

LMU GUT Course

Lecture IV

13/11/2020

Fall 2020



Non-Abelian Monopoles

Derived:

$$\vec{A}_s = -\frac{g}{4\pi v} \frac{1 + \cos\theta}{\sin\theta} \hat{\varphi}$$

$\theta = 0$
regular

Dirac:

(no string)

$\theta = 0$: $\vec{A}_s = \frac{1}{e} \nabla \alpha \quad (\gamma, \theta, \varphi)$

$$\alpha = \frac{g}{2\pi} \varphi$$

$$\frac{1}{e} \nabla \alpha = \frac{1}{e} U^\dagger \nabla U$$

$$U = e^{i\alpha}$$

$$U(2\pi) = U(0)$$

A must!

SU(2)

$$\psi \rightarrow e^{i T_a \theta_a} \psi$$

$$T_a = \sigma_a / 2 \quad (\theta_3, T_3)$$

$$\begin{aligned} \Rightarrow \psi(2\pi) &= e^{i \frac{\sigma_3}{2} 2\pi} = -1 \\ &= e^{i \sigma_3 \pi} \end{aligned}$$

$$V_{(1)} \otimes U = e^{i\alpha(x)} \textcircled{2}$$

$$U(2\pi) \neq 1$$

$$\xi = 1 \quad U = e^{i\alpha}$$

$$\xi = 0.1345 \quad U = e^{i\alpha \cdot 0.1345}$$

$$\xi = \frac{1}{3} \quad U = e^{i\alpha/3}$$

$$SU(2) \quad ! \quad Q = \bar{T}_3$$

$$U = e^{i\bar{T}_3 \theta}$$

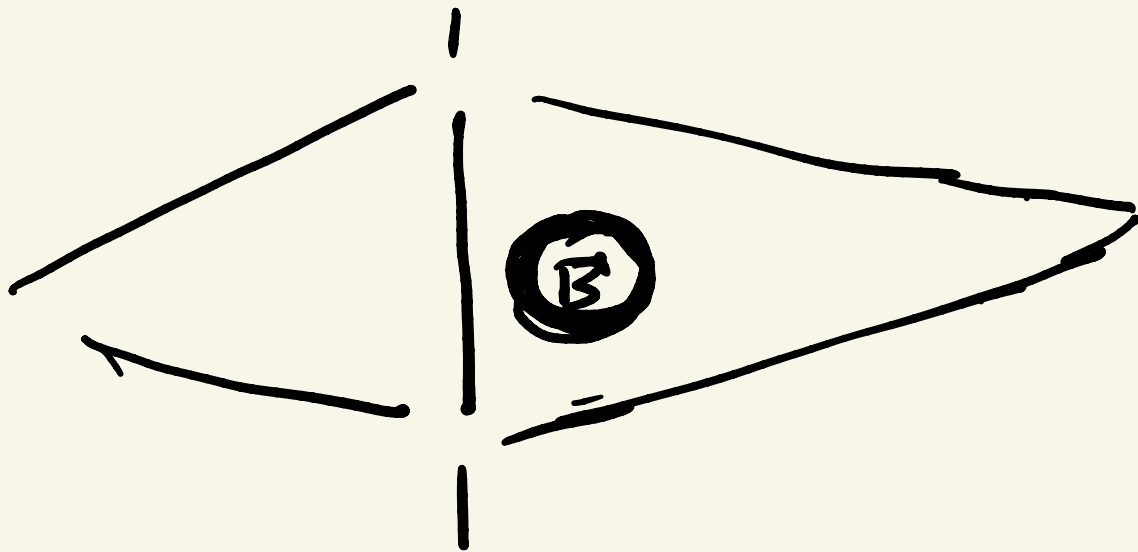
A hermit - Bohm

$$\vec{A} = 0 \quad \psi(x) \leftrightarrow \bar{V}(x)$$

$$\vec{A} \neq 0 \Rightarrow \psi'(x) = e^{i g(x)} \psi(x)$$

$$g(x) = e \int \vec{A} dx'$$

$$\bar{V} \rightarrow \bar{V} - ie \vec{A}$$



$$\begin{aligned} \text{difference} &= e \oint \vec{A} \cdot d\vec{x} \\ &= e \int \vec{B} \cdot d\vec{S} = e g \mu \end{aligned}$$

