

LMU GUT Course

Lecture XVI

8/1/2021

LMU

Winter 2021



SU(5): unification

and phenomenology

- matter = fermions =
= quarks + leptons

$$5_F = \begin{pmatrix} d \\ \dots \\ e^c \\ -\nu^c \end{pmatrix}_R \left. \begin{array}{l} \} SU(3) \rightarrow \text{gluons} \\ \} SU(2) \rightarrow W^\pm, Z, \gamma \end{array} \right\}$$

$$10_F = \begin{pmatrix} 0 & u^c & \dots & u & d \\ & 0 & \dots & \dots & \dots \\ & & 0 & \dots & \dots \\ \dots & \dots & \dots & 0 & e^c \\ & & & & 0 \end{pmatrix}_L \left. \begin{array}{l} \} SU(3) \\ \} SU(2) \end{array} \right\}$$

$$\Rightarrow Q_u = 0, \quad Q_e = 3 Q_d$$

• groupe bosons = 24

$$24 = 12 + 12$$

SM (glue, w, z, γ)

new

$$12 = 6 + \bar{6} = (x, y)^a + (\bar{x}, \bar{y})^a$$

$$M_x = M_y + O(M_w)$$

$$\mathcal{L}_f = i \bar{\psi}_F \gamma^\mu D_\mu \psi_F + \dots$$

$$D_\mu = \partial_\mu - ig T_i A_\mu^i$$

$$i = 1, 2, \dots, 24$$

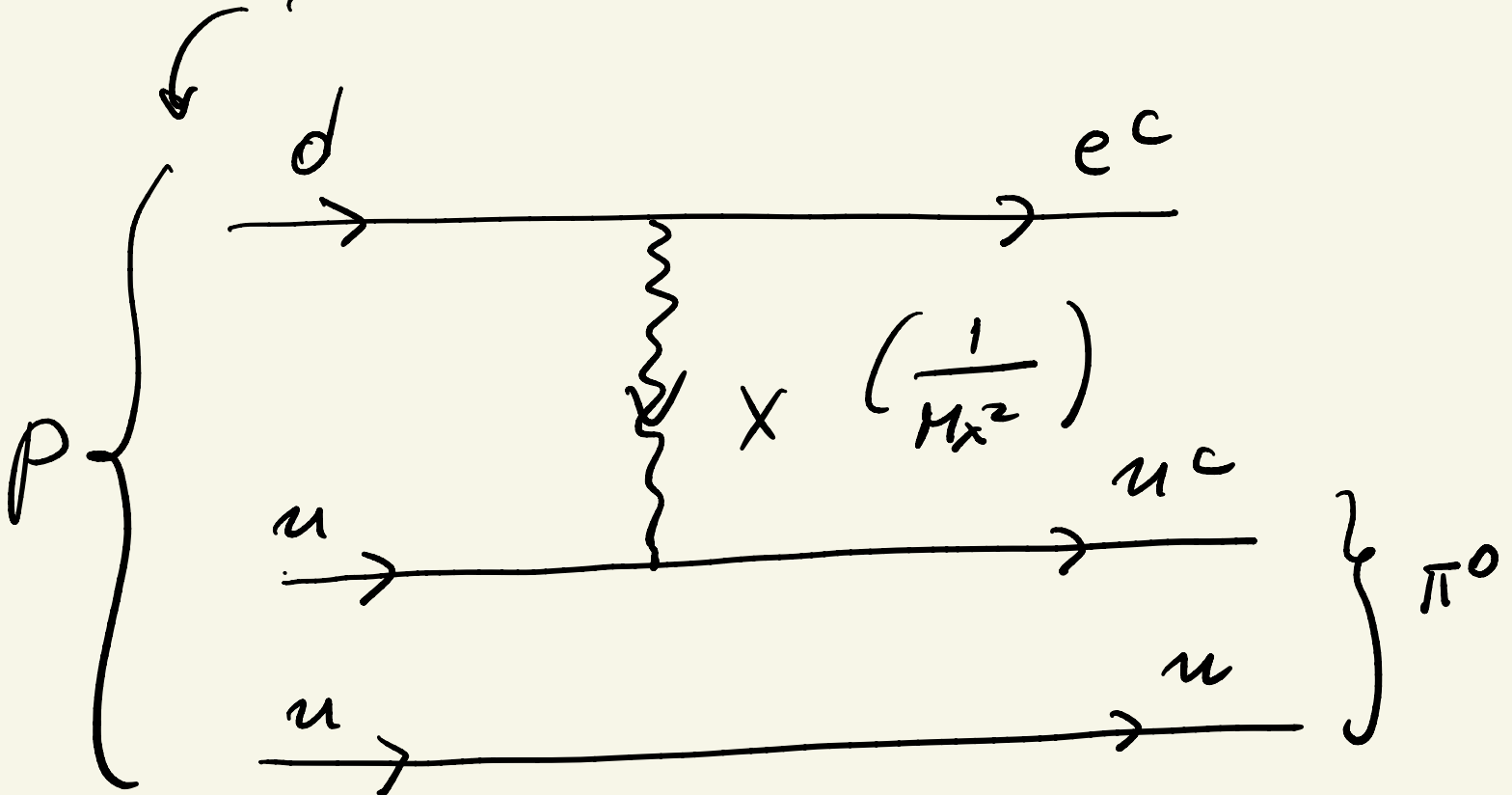
$$10_F = (u^c, u, d, e^c)$$

$$M_L^c \equiv C \bar{u}_R^T$$

$$(-4/3) \chi_{\mu} \left[\bar{u}_L^c \gamma^{\mu} u_L + \bar{d}_L^c \gamma^{\mu} e_L^c \right]$$

$$Q : \quad 2/3 \quad 2/3 \quad 1/3 \quad 1$$

$$\chi_{\mu} \bar{e}_L^c \gamma^{\mu} d_L + \dots$$



$$\boxed{P \rightarrow e^c + \pi^0}$$

$$\tau_P > 10^{34} \text{ yr (SK)}$$

$\Rightarrow X$ very heavy

$$\Gamma_p \propto g^4 M_x^{-4} m_p^5$$

$$\Gamma_\mu \propto g^4 M_w^{-4} m_\mu^5$$

$$\Rightarrow \Gamma_p / \Gamma_\mu = \left(\frac{m_p}{m_\mu} \right)^5 \left(\frac{M_w}{M_x} \right)^4$$

