

GUT and Neutrinos


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Lecture III

28 / 10 / 2022

LMU

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# SM : The theory

$$U(1) \rightarrow SU(2)_L \times U(1)_Y$$

QED                      ew gauge

1961

$i=1,2,3$

$$D_\mu = \partial_\mu - ig T_i A_\mu^i - ig' \frac{Y}{2} B_\mu$$

$$T_{iL} = \frac{\sigma_i}{2} \quad T_{iR} = 0$$

$$Q = T_3 + \frac{Y}{2}$$

$$Q \in C = \{ [T_\alpha, T_\beta] = 0 \}$$



$$\begin{array}{l|l} \begin{pmatrix} u \\ d \end{pmatrix}_L \equiv q_L & u_R, d_R \\ \begin{pmatrix} \nu \\ e \end{pmatrix}_L \equiv l_L & e_R \end{array}$$

$$\Rightarrow \mathcal{L}_{\text{weak}}(W) =$$

$$= \frac{g}{\sqrt{2}} (\bar{u}_L \gamma^\mu d_L + \bar{\nu}_L \gamma^\mu e_L) W_\mu^+$$

$$W_\mu^\pm = \frac{(A_1 \mp i A_2)_\mu}{\sqrt{2}}$$

$$\Rightarrow \text{if } \exists A_\mu \therefore$$

$$\mathcal{L}_A = e \bar{f} \gamma^\mu f Q A_\mu$$

$\Downarrow$

$$\mathcal{L}_2 = \frac{g}{\cos \theta_w} \bar{\psi} \gamma^\mu [T_3 - Q \sin^2 \theta_w] \psi + \mathcal{L}_\mu$$

$$\tan \theta_w = g'/g$$

$$e = g \sin \theta_w$$

$$\Leftrightarrow \theta_w^{\text{exp}} \approx 30^\circ$$

$$\sin^2 \theta_w^{\text{exp}} = 0.23$$

$$e^2/4\pi \equiv \alpha = \alpha_{em} \approx 1/100$$

$$g^2/4\pi \equiv \alpha_2 = \alpha_w \approx 1/25$$

renormalizable = good theory

$\Rightarrow$  finite amplitudes

$\Downarrow$   
gauge symmetry

$\Downarrow$  Weinberg '67

Higgs mechanism is SM

(t Hooft '71-'73)

$\Rightarrow$  add a scalar multiplet  $\phi$   
to the theory  $\therefore$

$$\langle \phi \rangle \neq 0$$

$$\Rightarrow M_{W,Z} \propto \langle \phi \rangle$$

Q. why scalar ( $v=0$ )?

why not  $\langle \psi \rangle \neq 0$  ?

$\langle A_\mu \rangle \neq 0$  ?

$\langle e \rangle \neq 0$  ?

$\langle \nu \rangle \neq 0$  ?

break Lorentz !

but Lorentz = sacred

⇓  
scalar: Lorentz!

Task of Weinberg:

$$M_W \neq 0, \quad M_Z \neq 0$$

$$M_A = 0$$

$$m_e \neq 0, \quad \omega_e \neq 0$$

$$e: \quad m_e \bar{e} e = m_e e^\dagger \gamma_0 e$$

$$e = e_L + e_R \equiv \underbrace{L e + R e}_{\uparrow}$$

$$L(R) = \frac{1 \pm \gamma_5}{2} \quad (\gamma_5^2 = 1)$$

$$\{\gamma_5, \gamma_\mu\} = 0$$

$$\Rightarrow [\gamma_5, \Sigma_{\mu\nu}] = 0$$

$\Rightarrow L, R =$  Lorentz invariant

$$m_e \bar{e} e = m_e (e_L^\dagger + e_R^\dagger) \gamma_0 (e_L + e_R)$$

$$= m_e e^\dagger (L + R) \gamma_0 (L + R) e$$

$$= m_e e^\dagger \gamma_0 R e + h.c.$$

$$= m_e (\bar{e}_L e_R + h.c.)$$





mass term :  $L \leftrightarrow R$

$\bar{e}_L e_R \leftarrow$  breaks  $SU(2)$   
FORBIDDEN

$\Downarrow$  instead

$$\mathcal{L}_Y^{(e)} = Y_e \bar{l}_L \Phi e_R + \text{h.c.}$$

$$\uparrow \quad \Phi \rightarrow U \Phi$$

$$l_L \rightarrow U l_L \quad \therefore U^\dagger U = 1$$

$$Y = 2 [Q - T_3]$$

$\Downarrow$

$$Y(e_R) = -2, \quad Y(l_L) = -1$$

$$\Rightarrow \boxed{Y(\Phi) = 1}$$

$$\begin{aligned} \Phi &= \begin{pmatrix} \varphi^+ \\ \varphi^0 \end{pmatrix} & Q &= T_3 + Y/2 \\ \Downarrow & & & \\ l &= \begin{pmatrix} \nu \\ e \end{pmatrix} \end{aligned}$$

