Homepage: ...TVI\_TMP-TA4\_-Condensed-Matter-Field-Theory/

Nov. 20, 2019

Prof. Dr. J. von Delft, Dr. O. Yevtushenko, F. Kugler

## Problem Set 6: Spin representations

**Exercise 1.** Fermionic representation of S = 1/2 magnets

We consider a spin system on a lattice, with a spin-1/2 representation on each lattice site. Hence, the spin operators on sites r, r' fulfill

$$[S_a(r), S_b(r')] = i\epsilon_{abc}S_c(r)\delta_{r,r'},\tag{1}$$

$$\sum_{a \in \{x, y, z\}} S_a(r) S_a(r) = 3/4. \tag{2}$$

(a) Consider a Schwinger-fermion representation of the spin operators, defined by

$$S_a(r) = \frac{1}{2} f_{\alpha}^{\dagger}(r) \sigma_{\alpha\beta}^a f_{\beta}(r), \quad \alpha, \beta \in \{1, 2\},$$
(3)

where  $f_{\alpha}$  are fermionic operators,  $\{f_{\alpha}(r), f_{\beta}^{\dagger}(r')\} = \delta_{\alpha\beta}\delta_{r,r'}$ , and  $\sigma^a$  are Pauli matrices. Check that  $[S_x(r), S_y(r)] = iS_z(r)$  holds in this representation. Is there a gauge redundancy in this mapping?

- (b) Check that (2) does not hold for all states in the fermionic on-site Hilbert space. Formulate a constraint on the Hilbert space such that (2) is fulfilled.
- (c) Consider now a Majorana-fermion representation, defined by

$$S_a(r) = -\frac{i}{2} \epsilon_{abc} \gamma_b(r) \gamma_c(r). \tag{4}$$

Here,  $\gamma_a, a \in \{x, y, z\}$ , are "real" (i.e., self-adjoint) fermions in the sense that they satisfy

$$\{\gamma_a(r), \gamma_b(r')\} = \delta_{ab}\delta_{r,r'}, \qquad \gamma^{\dagger}(r) = \gamma(r).$$
 (5)

Check that  $[S_x(r), S_y(r)] = iS_z(r)$  also holds in this representation. Does (2) hold as well? Is there a gauge redundancy in this mapping?

- (d) For the construction of a Majorana path integral, we want to characterize the local Hilbert space in terms of Majorana operators  $\gamma_i$ . In order to do this, it is convenient to introduce a further fourth "Majorana" operator  $\gamma_4$  that completes the operator algebra (5). Express now  $\gamma_i$  as real or imaginary parts of two complex fermions c, d.
- (e) Based on (d), argue that the Majorana Lagrangian is of the form

$$\mathcal{L} = \frac{1}{2} \sum_{r} \gamma_{a,r} \partial_{\tau} \gamma_{a,r} + H[\gamma]. \tag{6}$$

(f) To construct  $H[\gamma]$ , we need the Hamiltonian to be normal ordered. We want to show that a quadratic Hamiltonian in the Majorana basis is quite often trivial. Let us assume the following rather general Hamiltonian

$$H_r = J \sum_{i,j \in \{x,y\}} \gamma_i(r) \gamma_j(r). \tag{7}$$

and bring it into a normal ordered form by using the Dirac fermions c, d at first. Can you think of a possibility to construct a quadratic Hamiltonian in Majorana representation that is nontrivial?

(g) As a minimal example, we consider a single spin in a magnetic field h:

$$H = -hS_z$$
.

Write down the action for the Majorana fermions  $\gamma_x$  and  $\gamma_y$ , and transform it to the Matsubara-frequency representation. Compute the magnetization  $M=\langle S_z\rangle$  as a Majorana-fermion correlator. Finally, compute the susceptibility  $\chi=\mathrm{d}M/\mathrm{d}h|_{h=0}$  and demonstrate the Curie law for high temperatures.

You can check your results by solving this problem with standard (quantum-mechanical) techniques.