$$\Rightarrow \bar{\tau}_p / \bar{\tau}_\mu = \left(\frac{M_x}{M_w} \right)^4 \left(\frac{m_\mu}{m_p} \right)^4$$

$$= 10^{-5} \left(\frac{M_x}{M_w} \right)^4$$

$$\bar{\tau}_\mu \approx 10^{-6} \text{ sec}$$

$$\Rightarrow \tau_p = 10^{-11} \text{ sec} \left(\frac{M_x}{M_W} \right)^4 \gtrsim 10^{41} \text{ sec}$$

$$\left(M_x / M_W \right)^4 \gtrsim 10^{52}$$

$$\Rightarrow M_x / M_W \gtrsim 10^{13}$$

$$M_x > 10^{15} \text{ GeV}$$

\$U(5)\$: unified

unification scale M_U

$$M_U \gtrsim M_x > 10^{15} \text{ GeV}$$

$$10^{15} \text{ GeV} < M_U \ll M_{pl}$$

M_{pe} = scale when gravity
gets strong

$$\alpha(\text{gravity}) \propto G_N E_1 E_2$$

$$\approx \frac{1}{M_p^2} E_1 E_2$$

$$M_p \approx 10^{18} \text{ GeV} - 10^{19} \text{ GeV}$$

$$10^{15} \text{ GeV} < M_U \ll 10^{18} \text{ GeV}$$

$\alpha_1, \alpha_2, \alpha_3$ - couplings of
 $U(1), SU(2), SU(3)$

• all 3 must unify at M_U

= miracle

- in a small window:
 $10^{15} - 10^{17}$ GeV

- Symmetry breaking (I)

$$SU(5)$$

$$\Sigma \equiv 24_H$$

$$\downarrow M_\nu \approx M_X$$

$$SU(3) \times SU(2) \times U(1)$$

$$\Sigma \rightarrow U \Sigma U^\dagger$$

Adjoint

$$\Sigma^+ = \Sigma$$

$$T_3 \Sigma = 0$$

$$U \equiv e^{i\theta_i T_i} = 1 + i\theta_i T_i + \dots$$

$$\Sigma \rightarrow \Sigma + i [\tau_i, \Sigma] \theta_i$$

\Downarrow

$$D_\mu \Sigma = \partial_\mu - ig [\tau_i, \Sigma] A_\mu^i$$

Reminder: $\psi \rightarrow e^{iQ\alpha} \psi, U(1)$

$$D_\mu = \partial_\mu - ig A_\mu Q$$

• $D \rightarrow e^{iA_a \sigma_a / 2} D$ SU(2)

