

Black Holes and their Thermodynamics¹

Problem Sheet 10

Prof. Dieter Lüst² and Marvin Lüben³

December 16, 2019

Tutorials take place on Mondays, 2-4pm (c.t.) in room A 449, Theresienstr. 37. There will be no tutorial on December 23. Please hand in your solutions at the next tutorial on January 10, 2020 or send them in a single pdf in an email before the tutorial. Note the unusual date and time.

Exercise 1 *The Kerr-Newman black hole with electrical charge Q is given by*

$$ds^2 = -\frac{\rho^2 \Delta}{\Sigma} dt^2 + \frac{\Sigma}{\rho^2} \sin^2 \theta (d\phi - \omega dt)^2 + \frac{\rho^2}{\Delta} dr^2 + \rho^2 d\theta^2 \quad (1)$$

$$A_t = \frac{Qr}{\rho^2}, \quad A_\phi = -\frac{Qar \sin^2 \theta}{\rho^2} \quad (2)$$

where

$$\Delta = r^2 - 2Mr + a^2 + Q^2, \quad \rho^2 = r^2 + a^2 \cos^2 \theta, \quad (3)$$

$$\Sigma = (r^2 + a^2)^2 - a^2 \Delta \sin^2 \theta, \quad \omega = \frac{(2Mr - Q^2)a}{\Sigma} \quad (4)$$

1. *The co-rotating electric horizon potential is defined by*

$$\Phi_H = \xi^\mu A_\mu \Big|_{r=r_+}, \quad \text{with } \xi = \partial_t + \Omega_H \partial_\phi, \quad (5)$$

where $\Omega_H = \frac{a}{r_+^2 + a^2}$. Here we assume that $M^2 > a^2 + Q^2$ and r_+ is the radial position of the outer horizon of the Kerr-Newman black hole. Determine r_+ . Show that

$$\Phi_H = \frac{Qr_+}{r_+^2 + a^2} \quad (6)$$

2. *Convince yourself that the area of the (outer) event horizon is still given by $A = 4\pi(r_+^2 + a^2)$ and verify Smarr's formula*

$$M = \frac{\kappa A}{4\pi} + 2\Omega_H J + \Phi_H Q, \quad (7)$$

where $\kappa = \frac{r_+ - r_-}{2(r_+^2 + a^2)}$ is the surface gravity of the outer event horizon.

¹ www.physik.uni-muenchen.de/lehre/vorlesungen/wise_19_20/bh_info_wise_2019_20

² dieter.luest@lmu.de

³ mlueben@mpp.mpg.de

3. Use this formula for M to deduce the first law of black hole mechanics for charged rotating black holes:

$$dM = \kappa dA + \Omega_H dJ + \Phi_H dQ. \quad (8)$$

Exercise 2 In this exercise we will study implications of Hawking's area theorem.

1. The area law implies a limit to the efficiency of mass/energy conversion in black hole collisions. Consider two black holes that initially are well separated so that we can approximate them as Schwarzschild black holes (i.e. uncharged and non-rotating) of masses M_1 and M_2 , respectively. Assume further that the final black hole is again a Schwarzschild black hole of mass M_3 . The energy difference $E = M_1 + M_2 - M_3$ is radiated away by gravitational radiation. Use the black hole area law to show

$$\eta \equiv \frac{E}{M_1 + M_2} \leq 1 - \frac{1}{\sqrt{2}} \quad (9)$$

for the efficiency of the process.

2. Assuming the area law show that a Schwarzschild black hole of mass M_3 can never bifurcate into two Schwarzschild black holes of masses M_1 and M_2 , respectively.

Exercise 3 Identifying the surface gravity κ of the (outer) event horizon with the temperature of a black hole via

$$T = \frac{\hbar}{2\pi} \kappa \quad (10)$$

leads to a negative heat capacity for a Schwarzschild black hole (Check this!), i.e. the temperature decreases with increasing mass. Plot T/\hbar as a function of M/Q for a Reissner-Nordström black hole for a fixed charge Q , and discuss the heat capacity in that case.

(**Hint:** The heat capacity of a charged black hole of mass M and fixed charge Q is $C = T \frac{\partial S}{\partial T} |_Q$ where $S = A/4\hbar$ is the Bekenstein-Hawking entropy of the black hole.)

Exercise 4 Consider the Klein-Gordon equation of a massive, complex scalar field in Minkowski spacetime,

$$(\nabla^2 - m^2)\phi = 0. \quad (11)$$

1. Show that the inner product on the space of solutions to the Klein-Gordon equation, i.e.

$$\langle \phi_1, \phi_2 \rangle = -i \int_{\Sigma_t} d^{D-1}x (\phi_1 \partial_t \phi_2^* - \phi_2^* \partial_t \phi_1) \quad (12)$$

is independent of the constant-time hypersurface Σ_t over which the integral is taken.

2. Use the inner product (11) to show

$$\langle f_{\vec{k}_1}, f_{\vec{k}_2} \rangle = -\langle f_{\vec{k}_1}^*, f_{\vec{k}_2}^* \rangle = \delta^{D-1}(\vec{k}_1 - \vec{k}_2), \quad \langle f_{\vec{k}_1}, f_{\vec{k}_2}^* \rangle = 0, \quad (13)$$

with

$$f_{\vec{k}} = \frac{e^{ik_\mu x^\mu}}{\sqrt{(2\pi)^{D-1} 2\omega}}, \quad \omega^2 = \vec{k}^2 + m^2, \quad \omega > 0 \quad (14)$$

3. Argue that the vacuum state does not change under a Lorentz transformation.