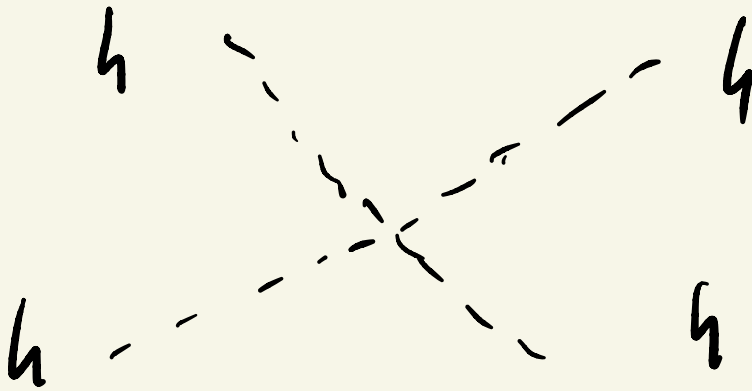
$$M_h^2 = 8 \lambda / g^2 M_W^2$$

$$(g \approx 0.6)$$

$$(\lambda \approx 1/10?)$$

$M_h \leftrightarrow \lambda$ connection

but: λh^4



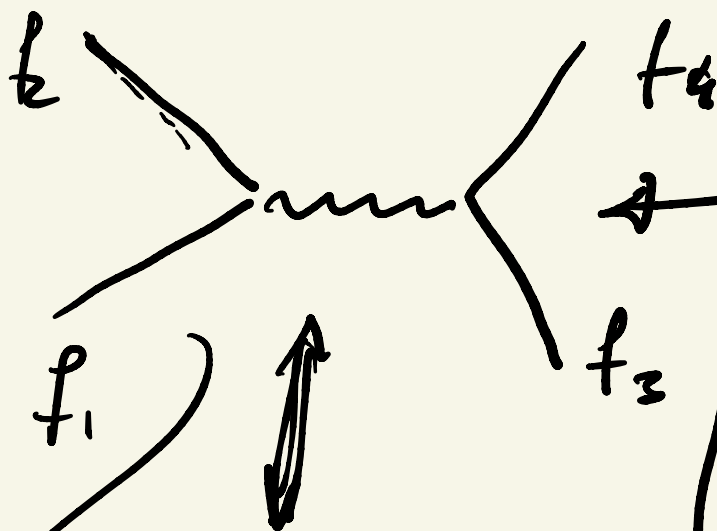
$$h + h \rightarrow h + h$$

$$\sigma(h+h) \propto \lambda^2$$

if you measure \rightarrow

\Rightarrow predict M_h !

compare with \bar{W} users



$$\Delta_{\mu\nu} \propto \frac{g_{\mu\nu}}{\cancel{k^2} - M_W^2}$$

$(k \ll M_W)$

$$\mathcal{L}_{\text{int}} = g/\sqrt{2} \bar{u}_L \gamma^\mu d_L W_\mu^+ + \text{h.c.}$$

$$\rightarrow \frac{g^2}{2} \frac{1}{M_W^2} \bar{u}_L \gamma^\mu d_L \bar{\nu}_L \gamma_\mu e_L$$

($d \rightarrow u + e + \bar{\nu}$)

$$= \frac{g^2}{8M_W^2} \bar{u} \gamma_\mu (1 + \gamma_5) d \bar{\nu} \gamma^\mu (1 + \gamma_5) e$$

||

$$\frac{G_F}{\sqrt{2}} = \frac{g^2}{8M_W^2}$$

← measure

↑
measure

↓

"predict" M_W

Hierarchy (Higgs mass)

issue

$$V = \frac{\lambda}{4} (\Phi^\dagger \Phi - v^2)^2$$

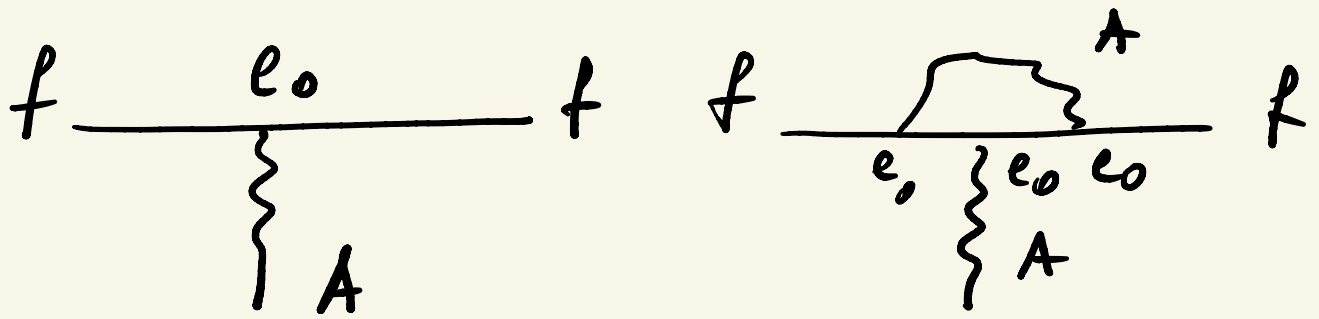
$$= \frac{\lambda}{4} (\Phi^\dagger \Phi)^2 - \frac{\lambda}{2} v^2 \Phi^\dagger \Phi + \dots$$

$$= \frac{\lambda}{4} (\Phi^\dagger \Phi)^2 - \underbrace{\frac{\lambda}{2} v^2 \Phi^\dagger \Phi}_{-\frac{\mu^2}{2} \Phi^\dagger \Phi} + \dots$$