\Downarrow

$$D_\mu = \partial_\mu - ig \frac{\sigma_a}{2} A_\mu^a$$

$$\langle Z \rangle \rightarrow U \langle Z \rangle U^\dagger$$

$$= \text{diag} (1, 1, 1, -3/2, -3/2) \sqrt{x}$$

$$v_x = ?$$

$$\cdot T_\nu D_\nu \langle \Sigma \rangle D^\mu \langle \Sigma \rangle \Rightarrow$$

$$M_x = -M_y = \frac{5}{2} g v_x$$

$$M_{gluon} = M_w = M_z = M_\gamma = 0$$

Stage 1 of sym. breaking

$$(II) \quad SU(3) \times SU(2) \times U(1)$$

$$\downarrow \langle S_H = \Phi \rangle = M_w$$

$$SU(3) \times U(1)_{em}$$

$$S_H = \Phi = \begin{pmatrix} T^+ \\ \dots \\ \phi^+ \\ \phi^0 \end{pmatrix} \leftarrow \text{new} \neq \Phi$$

$$V = V_{\Sigma} + V_{\Phi} +$$

$$+ \alpha (T_V \Sigma^2) \bar{\Phi}^+ \bar{\Phi} + \left[\lambda \bar{\Phi}^+ \Sigma^2 \bar{\Phi} \right]$$

irrelevant

crucial

$$\boxed{\Sigma \rightarrow \langle \Sigma \rangle}$$

$$\alpha: (T_V \langle \Sigma \rangle^2) \bar{\psi}_H^+ \bar{\psi}_H \rightarrow \delta m_S^2 \bar{\psi}_H^+ \bar{\psi}_H$$

(change in V_{Φ})

$$V_S = \mu^2 \bar{\psi}_H^+ \bar{\psi}_H + \lambda (\bar{\psi}_H^+ \bar{\psi}_H)^2$$

$$\mu^2 \rightarrow \mu^2 + \alpha \text{Tr}(\Sigma)^2 = \mu_{\text{new}}^2$$

$$\beta: \quad \bar{\psi}_H^+ (\Sigma^2) \psi_H$$

$$= (\bar{T}^+ \bar{\Phi}^+) v_x^2 (1, 1, 1, g/4, g/4) \begin{pmatrix} T \\ \Phi \end{pmatrix}$$

↑
doublet (SM)

$$= \beta v_x^2 [T^+ T + g/4 \bar{\Phi}^+ \Phi]$$

⇓

$$m_T^2 = m_\Sigma^2 + \beta v_x^2 \approx (10^{15} \text{ GeV})^2$$

$$m_\Phi^2 = m_\Sigma^2 + g/4 \beta v_x^2 \approx (100 \text{ GeV})^2$$

• $T = ?$ Who is this?

} Symmetry arguments

• weak singlet, $SU(3)$ triplet

$$Y = -2/3 \quad (Q_T = Y/2 = -1/3)$$

• Lorentz scalar

⊗ $T \left[\begin{array}{c} \text{? coeff.} \\ u_R^T C d_R^\gamma \epsilon_{\alpha\beta\gamma} \\ \text{+ 1/3} \end{array} + \underbrace{u_L^T C d_L^\gamma \epsilon_{\alpha\beta\gamma}}_{\text{singlet?}} \right]$

Q: $-1/3$ $+1/3$

$$f_1^T C f_2 = \text{Lorentz singlet}$$

$$q_L \equiv \begin{pmatrix} u \\ d \end{pmatrix}_L \Rightarrow q_L^T C i \sigma_2 q_L \quad \rightarrow$$

