

Exercises on General Relativity TVI TMP-TC1

Sample solution to problem set 4 of PS 13

Solution

(i) The corresponding metric is given by

$$g_{\mu\nu} = \text{diag} \left(-1, \frac{a^2}{1-kr^2}, a^2 r^2, a^2 r^2 \sin^2 \theta \right) \quad (1)$$

$$g^{\mu\nu} = \text{diag} \left(-1, \frac{1-kr^2}{a^2}, \frac{1}{a^2 r^2}, \frac{1}{a^2 r^2 \sin^2 \theta} \right). \quad (2)$$

Christoffel symbols; with help of an algebraic program we compute all non-vanishing components:

$$\Gamma_{rr}^t = \frac{\dot{a}a}{1-kr^2}, \quad \Gamma_{\theta\theta}^t = r^2 a \dot{a}, \quad \Gamma_{\phi\phi}^t = r^2 a \dot{a} \sin^2 \theta \quad (3)$$

$$\Gamma_{tr}^r = \frac{\dot{a}}{a}, \quad \Gamma_{rr}^r = \frac{kr}{1-kr^2}, \quad \Gamma_{\theta\theta}^r = -r(1-kr^2) \quad (4)$$

$$\Gamma_{\phi\phi}^r = -r(1-kr^2) \sin^2 \theta, \quad \Gamma_{r\theta}^\theta = \frac{\dot{a}}{a}, \quad \Gamma_{r\theta}^\theta = \frac{1}{r} \quad (5)$$

$$\Gamma_{\phi\phi}^\theta = -\sin \theta \cos \theta, \quad \Gamma_{t\phi}^\phi = \frac{\dot{a}}{a}, \quad \Gamma_{r\phi}^\phi = \frac{1}{r}, \quad \Gamma_{\phi\theta}^\theta = \frac{\cos \theta}{\sin \theta}. \quad (6)$$

Riemann tensor:

$$R_{rtr}^t = \frac{\ddot{a}a}{1-kr^2}, \quad R_{\theta t\theta}^t = r^2 a \ddot{a}, \quad R_{\phi t\phi}^t = r^2 a \ddot{a} \sin^2 \theta \quad (7)$$

$$R_{ttr}^r = \frac{\ddot{a}}{a}, \quad R_{\theta r\theta}^r = r^2(k + \dot{a}^2), \quad R_{\phi r\phi}^r = r^2 \sin^2 \theta(k + \dot{a}^2) \quad (8)$$

$$R_{t\theta\theta}^\theta = \frac{\ddot{a}}{a}, \quad R_{rr\theta}^\theta = -\frac{k + \dot{a}^2}{1-kr^2}, \quad R_{\phi\theta\phi}^\theta = r^2 \sin^2 \theta(k + \dot{a}^2) \quad (9)$$

$$R_{t\phi\phi}^\phi = \frac{\ddot{a}}{a}, \quad R_{rr\phi}^\phi = -\frac{k + \dot{a}^2}{1-kr^2}, \quad R_{\theta\theta\phi}^\phi = -r^2(k + \dot{a}^2). \quad (10)$$

Ricci tensor:

$$R_{tt} = \frac{3\ddot{a}}{a}, \quad R_{rr} = \frac{2k + 2\dot{a}^2 + a\ddot{a}}{1-kr^2} \quad (11)$$

$$R_{\theta\theta} = r^2(2k + 2\dot{a}^2 + a\ddot{a}), \quad R_{\phi\phi} = r^2(2k + 2\dot{a}^2 + a\ddot{a}) \sin^2 \theta. \quad (12)$$

Ricci scalar:

$$R = 3\frac{\ddot{a}}{a} + \left(\frac{2k + 2\dot{a}^2 + a\ddot{a}}{a^2} \right) \cdot 3 = 6 \left(\frac{k}{a} + \frac{\dot{a}^2}{a^2} + \frac{\ddot{a}}{a} \right). \quad (13)$$

(ii) If $k = 0$, defining $dt = a d\eta$ we have

$$ds^2 = -dt^2 + a^2(dr^2 + r^2 d\Omega^2) \quad (14)$$

$$= -a^2 d\eta^2 + a^2(dr^2 + r^2 d\Omega^2) \quad (15)$$

$$= a^2(-d\eta^2 + dr^2 + r^2 d\Omega^2). \quad (16)$$

This metric is proportional to the flat metric, with a proportionality factor depending on the coordinate t . We say that g is conformally flat.

(iii) Let x^μ be the position of some galaxy in the universe. The galaxy is not moving. Let us derive the components of the $T_{\mu\nu}$ in spherical coordinates:

$$T_{00} = (\epsilon + p)(-1)^2 + p(-1) = \epsilon \quad (17)$$

$$T_{rr} = 0 + pg_{rr} = p \frac{a^2}{1 - kr^2} \quad (18)$$

$$T_{\theta\theta} = pr^2 a^2 \quad (19)$$

$$T_{\phi\phi} = pr^2 a^2 \sin^2 \theta \quad (20)$$

Taking the trace of the Einstein equations and solving for R , we obtain

$$R_{\mu\nu} = 8\pi G_N (T_{\mu\nu} - \frac{1}{2} g_{\mu\nu} T^\kappa{}_\kappa) + g_{\mu\nu} \Lambda \quad (21)$$

with

$$T^\mu{}_\mu = g^{\mu\nu} T_{\mu\nu} = -\epsilon + 3p \quad (22)$$

and for the component 00

$$-\frac{\ddot{a}}{a} = 4\pi G_N (p + \frac{1}{3}\epsilon) - \frac{1}{3} \quad (23)$$

and for the component rr

$$\frac{\ddot{a}}{a} + 2\frac{\dot{a}^2}{a^2} + 2\frac{k}{a} = -4\pi G_N (p - \epsilon) + 1. \quad (24)$$

Using the first equation into the second we derive the Friedmann equations:

$$\frac{\dot{a}^2}{a^2} + \frac{k}{a^2} = \frac{8\pi G_N}{3} \epsilon + \frac{1}{3} \Lambda \quad (25)$$

$$\frac{\ddot{a}}{a} = -\frac{4\pi G_N}{3} (\epsilon + 3p) + \frac{1}{3} \Lambda \quad (26)$$