NOT physical

$$e g \mu = 2\pi n \quad (1)$$

if true \Rightarrow no way of
detecting "strings"

\vec{A}_N, \vec{A}_S

Wu-Yang

$$\bar{A}_N - \bar{A}_C = \frac{1}{e} \int \mathcal{D}d \quad (2)$$



pure gauge

$$\text{iff } U = e^{i\alpha} \quad (3)$$

$$U(2\pi) = U(0)$$

$$(\bar{A}_N, \bar{A}_S) = \text{potential}$$

$$e\oint = 2\pi n$$

$$\text{of } : \bar{\psi} \psi$$

$$\psi \rightarrow \psi' = e^{i\beta} \psi$$

$$\oint \bar{A} d\bar{L} = \text{quantized}$$

$$\oint \bar{A} d\bar{L} \approx u$$

$$\Rightarrow \bar{A} = 0 \Leftrightarrow \bar{A} \neq 0$$

$$U(1) \subseteq SU(2)$$

$$\underbrace{ew}_{(P)}$$

$$Q = T_3$$

$$SU(2) \rightarrow U(1)$$

$$SO(3) \longrightarrow U(1) = SO(2)$$

$$\phi_i \hat{x}_i = \vec{\Phi}, \quad \vec{\Phi} = \begin{pmatrix} \phi_1 \\ \phi_2 \\ \phi_3 \end{pmatrix}$$

$$\vec{\Phi} \rightarrow O \vec{\Phi} \quad O^T O = O O^T = 1 \\ \det O = 1$$

$$\phi_i \rightarrow \phi_i + \epsilon_{ijk} \theta_j \phi_k$$

$$O = e^{i L_a \phi_a}$$

$$a = 1, 2, 3$$

$$i = 1, 2, 3$$

$$L_a^\dagger = L_a, \quad T_a L_a = 0$$

$$(L_a)_{bc} = -i \epsilon_{abc}$$

SU(2)

$$U U^\dagger = 1, \quad \det U = 1$$

$$U = e^{i \theta_i T_i}$$

$$T_i = \frac{\sigma_i}{2}$$

(*)

$$\Sigma \rightarrow U \Sigma U^\dagger$$

$$T_\nu \Sigma = 0$$

$$\Sigma = \Sigma^\dagger$$

$$\Sigma = T_a \varphi_a$$

$$\Sigma \rightarrow \hat{U} \Sigma = e^{i\theta_a \hat{T}_a}$$

$$\hat{T} \Sigma = [\tau, \Sigma]$$

(**)

$$V = \frac{\lambda}{4} (\varphi_i \varphi_i - v^2)^2 \quad (4)$$

φ^2

$$V = \frac{\lambda}{4} (2T_\nu \Sigma^2 - v^2)^2 \quad (5)$$

$$T_\nu \Sigma^2 = \varphi_i \varphi_j T_\nu T_i T_j = \varphi_i \varphi_j \frac{1}{2} \delta_{ij}$$

$$= \frac{1}{2} \varphi^2$$

$$U = e^{i\theta_a T_a}$$

$$\Sigma \rightarrow e^{i\theta_a T_a} \Sigma e^{-i\theta_a T_a}$$

$$= (1 + i\theta_a T_a + \dots) \Sigma (1 - i\theta_a T_a)$$

$$= \Sigma + i [T_a, \Sigma] \theta_a$$

$$= \Sigma + i \hat{T}_a \theta_a \Sigma = \hat{U} \Sigma$$

$$\hat{T}_a \Sigma = [T_a, \Sigma]$$

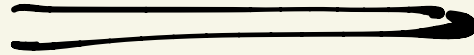
$$\Sigma \rightarrow \Sigma + i [T_a, T_b] \psi_b \theta_a + \dots$$

$$= \Sigma + i \epsilon_{abc} T_c \psi_b \theta_a + \dots$$

$$= T_c [\psi_c + i \epsilon_{abc} \theta_a \psi_b + \dots]$$

$$\psi_c' = \psi_c + i \epsilon_{abc} \psi_b \theta_a$$

checks ✓



$$V = \frac{\lambda}{4} (2T, \Sigma^2 - \varrho^2)^2$$

$$M_0 = \{ \text{vacuum manifold} \}$$

$$= \{ \tau_0, V = V_{\text{min}} \}$$

$$= \{ \Sigma^0, 2T, \tau_0^2 = \varrho^2 \}$$



$$i=1,2,3 \quad \psi_0^i \psi_0^i = \varrho^2$$

$$= S_2$$

$$\Sigma = T_a \psi_a$$

Vectors. $\psi_0^3 = \psi, \quad \psi_0^1 = \psi_0^2 = 0$

$\Sigma^0 = L_3 \psi$ $SO(3)$
 $Q = L_3$ $\begin{pmatrix} 0 \\ 0 \\ \psi \end{pmatrix}$

