

Black Holes and their Thermodynamics<sup>1</sup>  
**Problem Sheet 4**

Prof. Dieter Lüst<sup>2</sup> and Marvin Lüben<sup>3</sup>

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Tutorials take place on Mondays, 2-4pm (c.t.) in room A 449, Theresienstr. 37. Please hand in your solutions at the next tutorial on November 11, 2019 or send them in a single pdf in an email before the tutorial.

**Exercise 1** *A general, static and spherically symmetric metric in coordinates  $(t, r, \theta, \phi)$  can be written as*

$$ds^2 = -A(r)dt^2 + \frac{1}{B(r)}dr^2 + r^2d\Omega^2, \quad (1)$$

where  $d\Omega^2 = d\theta^2 + \sin^2\theta d\phi^2$  is the metric on a unit 2-sphere. Assume that the  $A(r)$  and  $B(r)$  are analytic functions of  $r$  such that both have a simple zero at  $r = r_+ > 0$  and are positive for  $r > r_+$ .

1. Show that radial null geodesics are given by  $t \pm r^* = \text{constant}$ , where

$$r^* = \int_{r_0}^r \frac{dr'}{\sqrt{A(r')B(r')}} , \quad (2)$$

with  $r_0 > r_+$  an arbitrary constant. Show that  $r^* \rightarrow -\infty$  as  $r \rightarrow r_+$ .

2. Ingoing Eddington-Finkelstein coordinates are defined by

$$v = t + r^* . \quad (3)$$

Obtain the metric in these coordinates. Explain why this metric can be analytically continued through  $r = r_+$ .

3. For  $A(r) = B(r) = 1 - 2M/r$  the metric in eq. (1) is the Schwarzschild metric. Which describes a black hole of mass  $M$ . The Kretschmann invariant is defined by

$$K = R^{abcd}R_{abcd} \quad (4)$$

and given by  $K = 48M^2/r^6$  for the spacetime defined by the Schwarzschild metric. Comment on the definitions of spacetime and coordinate singularities.

<sup>1</sup> [www.physik.uni-muenchen.de/lehre/vorlesungen/wise\\_19\\_20/bh\\_info\\_wise\\_2019\\_20](http://www.physik.uni-muenchen.de/lehre/vorlesungen/wise_19_20/bh_info_wise_2019_20)

<sup>2</sup> [dieter.luest@lmu.de](mailto:dieter.luest@lmu.de)

<sup>3</sup> [mlueben@mpp.mpg.de](mailto:mlueben@mpp.mpg.de)

**Exercise 2** *In this exercise we will show, that the region  $r < 2M$  of the Schwarzschild solution indeed describes a black hole. First, let us introduce the following definitions*

- *A vector is called causal if it is timelike or null. A curve is called causal if its tangent vector is everywhere causal.*
- *Given a time-orientation of the spacetime: a causal vector field  $T^a$ . A vector field  $X^a$  is called future-directed if it lies in the same lightcone as  $T^a$ .*

*For the Schwarzschild spacetime let  $\partial_t$  denote the time-orientation for  $r > 2M$ . Show that every future-directed causal curve starting inside the event horizon  $r < 2M$  will not cross the event horizon. The following steps may provide a guideline:*

1. *Argue that  $-\partial_r$  is future-oriented causal.*
2. *Let  $x^\mu(\lambda)$  be a future-directed causal curve. Contract the tangent vector  $dx^\mu/d\lambda$  with  $-\partial_r$  to determine the sign of  $\dot{v}$ .*
3. *After these preliminaries, inspect  $\dot{x}^2$  for  $r < 2M$ .*
4. *In addition, what can you say about future-directed causal curves that start on the event horizon?*

**Hint:** *Work in ingoing Eddington-Finkelstein coordinates from the previous exercise.*

**Exercise 3** *Take the Schwarzschild metric, i.e. set  $A(r) = B(r) = 1 - 2M/r$  in eq. (1).*

1. *Write the Lagrangian of the geodesic equation for the Schwarzschild metric. Derive the components of the geodesic equation. Determine the Christoffel symbols from the geodesic equation.*
2. *Note that the Lagrangian from which geodesics are obtained is independent of  $t$  and  $\phi$ . What consequences does this have?*