// $SU(2)$ singlet

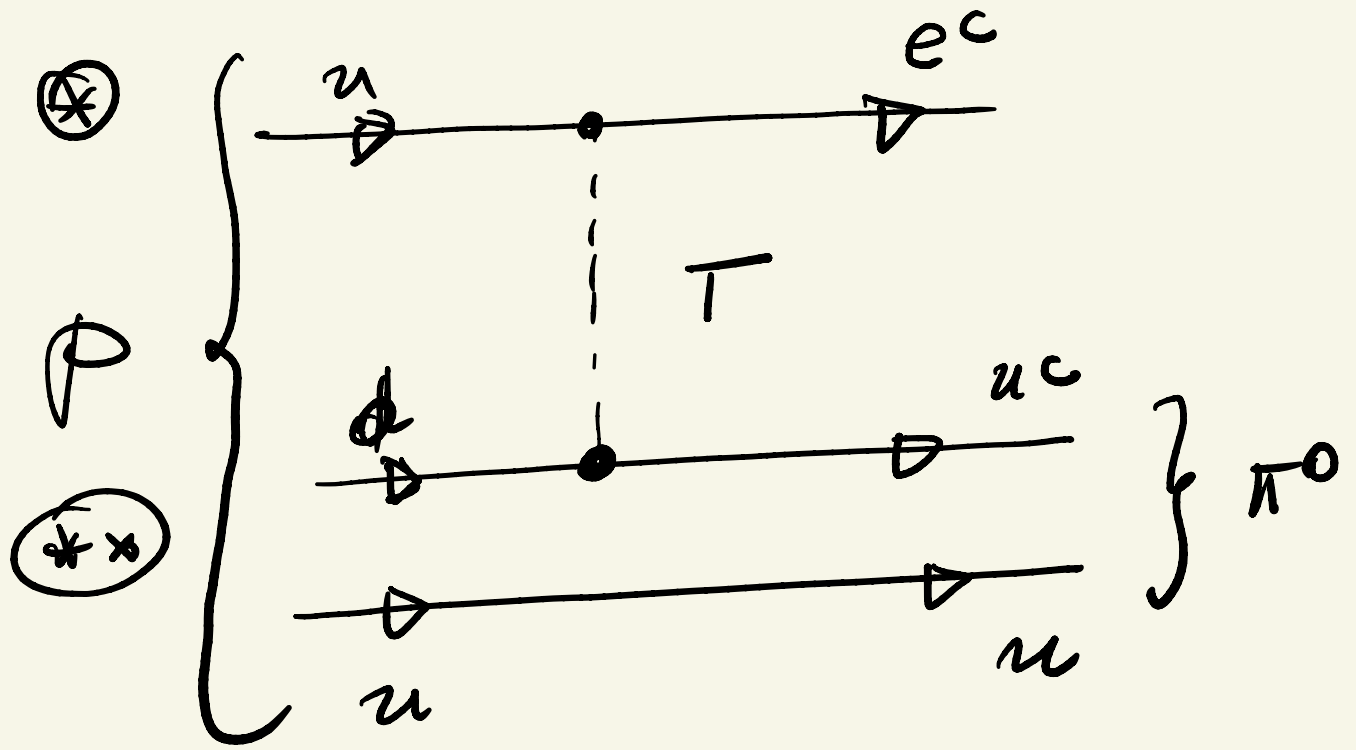
$$u_L^T C d_L$$

quark - lepton

coeff. ?

$$(1/3) \underbrace{T^*}_{-1/3} \underbrace{u e}_{-1/3} = T^*_\alpha \left[u_R^{\alpha T} C e_R \otimes \right. \\ \left. + u_L^{\alpha T} C e_L \otimes d_L^{\alpha T} C \nu_L \right]$$

$$l_L \equiv \begin{pmatrix} \nu \\ e \end{pmatrix}_L \Rightarrow q_L^T C i \sigma_2 l_L \\ = u_L^T C e_L \otimes d_L^T C \nu_L$$



$\Rightarrow m_T > 10^{12} \text{ GeV}$
 (unit.) $\Rightarrow m_T \gtrsim M_X$

• mass terms (T, Φ)

\Downarrow
 T: $m_S^2 + \rho v_X^2 \approx (10^{15} \text{ GeV})^2$ (1)
 Φ $m_S^2 + 9/4 \rho v_X^2 = 0$ (2)

$$m_S^2 \approx -9/4 \beta v_x^2$$

$$\Rightarrow \boxed{m_T^2 = -5/4 \beta v_x^2}$$

$(\beta < 0)$

T is heavy

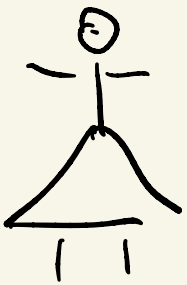
price: fine-tuning (FT)

$$\left. \begin{array}{l} m_{15} \approx 10^{15} \text{ GeV} \\ \beta v_x \approx 10^{15} \text{ GeV} \end{array} \right\} \boxed{\begin{array}{l} m_S^2 + \\ + 9/4 \beta v_x^2 = 0 \end{array}}$$