$$(L_3)_{ij} = -i\epsilon_{3ij}$$

$$L_3 = \begin{pmatrix} 0 & -i & 0 \\ i & 0 & 0 \\ 0 & 0 & 0 \end{pmatrix}$$

photon

$$A = A_3$$

$$D_\mu = \partial_\mu - ig L_3 \hat{t}^a \\ = \partial_\mu - ig (L_3 A_\mu^3 + \dots)$$

$$D_\mu \Phi_0 = -ig L_3 A_\mu^3 \Phi_0 + \dots$$

$$L_3 \Phi_0 = 0$$

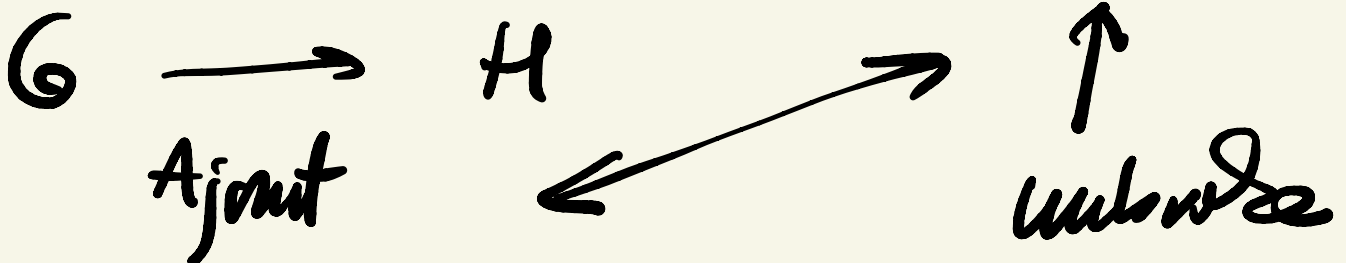
$$\Rightarrow A_3^3 = \text{massless}$$

$$\boxed{SU(2)} \quad \hat{T}_a \Sigma = [T_a, \Sigma] \quad \text{Pauli'}$$

$$\Sigma_0 = v T_3 \Rightarrow Q = T_3$$

$$\Rightarrow Q = \frac{\Sigma_0}{v}$$

$$[Q, \Sigma_0] = 0$$



$$g = e$$

↓

$$\textcircled{e} \left[\begin{array}{l} A = A_3 \\ u_A = 0 \end{array} \right.$$

$$M_{A_1} = M_{A_2} = e v$$

$$W^\pm = \frac{A_1 \mp i A_2}{\sqrt{2}}$$

check

$$T_\nu (D_\mu \bar{\psi})^\dagger (D^\mu \psi)$$

$$A \rightarrow F_{\mu\nu} = \partial_\mu A_\nu - \partial_\nu A_\mu \neq F_{\mu\nu}^3$$

$$F_{\mu\nu}^a = \partial_\mu A_\nu^a - \partial_\nu A_\mu^a + g \epsilon_{abc} A_\mu^b A_\nu^c$$

$$F_{\mu\nu}^3 = \underbrace{(\partial_\mu A_\nu - \partial_\nu A_\mu)}_{F_{\mu\nu}} + g (A_\mu^1 A_\nu^2 - A_\mu^2 A_\nu^1)$$

$$\Sigma_0 = T_a \phi_a^0$$

$$Q = \frac{\Sigma_0}{v} \quad [Q, \Sigma_0] = 0$$

↑ unbroken

$$\boxed{SU(2) \xrightarrow{\Sigma} U(1)}$$

→ \odot → denote vacuum value of ...

$$\boxed{\Sigma_0 = \text{vev of } \Sigma}$$

$$E_0 = 0 \quad (\text{vacuum})$$

$$V = \frac{\lambda}{4} (2T_v \Sigma^2 - v^2)^2$$

$$T_{\nu} \Sigma_0^2 = \omega^2/2$$

Can I find a finite energy
solution (classical) of
equation of motion?

