


Neutrino Physics Course

Lecture III

2/5/2023

LMU

Spring 2023



Majorana mass for neutrino?

$$\psi_L = \begin{pmatrix} u_L \\ 0 \end{pmatrix}, \quad \psi_R = \begin{pmatrix} 0 \\ u_R \end{pmatrix}$$

$$u_{L,R} \rightarrow e^{i\bar{\sigma}/2} (\vec{\sigma} \pm i\vec{\varphi})$$

ROT BOOST

$$\begin{aligned} \bullet \quad m_D \bar{\psi} \psi &= m_D (\bar{\psi}_L \psi_R + \bar{\psi}_R \psi_L) \\ &= m_D (u_L^\dagger u_R + u_R^\dagger u_L) \end{aligned}$$

$m_D =$ Dirac mass term

• $m_H (u_L^T i \sigma_2 u_L + h.c.) \left(\frac{1}{2} \right)$

Hejlskov 1937

↑
see below

$u_L \rightarrow e^{i \vec{\sigma}/2 \cdot \vec{\chi}} u_L, \quad \vec{\chi} = \vec{\theta} + i \vec{\varphi}$

$\Rightarrow u_L^T i \sigma_2 u_L \rightarrow u_L^T e^{i \vec{\sigma}^T/2 \cdot \vec{\chi}} (i \sigma_2) e^{i \vec{\sigma}/2 \cdot \vec{\chi}} u_L$

$= u_L^T i \sigma_2 \underbrace{e^{-i \vec{\sigma}^T/2 \cdot \vec{\chi}} e^{i \vec{\sigma}/2 \cdot \vec{\chi}}}_1 u_L$

$= u_L^T i \sigma_2 u_L$
Q.E.D.

• $u_L^T i \sigma_2 u_L = \boxed{\psi_L^T C \psi_L} =$

$= (u_L^T \ 0) i \sigma_2 \begin{pmatrix} u_L \\ 0 \end{pmatrix}$

$$= \begin{pmatrix} u_L^T & 0 \end{pmatrix} \begin{pmatrix} i\sigma_2 & 0 \\ 0 & -i\sigma_2 \end{pmatrix} \begin{pmatrix} u_L \\ 0 \end{pmatrix}$$



Q.E.D.

$$\psi_L^T C \psi_L = \text{inv.} \Leftrightarrow \psi^c = C \bar{\psi}^T$$

Proof:

is a spinor

$$\psi = \psi_L \Rightarrow \psi^c = i\sigma_2 \gamma_0 \gamma_0 \psi_L^*$$

$$= \begin{pmatrix} 0 & i\sigma_2 \\ -i\sigma_2 & 0 \end{pmatrix} \begin{pmatrix} u_L^* \\ 0 \end{pmatrix}$$

$$(\psi^c)_R = \begin{pmatrix} 0 \\ -i\sigma_2 u_L^* \end{pmatrix} = C \bar{\psi}_L^T$$



$$\overline{(\psi^c)_R} \psi_L = iuv. \quad (*)$$

\uparrow \uparrow
 spinor spinor

$\underbrace{\hspace{10em}}$
 Lorentz

$$\bar{\psi} \psi \longrightarrow \bar{\psi}_1 \psi_2 = iuv.$$

\parallel
 inv.

$$\begin{aligned}
 (*) \quad \overline{(\psi^c)_R} \psi_L &= \overline{C \bar{\psi}_L^T} \psi_L = \\
 &= (C \gamma_0 \psi_L^*)^+ \gamma^0 \psi_L \\
 &= \psi_L^T \gamma_0 C^+ \gamma^0 \psi_L = \frac{1}{2} C^+ = -C \zeta
 \end{aligned}$$