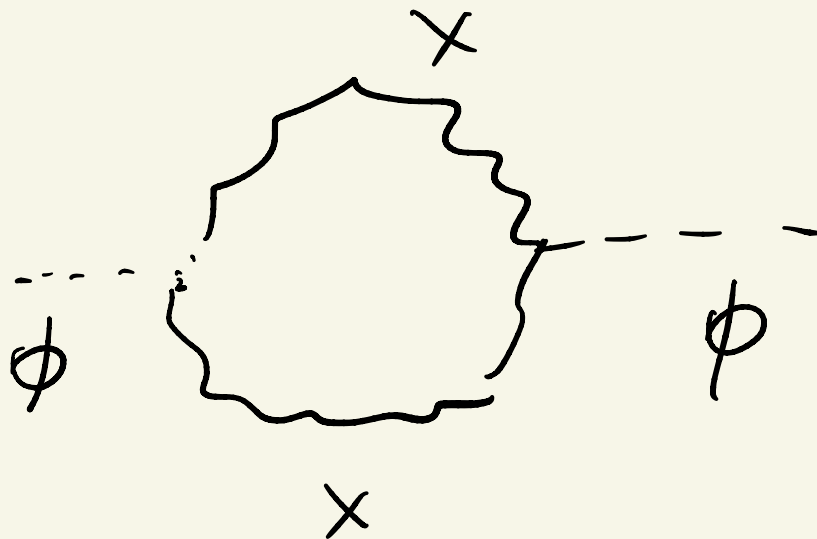
Is this a
problem??

Alice: FT is impossible
to maintain!

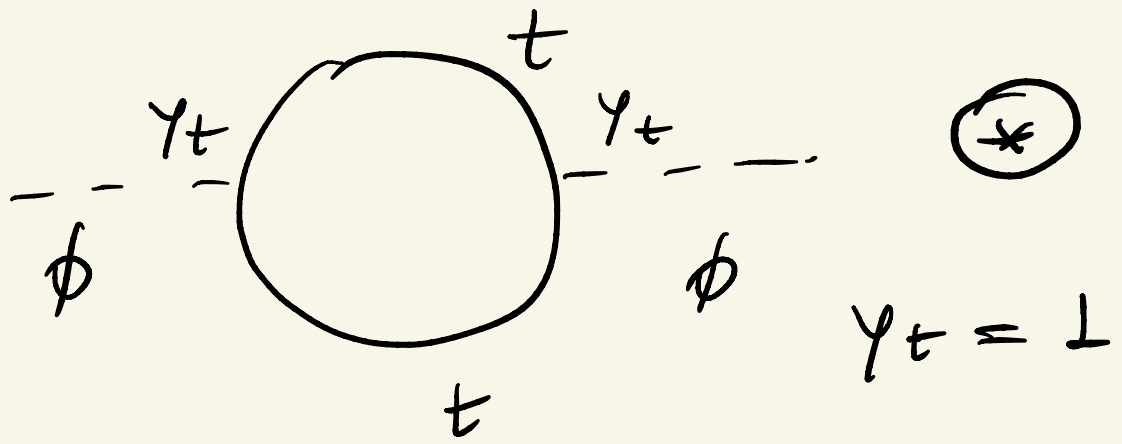
Bob: Just do it!



tree-level?



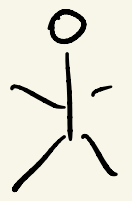
$$M \phi^2 \propto \frac{\alpha}{\pi} M_x^2$$



$$\gamma_t = 1$$

$$\Rightarrow m_\phi^2 \approx \frac{\gamma_t^2}{16\pi^2} \left(\Lambda^2 + m_t^2 \right)$$

$$\approx M_x$$



Bob :