STATIC

$$D_0 \Phi_{cl} = 0$$

$$D_0 \Psi_{cl} = 0$$

$$\mathcal{L} = \frac{1}{2} |D_{\mu} \Sigma|^2 - V(\Sigma) \\ - \frac{1}{4} F_{\mu\nu}^a F^{\mu\nu a}$$

$$\begin{aligned}
 & \downarrow \rightarrow 0 \text{ at } \infty \\
 E = \int \mathcal{H} = & \left[\frac{1}{2} |\mathbf{D}_i \Sigma|^2 + V(\Sigma) \right. \\
 & \left. + \sum_a \frac{1}{2} (\vec{E}_a^2 + \vec{B}_a^2) \right] dV \\
 & = \text{finite} \quad (< \infty)
 \end{aligned}$$

$$\Rightarrow \boxed{\mathcal{H}(\Sigma) \xrightarrow{\infty} 0 \text{ (} \mathcal{M}_0 \text{)}}$$

$$\Rightarrow \therefore V(\Sigma_0) = 0 \\
 \Sigma_0 \subseteq \mathcal{M}_0$$

$$\boxed{\text{map: } \mathcal{M}_0 \rightarrow \mathcal{M}_0}$$

$$M_0 = S_2 \Rightarrow M_\infty = S_2$$

spherically symmetric Σ_0

$$M_\infty (r \rightarrow \infty) (\theta, \varphi)$$

$\Sigma_0 = vT_3$ exosphere \rightarrow
vacuum at M_∞

$M_\infty = S_2 \rightarrow 1 \text{ part}$
Trivial

non-trivial map

$$\Sigma_\infty = T_a \varphi_a^\infty \text{ (value at } \infty)$$

$$Q_a = \frac{\Sigma_{a0}}{v} \quad (\text{canceled})$$

$$\hat{Q}_{a0} \Sigma_{a0} = [Q_{a0} \Sigma_{a0}] = 0$$

Vacuum:

$$Q = v T_3 \Rightarrow A = A_3$$

$$Q = \sum_{a=1}^3 \psi_a^\dagger T_a \Rightarrow A = \sum_{a=1}^3 A_a \frac{\psi_a^\dagger}{v} (?)$$

$$\sum_{a=1}^3 \psi_a^\dagger \psi_a = v^2$$

↑ canceled

$$\left\{ \begin{array}{l} \psi_a^\dagger = v \delta_{a3} \\ \Rightarrow A = A_3 \end{array} \right.$$

$$\left\{ \begin{array}{l} \psi_a^\dagger = v \delta_{a1} \\ \Rightarrow A = A_1 \end{array} \right.$$

$$T_\nu \frac{1}{2} |\partial_\mu \Sigma|^2 \rightarrow A_a A_b \underline{M}_{ab}^2$$