$$= -\psi_L^T \gamma_0 C \partial_0 \psi_L = \psi_L^T C \psi_L$$

Q. E. D.

$$m_M \underbrace{\psi_L^T C \psi_L}$$

breaks e.g. global symmetry

$$\psi_L \rightarrow e^{i\alpha} \psi_L \quad \therefore$$

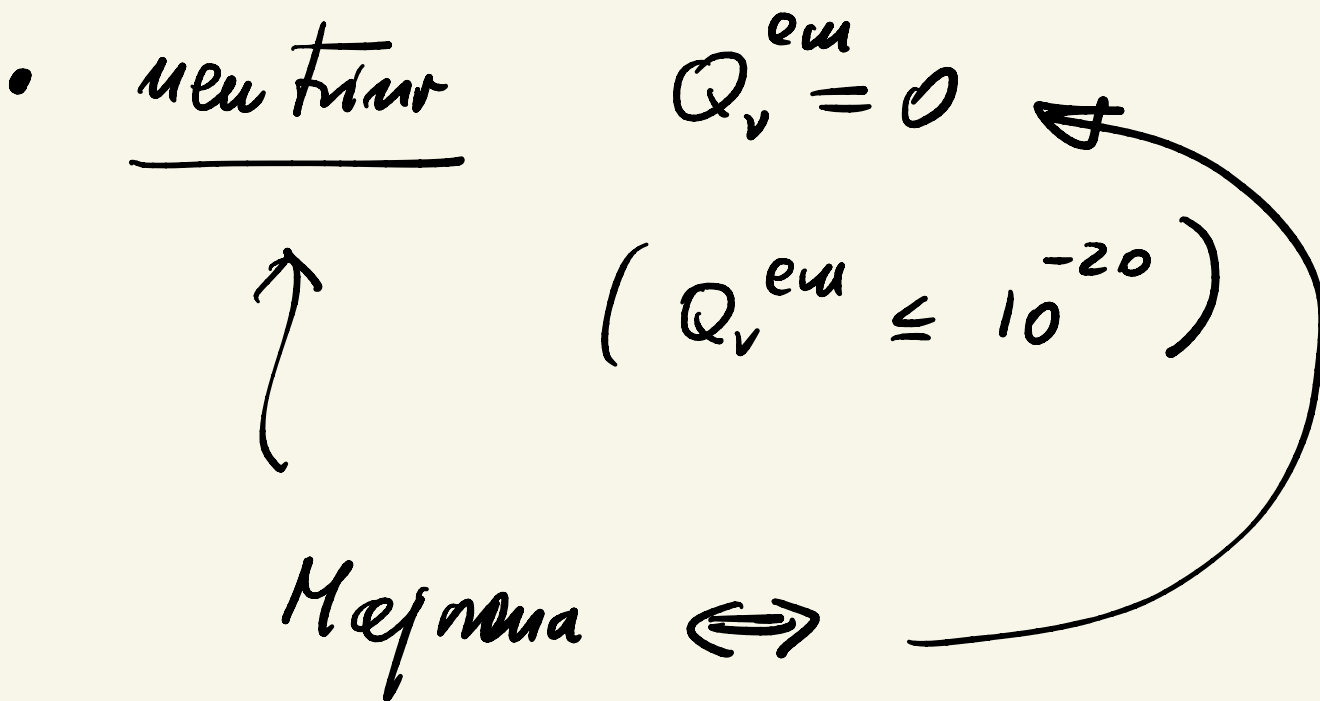
$$\psi_L^T C \psi_L \rightarrow e^{2i\alpha} \psi_L^T C \psi_L$$