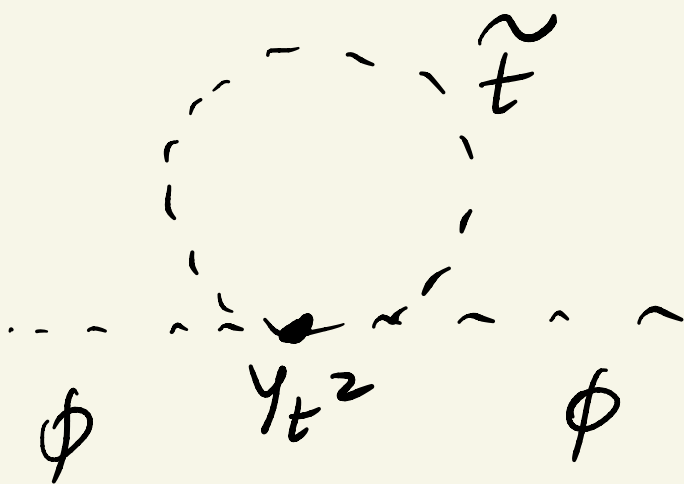
do it at t-loop!

xx

\tilde{t} (scalar) colored triplet

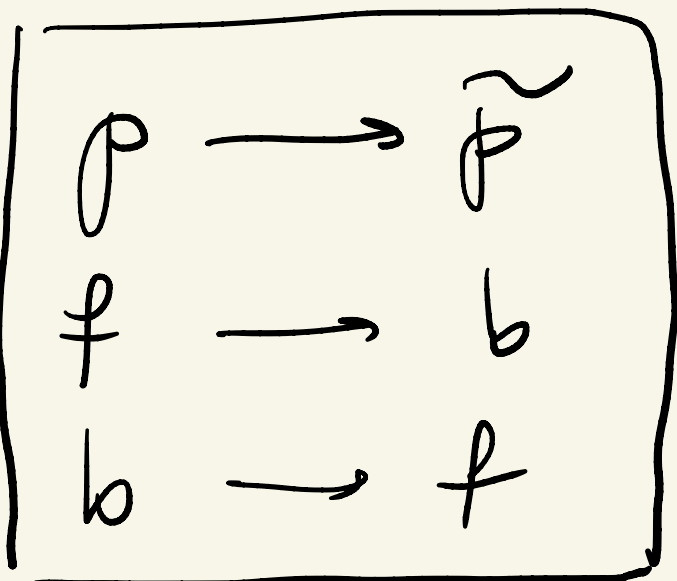
scalar top = stop

same couplings!



$$m_{\phi}^2 = \frac{\gamma_{\tilde{\tau}^2}}{16\pi^2} \left(\cancel{A^2 + m_{\tilde{\tau}}^2} - \cancel{A^2 + m_{\tilde{t}}^2} \right)$$

$m_{\tilde{E}} \leq \text{TeV} \Rightarrow$ Higgs
mass protected



1970's
Super Symmetry

$$e \longrightarrow \tilde{e} \text{ (selectron)}$$

$$m_{\tilde{e}} \approx 100 \text{ GeV}$$

(broken) supersymmetry

$\Lambda_{SS} \approx$ scale of $\tilde{S} \tilde{S}$
breaking

= masses of superpartners
($\leq 1 \text{ TeV}$)

- beautiful protective mechanism, but
- a collection of models

$$M_{\tilde{E}} > m_{\tilde{e}} < M_{\text{GUT}} !?$$

~ 100 new parameters

- $M_{\text{particles}} = \text{the same}$
 $= \text{at same large scale}$

Super symmetry =

the same FT

but

do it at tree level and you
are free forever



Can I find a rationale to
Higgs = light?

Idea: Higgs = (pseudo)
= Nambu - Goldstone boson

• Anselm ---

• Bertolini, Dvali, ---

SU(6)

⇒ change physics at the
unif. scale (M_0)

⇒ this theory \Leftrightarrow usual
FT

$$W \longrightarrow \tilde{W} \text{ (winos)}$$

$$Z \longrightarrow \tilde{Z} \text{ (zinos)}$$

$$\gamma \longrightarrow \tilde{\gamma} \text{ (photino)}$$

$$f \longrightarrow \tilde{f} \text{ (sfermions)}$$

$$h \longrightarrow \tilde{h} \text{ (Higgsinos)}$$

(Higgs)

2 Higgs
doublets

FT — because

$$m_T \gg m_\Phi$$

(T mediates p decay)

