

# STRING THEORY I

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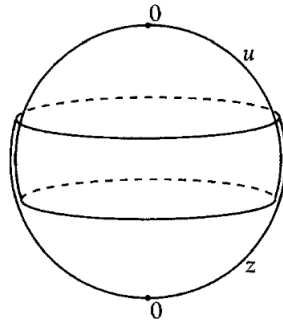
## Exercise sheet 5

DUE ON NOVEMBER 22, 2019

### Exercise 1: Metric on the sphere

The sphere  $S_2$  can be covered by two coordinate patches. One can introduce complex coordinates in each patch,  $z = z_1 + iz_2$  and  $u = u_1 + iu_2$  and take as the open neighborhoods of the north and south pole the open disks  $|z| < \rho$  and  $|u| < \rho$ , with  $\rho > 1$ . The transition function in the overlap region can be taken to be

$$u = 1/z . \quad (1.1)$$



A globally defined metric for a sphere of radius  $r$  is given by

$$ds^2 = \frac{4r^2 dz d\bar{z}}{(1 + |z|^2)^2} = \frac{4r^2 du d\bar{u}}{(1 + |u|^2)^2} , \quad (1.2)$$

(cf. exercise 2, sheet 2, for  $r = 1$ ). This metric has constant curvature  $R = 2/r^2$  (leading to the Euler number  $\chi = \frac{1}{4\pi} \int_{S_2} \sqrt{g} R = 2$ ).

(a) Show that (1.2) is indeed a well defined global metric, i.e. verify that the two expressions in the  $u$  and the  $z$  coordinates agree in the overlap region.

(b) Obviously, one can perform Weyl-transformations in the  $z$ - and the  $u$ -patches separately that bring (1.2) to the flat gauge in each of the patches. Show that this does not correspond to a globally well defined Weyl-transformation though, i.e. the two Weyl-transformations do not agree in the overlap region.

**Remark:** Put differently, choosing the flat gauge in both patches simultaneously does not lead to a well defined global metric, as the two expressions do not agree in the overlap region.

## Exercise 2: Angular momentum tensor

Show that for a closed string the angular momentum tensor

$$J^{\mu\nu} = T \int_0^{2\pi} d\sigma (X^\mu \partial_\tau X^\nu - X^\nu \partial_\tau X^\mu) \quad (2.1)$$

is given by

$$J^{\mu\nu} = x^\mu p^\nu - x^\nu p^\mu - i \sum_{n \neq 0} \frac{1}{n} (\alpha_{-n}^\mu \alpha_n^\nu + \tilde{\alpha}_{-n}^\mu \tilde{\alpha}_n^\nu) . \quad (2.2)$$

## Exercise 3: DD and ND boundary conditions for open strings

(a) Consider an open string coordinate  $X^a$  with Dirichlet boundary conditions on both sides, i.e.

$$X^a(\tau, 0) = x_1^a \quad , \quad X^a(\tau, \pi) = x_2^a . \quad (3.1)$$

Show that the most general solution to the open string equations of motion with these boundary conditions can be expanded as

$$X^a = x_1^a + \frac{x_2^a - x_1^a}{\pi} \sigma + \sqrt{2\alpha'} \sum_{n \neq 0} \frac{1}{n} \alpha_n^a e^{-in\tau} \sin(n\sigma) \quad (3.2)$$

with  $\alpha_{-n}^a = (\alpha_n^a)^*$ .

(b) Now consider an open string coordinate  $X^r$  with Neumann boundary conditions on one side and Dirichlet boundary conditions on the other side, i.e.

$$\frac{\partial X^r}{\partial \sigma}(\tau, 0) = 0 \quad , \quad X^r(\tau, \pi) = x_2^r . \quad (3.3)$$

Show that the most general solution to the open string equations of motion with these boundary conditions can be expanded as

$$X^r = x_2^r + i\sqrt{2\alpha'} \sum_{n \in \mathbb{Z}_{\text{odd}}} \frac{2}{n} \alpha_{\frac{n}{2}}^r e^{-i\frac{n}{2}\tau} \cos\left(\frac{n\sigma}{2}\right) \quad (3.4)$$

with  $\alpha_{-\frac{n}{2}}^r = (\alpha_{\frac{n}{2}}^r)^*$ .

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For questions:

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