• electron $\psi = e$

$$m_D \bar{e}_L e_R + \cancel{m_M^L e_L^T C e_L} + \cancel{m_M^R e_R^T C e_R}$$

$$m_M^L / m_D \leq 10^{-20}$$

$$m_M^R / m_D \leq 10^{-20}$$



Lepton number (L)

$$n \rightarrow p + e + \bar{\nu}$$

$$\Rightarrow L(\nu) = L(e) = 1$$

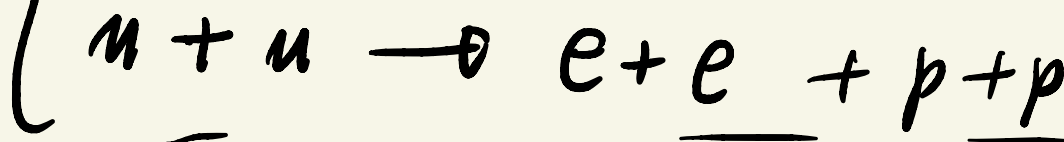
neutrino is a lepton

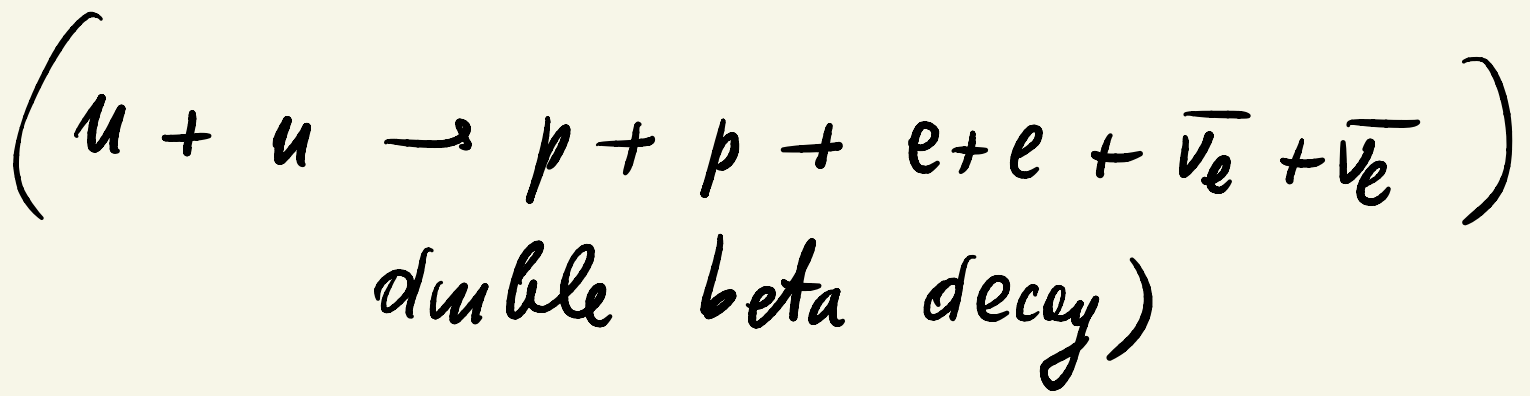
$$\nu \rightarrow e^{i\beta L} \nu = e^{i\beta} \nu$$

$$e \rightarrow e^{i\beta L} e = e^{i\beta} e$$

$$\therefore \psi_L^\dagger c \psi_L \Rightarrow \Delta L = 2$$

Racah 1937 } neutrino-less double
Furry 1938 } beta decay





Dirac $\psi_D = \begin{pmatrix} u_L \\ u_R \end{pmatrix}$ independent

$$\Leftrightarrow \psi = 4 \text{ components}$$

↑
Complex

Majorana

$$\begin{aligned} \psi_M &= \psi_L + c \bar{\psi}_L^T \\ &= \begin{pmatrix} u_L \\ 0 \end{pmatrix} + \begin{pmatrix} 0 \\ -i\sigma_2 u_L^* \end{pmatrix} \end{aligned}$$

$$= \begin{pmatrix} u_L \\ -i \sigma_2 u_L^* \end{pmatrix}$$

$\Leftrightarrow \psi_M = 2 \text{ complex comp.}$

$$\bullet \bar{\psi}_M \psi_M = (\bar{\psi}_L + c \bar{\psi}_L^T) (\psi_L + c \bar{\psi}_L^T)$$

