

# STRING THEORY I

LUDWIG-MAXIMILIANS-UNIVERSITÄT MÜNCHEN  
DR. MICHAEL HAACK

---

## Exercise sheet 9

DUE ON DECEMBER 20, 2019

### Exercise 1: Polarization tensor

Consider the closed string state

$$|\phi\rangle = \xi_{\mu\nu} \alpha_{-1}^{\mu} \tilde{\alpha}_{-1}^{\nu} |0; k\rangle, \quad (1.1)$$

where  $\xi_{\mu\nu}$  denotes the polarization tensor.

(a) What do the Virasoro constraints on physical states (for  $A = -1$ ),

$$(L_0 - 1)|\phi\rangle = (\tilde{L}_0 - 1)|\phi\rangle = L_m|\phi\rangle = \tilde{L}_m|\phi\rangle = 0, \quad m > 0, \quad (1.2)$$

imply for  $k$  and  $\xi_{\mu\nu}$ ? Show first that if a state is annihilated by  $L_1$  and  $L_2$  it is annihilated by all  $L_n$  with  $n \geq 1$ .

(b) Show the presence of the gauge equivalence

$$\xi_{\mu\nu} \cong \xi_{\mu\nu} + a_{\mu}k_{\nu} + k_{\mu}b_{\nu}, \quad a \cdot k = b \cdot k = 0 \quad (1.3)$$

by finding the relevant null states, i.e. physical states which are spurious (i.e. orthogonal to all physical states).

### Exercise 2: A ghost at the second excited level

Consider the open string state (with NN boundary conditions in all directions)

$$|\phi_k\rangle = \left( c_1 \alpha_{-1} \cdot \alpha_{-1} + c_2 \alpha_0 \cdot \alpha_{-2} + c_3 (\alpha_0 \cdot \alpha_{-1})^2 \right) |0; k\rangle. \quad (2.1)$$

Determine the relations between  $c_1, c_2$  and  $c_3$  which follow when demanding  $|\phi_k\rangle$  to be physical (use  $A = -1$ , i.e.  $(L_0 - 1)|\phi_k\rangle = 0$ ). Then compute

$$\langle \phi_k | \phi_{k'} \rangle = \frac{2c_1^2}{25} (D-1)(26-D) \delta^D(k-k'). \quad (2.2)$$

### Exercise 3: Field equations of the Kalb-Ramond field $B_{\mu\nu}$

Here we examine the field theory of a massless antisymmetric tensor gauge field  $B_{\mu\nu} = -B_{\nu\mu}$ . This gauge field is a tensor analog of the Maxwell gauge field  $A_{\mu}$ . In Maxwell theory one defines the field strength  $F_{\mu\nu} = \partial_{\mu}A_{\nu} - \partial_{\nu}A_{\mu}$ . For  $B_{\mu\nu}$  one defines a field strength  $H_{\mu\nu\rho}$  according to

$$H_{\mu\nu\rho} \equiv \partial_{\mu}B_{\nu\rho} + \partial_{\nu}B_{\rho\mu} + \partial_{\rho}B_{\mu\nu}. \quad (3.1)$$

(a) Show that  $H_{\mu\nu\rho}$  is totally antisymmetric.

(b) Prove that  $H_{\mu\nu\rho}$  is invariant under the gauge transformations

$$\delta B_{\mu\nu} = \partial_\mu \epsilon_\nu - \partial_\nu \epsilon_\mu . \quad (3.2)$$

(c) The above gauge transformations have a special feature: the gauge parameters themselves have a gauge invariance. Show that  $\epsilon'_\mu$  given by

$$\epsilon'_\mu = \epsilon_\mu + \partial_\mu \lambda \quad (3.3)$$

generates the same gauge transformation as  $\epsilon_\mu$ .

(d) Consider the spacetime action

$$S = \int d^D x \left( -\frac{1}{6} H_{\mu\nu\rho} H^{\mu\nu\rho} \right). \quad (3.4)$$

Find the field equation for  $B_{\mu\nu}$ .

---

For questions:

michael.haack@lmu.de