$$\mathcal{L}_Y^{(e)} = Y_d \bar{l}_L \Phi d_R +$$

$$Y: \quad -1/3 + 1 - 2/3 = 0$$

$$+ \bar{l}_L \Phi^* u_R$$

$$Y: \quad -\frac{1}{3} - 1 + \frac{4}{3} = 0$$

$i\sqrt{2}$

$$q_L \rightarrow U q_L, \quad \Phi \rightarrow U \Phi, \quad u_R \rightarrow u_R$$

$$\bar{q}_L \Phi^* u_R \rightarrow \bar{q}_L \underbrace{U^\dagger U^*}_{=1} \Phi^* u_R$$

$$S = 0 : |\uparrow \downarrow - \downarrow \uparrow\rangle$$

$$D = \begin{pmatrix} \uparrow \\ \downarrow \end{pmatrix} \Leftrightarrow \Phi$$

$$\Phi_1, \Phi_2 \Leftrightarrow$$

$$\Phi_1^T \in \Phi_2 \quad \therefore \epsilon^T = -\epsilon$$

$$\Downarrow$$

$$\boxed{\Sigma = i\sigma_2} \quad \Leftrightarrow \quad \Sigma = \begin{pmatrix} 0 & 1 \\ -1 & 0 \end{pmatrix}$$

$$\begin{aligned} \Phi_1^T \in \Phi_2 &\rightarrow \Phi_1^T U^T \in U \Phi_2 \\ &= \Phi_1^T U^T i\sigma_2 U \Phi_2 \end{aligned}$$

$$U = e^{i\theta \sigma_2 / 2}$$

$$(\sigma_1, \sigma_3)^T = (\sigma_1, \sigma_3)$$

$$\sigma_2^T = -\sigma_2$$

$$[\sigma_2, \sigma_2] = 0$$

$$\{\sigma_2, \sigma_{1,3}\} = 0$$

$$\Rightarrow U^T i\sigma_2 U = e^{i\theta_i \sigma_i^T} i\sigma_2 e^{i\theta_i \frac{\sigma_i}{2}}$$

$$= i\sigma_2 \underbrace{e^{-i\theta_i \frac{\sigma_i}{2}} e^{i\theta_i \frac{\sigma_i}{2}}}_{1}$$

$$= i\sigma_2$$



$$\Phi_1^T i\sigma_2 \Phi_2 = \text{SU}(2) \text{ INV}$$

 $\Leftrightarrow$ 

$$\bar{\Phi} \rightarrow U \bar{\Phi}$$

$$i\sigma_2 \bar{\Phi}^* \rightarrow U i\sigma_2 \bar{\Phi}^*$$

PROVE!

$$T_a^\dagger = T_a$$

$$\Leftrightarrow T_a^* = T_a^T$$

$$S \not\sim B$$

$$\mathcal{L}_{\Phi} = \frac{1}{2} (D_\mu \Phi^\dagger) (D^\mu \Phi) - V(\Phi)$$

$$V(\Phi) = \frac{\lambda}{4} (\Phi^\dagger \Phi - v^2)^2$$

Why only  $\Phi^\dagger \Phi$ ?

why not:  $\bar{\Phi}^T i \sigma_2 \bar{\Phi}$ ? (=0)

-11-  $(\bar{\Phi}^T \sigma_i \bar{\Phi}) / (\bar{\Phi}^T \sigma_i \bar{\Phi})$ ?

-12-  $(\bar{\Phi}^T i \sigma_2 \sigma_i \bar{\Phi}) (\bar{\Phi}^T i \sigma_2 \sigma_i \bar{\Phi}^*)$ ?

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INV  $\Leftrightarrow$   $f(\bar{\Phi})$  ∴

$$f(\bar{\Phi}) = f(U\bar{\Phi})$$

$$U\bar{\Phi} = \begin{pmatrix} 0 \\ \phi \end{pmatrix} = \nabla$$

$$\bar{\Phi} = \begin{pmatrix} \phi_1 \\ \phi_2 \end{pmatrix} \Rightarrow U \therefore U\bar{\Phi}$$

$$\Rightarrow \boxed{ON E \text{ inv.}} = f(\phi)$$

$$\boxed{f(\bar{\Phi}) = \bar{\Phi}^+ \bar{\Phi}}$$

•  $\mathcal{M}_0 = \text{vacuum manifold}$   
 $= \{ \bar{\Phi}_0 : V(\bar{\Phi}_0) = V_{\min} = 0 \}$   
 $= \{ \bar{\Phi}_0^+ \bar{\Phi}_0 = v^2 \}$

$$\bar{\Phi} = \begin{pmatrix} \phi_1 \\ \phi_2 \end{pmatrix}, \quad \phi_i \in \mathbb{C}$$