$$= \cancel{\bar{\psi}_L \psi_L} + \bar{\psi}_L c \bar{\psi}_L^T +$$

$$+ c \bar{\psi}_L^T \psi_L + \cancel{c \bar{\psi}_L^T c \bar{\psi}_L^T}$$

\downarrow

$$= (c \gamma_0 \psi_L^*)^\dagger \gamma^0 \psi_L + \bar{\psi}_L c \bar{\psi}_L^T$$

$$= \psi_L^T \gamma_0 c + \gamma^0 \psi_L + \psi_L^\dagger \gamma^0 c (\psi_L^\dagger \gamma^0)^T$$

$$= \psi_L^T (-c^+) \psi_L + \psi_L^T \gamma^0 c \gamma_0 \psi_L^*$$

$$= \underbrace{\psi_L^T C \psi_L}_{\text{Majorana}} + \underbrace{\psi_L^T c^+ \psi_L^*}_{\text{h.c.}}$$

$$\bullet i \bar{\psi}_M \gamma^\mu \partial_\mu \psi_M = i \overline{(\psi_L + c \bar{\psi}_L^T)} \gamma^\mu \partial_\mu (\psi_L + c \bar{\psi}_L^T)$$

$$= i \bar{\psi}_L \gamma^\mu \partial_\mu \psi_L + \overline{c \bar{\psi}_L^T} \gamma^\mu \partial_\mu c \bar{\psi}_L^T$$

$$- \partial_\mu \left(\frac{\parallel}{\text{---}} \right) -$$

$$- \partial_\mu \overline{c \bar{\psi}_L^T} \gamma^\mu c \bar{\psi}_L^T$$

||

$$= i \bar{\Psi}_L \gamma^\mu \partial_\mu \Psi_L + i \bar{\Psi}_L \gamma^\mu \partial_\mu \Psi_L$$

$$= 2i \bar{\Psi}_L \gamma^\mu \partial_\mu \Psi_L$$



$$\mathcal{L}_M = i \bar{\Psi}_M \gamma^\nu \partial_\nu \Psi_M - \underbrace{\bar{\Psi}_M \Psi_M}_{\text{mass term}} m_M \quad (1)$$

$$(E^2 = \vec{p}^2 + m_M^2)$$

$$= 2i \bar{\Psi}_L \gamma^\mu \partial_\mu \Psi_L - m_M (\Psi_L^T C \Psi_L + \text{h.c.})$$

$$= 2 \left[i \bar{\Psi}_L \partial^\mu \partial_\mu \Psi_L - \frac{m_M}{2} (\Psi_L^T C \Psi_L + \text{h.c.}) \right]$$

(2)

$$\bullet \bar{\psi}_M \psi_M \quad \therefore \quad \psi_M \rightarrow e^{i\alpha} \bar{\psi}_M \quad ???$$

$$\psi_M = \psi_L + c \bar{\psi}_L^T = \psi_L + c \gamma_0 \psi_L^*$$

$$\left. \begin{array}{ccc} & \downarrow & \downarrow \\ & e^{i\alpha} \psi_L & e^{-i\alpha} \dots \end{array} \right\}$$

50 % + 50 %

particle

anti-particle

Rise and fall of parity

$$\mathcal{L}_{QED} = i \bar{\psi} \gamma^\mu D_\mu \psi - m \bar{\psi} \psi$$

$$\boxed{F_{\mu\nu} = \partial_\mu A_\nu - \partial_\nu A_\mu} - \frac{1}{4} F_{\mu\nu} F^{\mu\nu} \quad (3)$$

$$D_\mu = \partial_\mu - ie Q^{em} A_\mu$$

\downarrow \uparrow photon

$$Q^{em} = \text{diag} (e_1, e_2, \dots)$$

$$e_e = -1, \quad e_p = +1, \quad e_u = 0$$

$$e_u = 2/3, \quad e_d = -1/3$$

$$|e_p|/|e_d| = 1 + O(10^{-20})$$

$$= 1$$

$$\boxed{e \propto n e_d}$$

charge
quantization

Dirac if \exists magnetic monopole

$$\Rightarrow \boxed{\oint_e \oint_m = 2\pi n}$$

eq. (31) $\psi \rightarrow e^{i\alpha(x)} Q \psi$

is inv.