\Rightarrow split m_T from $l\ell\phi$

FT - independent of
Super Symmetry

$$M_x > 10^{15} \text{ GeV}$$

$$M_y > 10^{15} \text{ GeV}$$

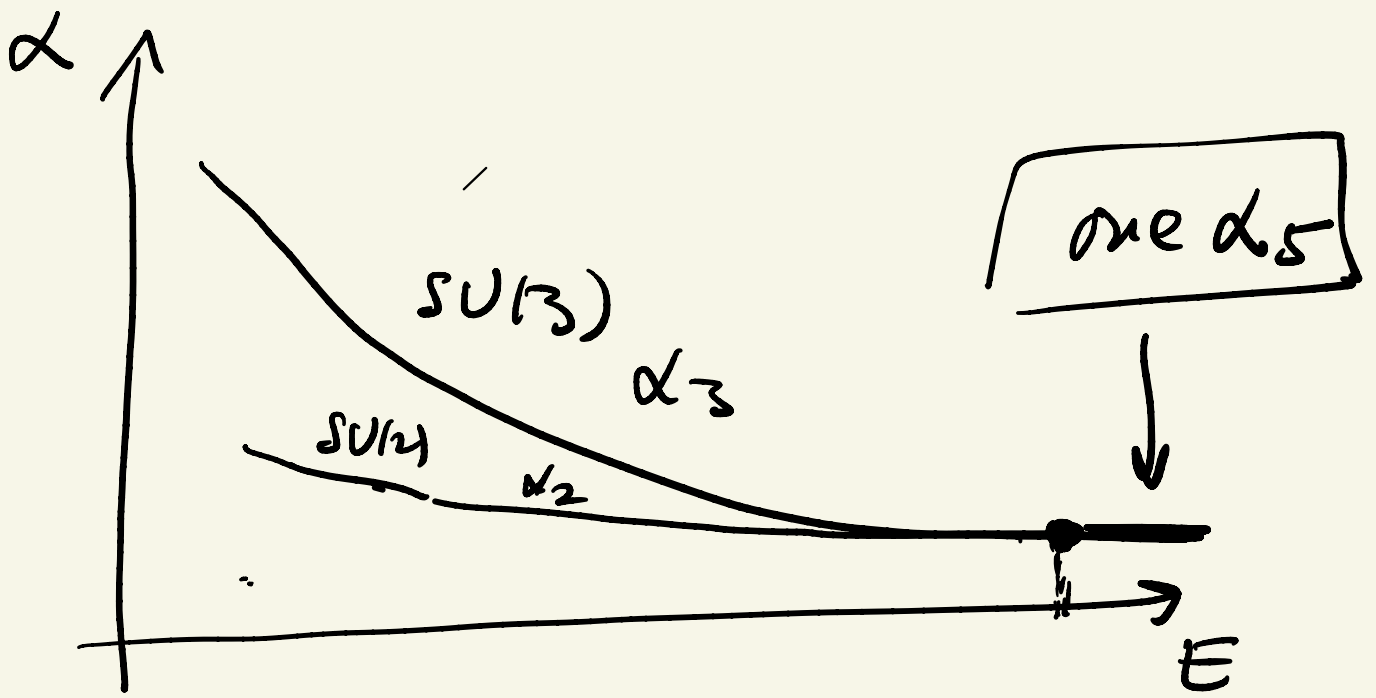
$$\left. \begin{array}{l} M_x = M_y \\ \text{SU}(2) \text{ sym.} \end{array} \right\}$$

$$M_x \propto g v_x$$

$$M_y \propto g v_x$$

• couplings "run" \Leftrightarrow

$$\alpha = \alpha(E)$$



$$M_U \approx M_X$$

$$M_U \approx M_X$$

- $$m_f = m_f^{\text{tree}} \left[1 + \frac{\alpha}{\pi} \ln \frac{\Lambda}{m} \right]$$

"small"

$$m_f^{\text{tree}} = 0 \Rightarrow \text{divergent log.}$$

$$\Rightarrow m_f = 0$$



Super Symmetry : mixes bosons
 \sim fermions

\Rightarrow Higgs \sim Higgsinos

\Rightarrow every body protected

• GUT \Rightarrow FT

(i) $S^1 S^1 \rightarrow$ NG picture

large scale $\Lambda \approx M_x$

$$\Rightarrow m_\phi^2 \approx m_{\phi 0}^2 + \frac{\alpha}{\pi} \Lambda^2$$

$\underbrace{\hspace{10em}}$
large correction

• no GUT

$\Lambda \sim M_{\text{Pl}}$ at $\approx M_{\text{Pl}}$

$\Lambda \approx 1-10 \text{ TeV}$

(ii)

unification of couplings

• ordinary $SU(5) \Rightarrow$

no unif.

• $S \& SU(5) \Rightarrow$
unif.