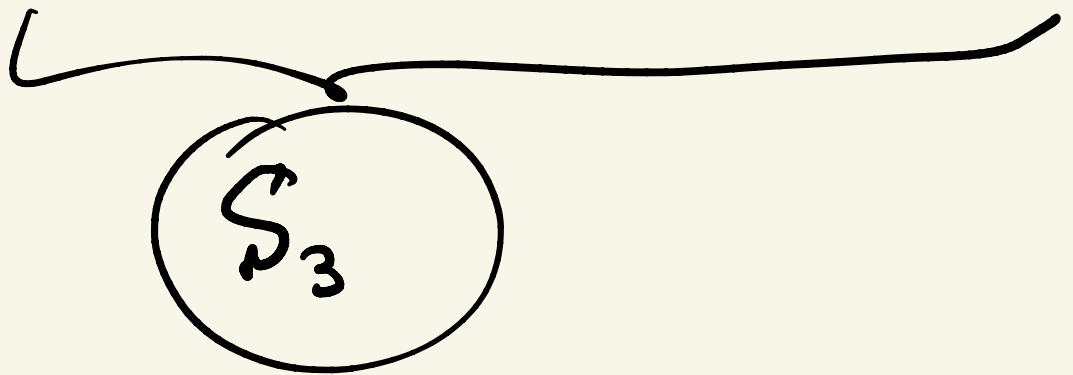
$$= \begin{pmatrix} R_1 + i R_2 \\ R_3 + i R_4 \end{pmatrix}$$



$$\Rightarrow \Phi^t \Phi = |\phi_1|^2 + |\phi_2|^2$$

$$= \sum_{i=1}^4 R_i^2$$

$$\Rightarrow \Phi_0^t \Phi_0^2 = \sum_{i=1}^4 R_{i0}^2 = v^2$$



$$\Phi_0^G = \begin{pmatrix} 0 \\ v \end{pmatrix}, \quad v \in \mathbb{R}$$



$$\Phi \rightarrow \Phi_0$$

$$\begin{aligned}
 \bullet \mathcal{L}_Y \rightarrow & \gamma_e \bar{l}_L \begin{pmatrix} 0 \\ \nu \end{pmatrix} e_R + \\
 & + \gamma_d \bar{q}_L \begin{pmatrix} 0 \\ \nu \end{pmatrix} d_R + \\
 & + \gamma_u \bar{q}_L \begin{pmatrix} \nu \\ 0 \end{pmatrix} u_R + h.c
 \end{aligned}$$

$$\Rightarrow \left[ \begin{array}{l} m_e = \gamma_e \nu \\ m_d = \gamma_d \nu \\ m_u = \gamma_u \nu \end{array} \right.$$

$$\bullet \frac{1}{2} (D_\mu \Phi_0)^+ (D^\mu \Phi_0) \quad Y(\Phi_0) = 1$$

$$D^\mu \Phi_0 = \left( -ig T_i A_\mu^i - ig' \frac{B_\mu}{2} \right) \Phi_0$$

a) charged glu. (A) ( $\sigma_1, \sigma_2$ )

$$D_\mu \Phi_0 \rightarrow -i \frac{g}{2} \begin{pmatrix} 0 & A_1 & -iA_2 \\ A_1 + iA_2 & 0 & \mu \end{pmatrix} \begin{pmatrix} 0 \\ \nu \end{pmatrix}$$

$$= -i \frac{g}{2} (A_1 - iA_2)_\mu \begin{pmatrix} \nu \\ 0 \end{pmatrix}$$

↓

$$(D_\mu \Phi_0)^\dagger D^\mu \Phi_0 = \frac{g^2}{4} \nu^2 (A_1^2 + A_2^2)$$

$$W^\pm = \frac{A_1 \mp iA_2}{\sqrt{2}}$$

$$M_W = \frac{g}{2} \nu$$

b) neutral glu. ( $\sigma_3, Y=1$ )

$$D_\mu \Phi_0 \rightarrow -i/2 \begin{pmatrix} gA_3 + g'B & 0 \\ 0 & -gA_3 + g'B \end{pmatrix} \Phi_0$$
$$\Phi_0 = \begin{pmatrix} 0 \\ v \end{pmatrix}$$

$$\Rightarrow |D_\mu \Phi_0|^2 = \frac{1}{4} v^2 (-gA_3 + g'B)^2$$

$$= \frac{1}{4} v^2 \left( \frac{gA_3 - g'B}{\sqrt{g^2 + g'^2}} \right)^2 (g^2 + g'^2)$$

$$\tan \theta_w \equiv g'/g$$

$\parallel$   
 $\neq$

$$M_Z^2 = \frac{(g^2 + g'^2)}{4} v^2$$

$$\Rightarrow Z = \cos\theta_w A_3 - \sin\theta_w B \quad \nearrow$$

$$A = \sin\theta_w A_3 + \cos\theta_w B$$

①

$$\therefore M_A = 0$$

②

$$M_Z \cos\theta_w = M_W$$

(1)

$CDF$   $\otimes$   $(D^0)$

Tevatron ( $E \approx TeV$ )

$\rightarrow$  Fermi

④ derivation from (1)

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Q1. why  $u_A = 0$ ?

Q2.  $\Phi_0^u = \begin{pmatrix} u \\ 0 \end{pmatrix} \rightarrow$

what happens?