$$A_\mu \rightarrow A_\mu + \frac{1}{e} \partial_\mu \alpha(x)$$

$$\mathcal{L}_{int} = e \bar{\psi} Q \gamma^\mu \psi A_\mu$$

$$P: A_i \xrightarrow{P} -A_i, \quad A_0 \xrightarrow{P} A_0$$

$$\psi \xrightarrow{P} \gamma^0 \psi$$

Proof: $\gamma^0 \gamma^i \gamma^0 = -\gamma^i$

$$\gamma^0 \gamma^0 \gamma^0 = \gamma^0 \quad \text{Q.E.D.}$$

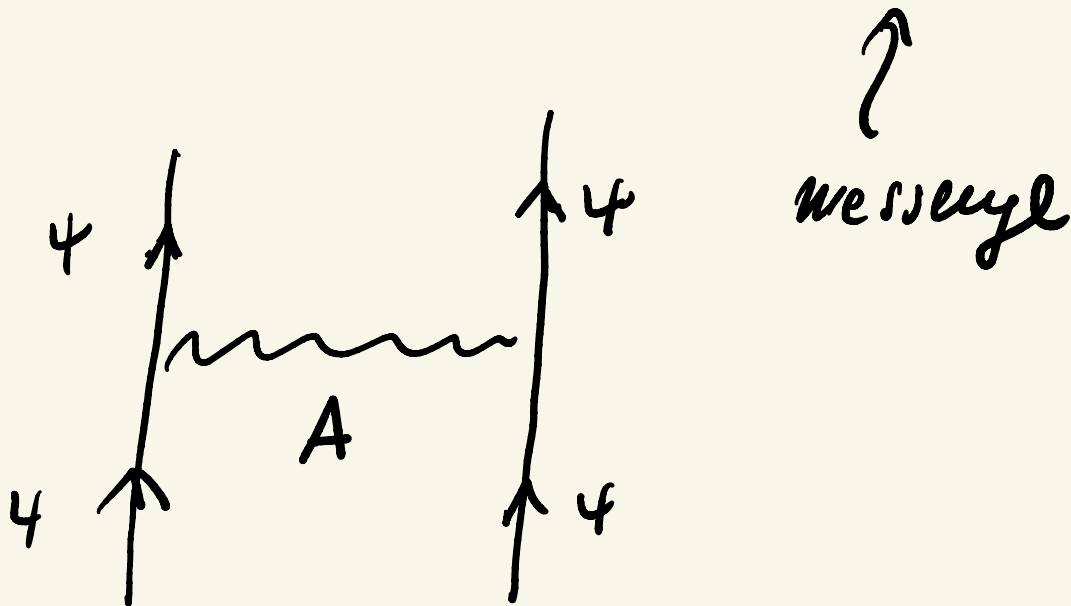
$$\begin{pmatrix} u_L \\ u_R \end{pmatrix} \xrightarrow{P} \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} \begin{pmatrix} u_L \\ u_R \end{pmatrix} \\ = \begin{pmatrix} u_R \\ u_L \end{pmatrix}$$

$$\boxed{u_L \xleftrightarrow{P} u_R}$$

$$\mathcal{H}_{\text{eff}}^{\text{em}} \propto \int_{\mu}^{\text{av}} J_{\mu}^{\text{em}} J_{\text{em}}^{\mu}$$

$$J_{\text{av}}^{\mu} = \bar{\psi} \gamma^{\mu} Q_{\text{em}} \psi$$

$$\Rightarrow \mathcal{L}_{\text{fund}} = e J_{\mu}^{\text{em}} A_{\mu} \text{ (photon)}$$



Fermi 1934

$$\mathcal{H}_{\text{eff}}^{\text{weak}} \propto J_{\nu}^{\mu} J_{\mu}^{-\nu}$$

$$\mathcal{L}_{\text{fund}} = \frac{g}{\sqrt{2}} J_{\nu}^{\mu} W_{\mu} + \text{h.c.}$$



What is the form of J_W^μ ?