Higgs mass term

$\sim v$



$$l = l_0 \left(1 + \frac{\alpha_0}{\pi} \ln \frac{\lambda}{u_f} \right)$$

$$\approx l_0 \left(1 + \frac{\alpha}{\pi} \ln \frac{\lambda}{u_f} \right)$$



$$l_0 \approx e \left(1 - \frac{\alpha}{\pi} \ln \frac{\lambda}{u_f} \right)$$

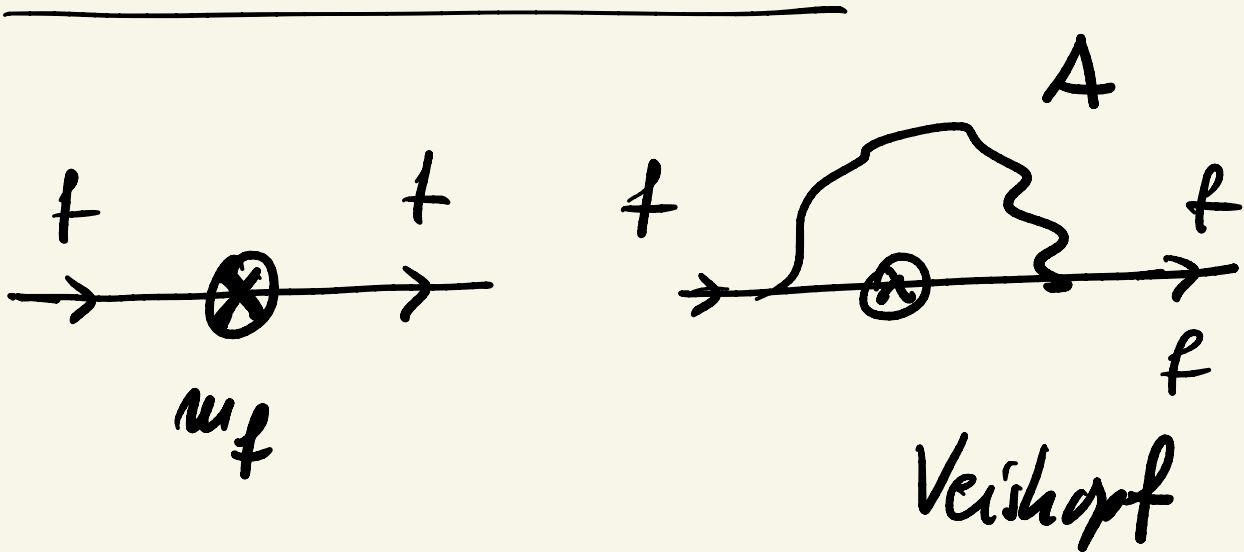


work with e

$$\text{Amplitudes} \propto e^{2\alpha} \left(1 + \frac{\kappa}{\lambda} \right)$$

z_{ev} ,

$$\lambda \rightarrow \infty$$



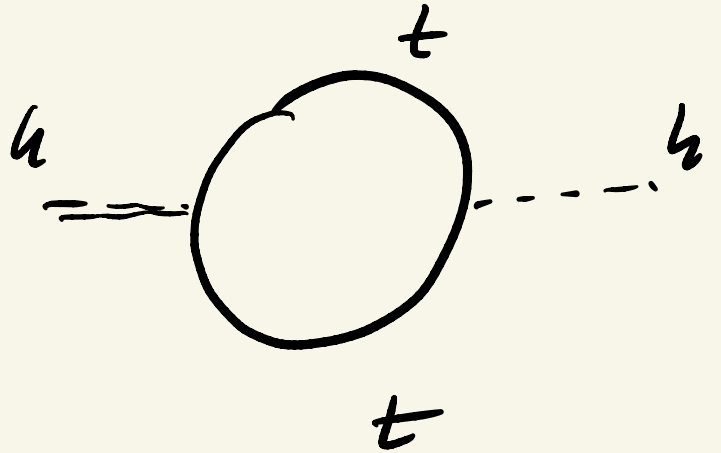
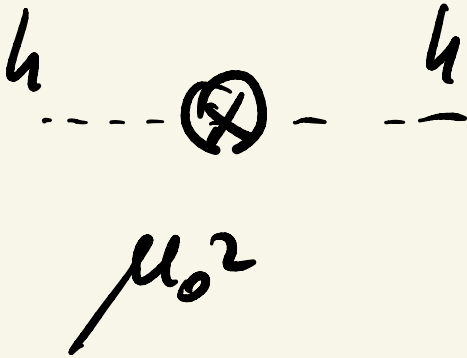
$$\hat{u}_f = u_f^0 \left[1 + \frac{\kappa}{\pi} \ln \frac{\lambda}{u_f} \right]$$

why?

small

($\ll 1$)

Scalar



$$\mu^2 = \mu_0^2 + \frac{4t^2}{16\pi^2} \Lambda^2 \approx \frac{1}{2} (125 \text{ GeV})^2$$

hierarchy "problem"

$$\mu^2 = \lambda v^2 = \frac{1}{2} m_h^2$$