f-ct of $(\Sigma_0 = T_a \phi_a^0)$

$$(\underline{M}^2 A^0)_a = M_{ab}^2 A_b^0 = 0$$

photo



derive M_{ab}^2 , prove

$$\text{that } A^0 = A_a \frac{\phi_a^0}{c} \Rightarrow$$

massless

$$Q = T_a \frac{\phi_a^0}{c} \Leftrightarrow A^0 = A_a \frac{\phi_a^0}{c}$$

III
A

$$Q = T_a \frac{\psi_a^\infty}{v} \iff A = \frac{A_a \psi_a^\infty}{v}$$

$$F_{\mu\nu} = F_{\mu\nu}^a \frac{\psi_a^\infty}{v}$$

$$\psi_a^\infty = v \frac{x^a}{r}$$

hedging

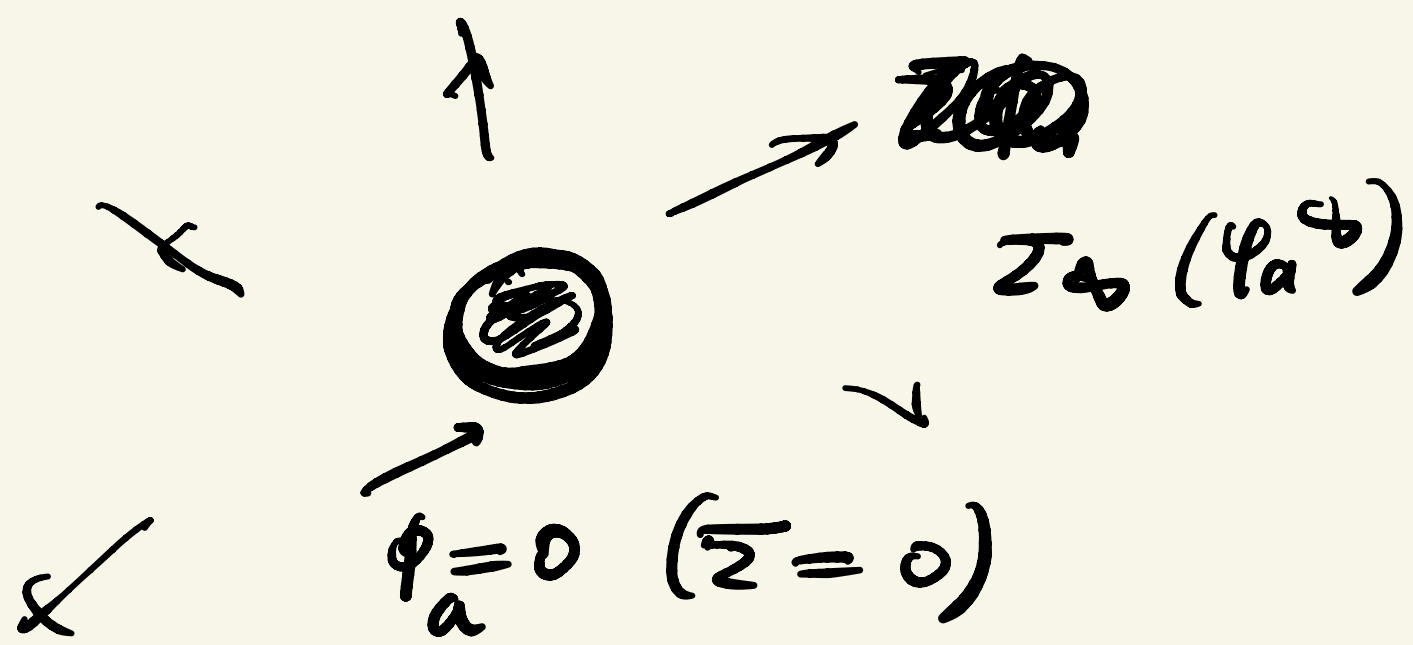
$$A = A_a \frac{x^a}{v}$$

$$F_{\mu\nu} = F_{\mu\nu}^a \frac{x^a}{v}$$

$$r \rightarrow 0 \quad ? \quad \cancel{\varphi_a = \varphi_a^\infty = a \frac{x_a}{r} ??}$$

$$\varphi_a^a \text{ monopole} \begin{matrix} \xrightarrow{\quad} \varphi_a^\infty (r \rightarrow \infty) \\ \xrightarrow{\quad} 0 (r \rightarrow 0) \end{matrix}$$

$$\boxed{V(r=0) = \frac{1}{a} a^4}$$



$$E = \text{finite} \Rightarrow D_i \quad Z_\infty = 0$$

$$\Downarrow$$

$$D_i \Sigma_a = \partial_i \Sigma_a g | [T_a, \Sigma] A^a_i |$$

$$\underline{\underline{\hspace{10em}}} = 0$$

$$g = e$$

$$A_i^a(\alpha) = \cancel{x \delta_{ia} \frac{1}{r}} + \cancel{\Sigma_{ij}^a \frac{x_i}{r^2 j}} \quad (\pm)$$

$$\Sigma_a = T_i \frac{x_i}{r} \cancel{e} + \cancel{\frac{x_i x_a}{r^2} z}$$

$$F_{ij}^a(\alpha) = \Sigma a_{ij} \frac{1}{r^2} \frac{1}{g}$$

$$F_{ij}^a = \Sigma a_{ij} \frac{1}{g r^2} \frac{x_a}{r}$$

$$= \Sigma i_j u B_u$$

$$\Rightarrow B_k^a = \frac{1}{g r^2} \frac{x_k}{r}$$

$$= \frac{1}{ev^2} \frac{x_u}{r}$$

$$g_m \propto \frac{1}{e}$$

$$g = e$$

$$M_m \propto v/g$$