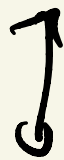
$$J_W^\mu = \underbrace{\bar{u} \gamma^\mu d + \bar{e} \gamma^\mu \nu}_{}$$

form?

$$\bar{u}_L \gamma^\mu d_L ?$$

$$\bar{u}_R \gamma^\mu d_R ?$$

$$\bar{u}_L \gamma^\mu d_L + \bar{u}_R \gamma^\mu d_R ?$$



$$QED: \quad \bar{e} \gamma^\mu e = \bar{e}_L \gamma^\mu e_L + \bar{e}_R \gamma^\mu e_R$$

$\underbrace{\hspace{10em}}_P$

1956 Lee + Yang

Summer

1956 Wu et al

Christmas



~~P~~ maximal

Parity!

$u_L \longleftrightarrow u_R$

$$z_L \equiv L \psi \equiv \frac{1 + \gamma_5}{2} \psi = \begin{pmatrix} u_L \\ 0 \end{pmatrix}$$

$$m_\psi = 0 \Rightarrow i \gamma^\mu \partial_\mu \psi = 0$$

$$\psi(x) = \tilde{\psi}(p) e^{i p x}$$

$$\Rightarrow \boxed{p_\mu \gamma^\mu \tilde{\psi}(p) = 0}$$



$$\begin{pmatrix} 0 & E - \vec{p} \cdot \vec{\sigma} \\ E + \vec{p} \cdot \vec{\sigma} & 0 \end{pmatrix} \begin{pmatrix} u_L \\ u_R \end{pmatrix} = 0$$

$$\Rightarrow \frac{1}{2} E u_{L,R} = \mp \vec{p} \cdot \vec{\sigma} \frac{1}{2} u_{L,R}$$

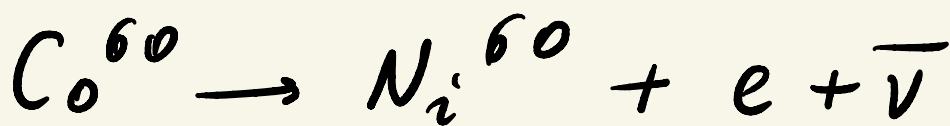
$$\Rightarrow \frac{1}{2} u_{L,R} = \mp \frac{\vec{p} \cdot \vec{\sigma}}{E} u_{L,R}$$

$$|\vec{p}| = E \quad (m=0)$$

$$\vec{S} = \vec{p} \frac{1}{2}$$

$$\Rightarrow \hbar u_{L,R} = \mp \frac{1}{2} u_{L,R}$$

$$\hbar \equiv \vec{S} \cdot \hat{p}$$



$$J=5$$

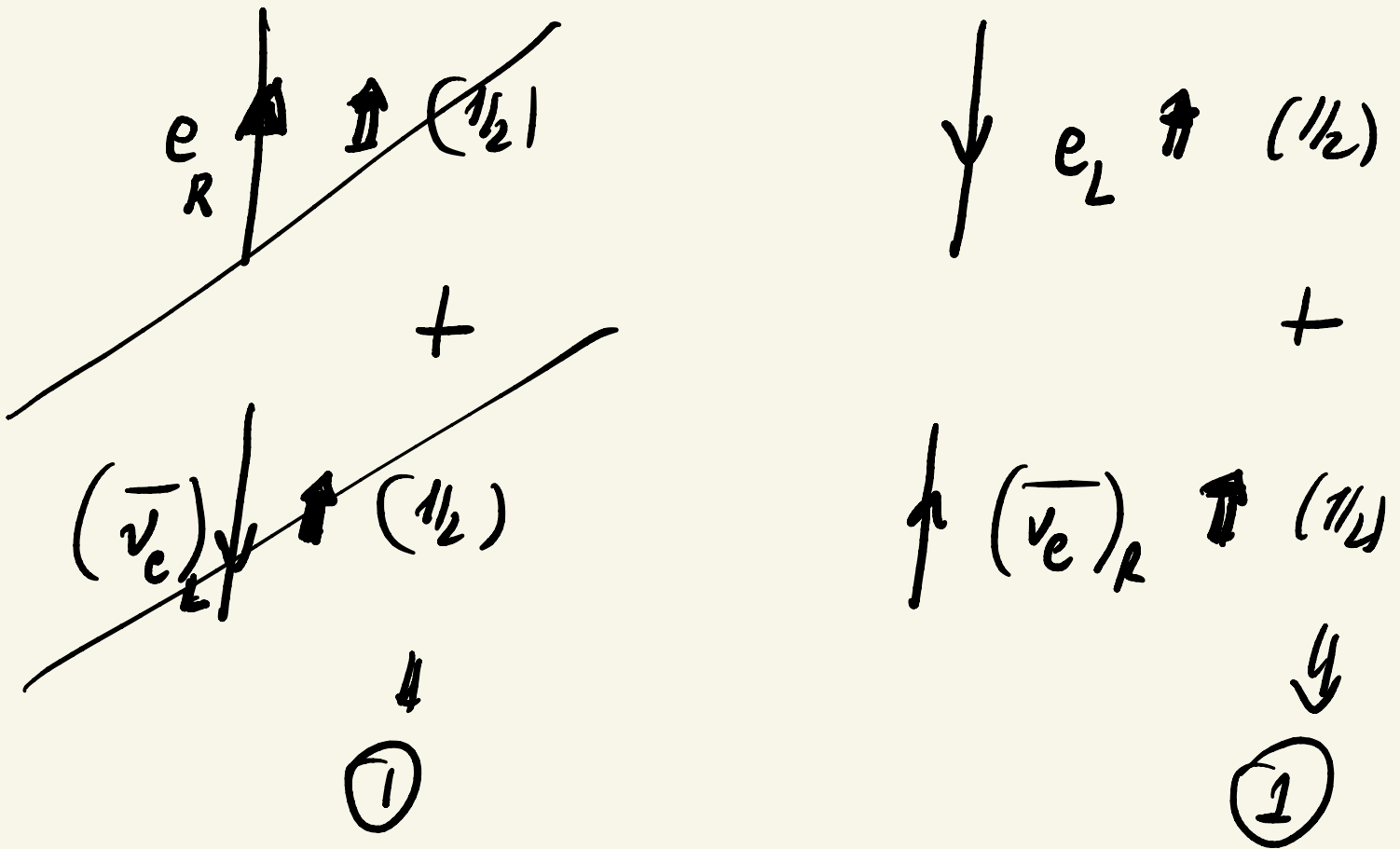
$$J=4$$



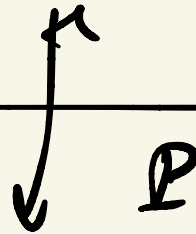
$$M_W = 80 \text{ GeV} = 80 \text{ mp}$$

$$J_z = +1$$

$$m_e = 0 = m_\nu$$



$$J_w^\mu = \bar{\nu}_L \gamma^\mu e_L + \bar{u}_L \gamma^\mu d_L$$



~~$$\bar{\nu}_R \gamma^\mu e_R + \bar{u}_R \gamma^\mu d_R$$~~

