

Problem 1

(a) The Kaon is the lowest mass particle, which contains an s-quark. Since only the weak interactions don't conserve flavor, the Kaon can only decay via weak interactions.

$$\begin{aligned}
 (b) \quad P|K^0\rangle &= -|K^0\rangle, & P|\bar{K}^0\rangle &= -|\bar{K}^0\rangle \\
 C|K^0\rangle &= |\bar{K}^0\rangle, & C|\bar{K}^0\rangle &= |K^0\rangle \\
 CP|K^0\rangle &= -|\bar{K}^0\rangle, & CP|\bar{K}^0\rangle &= -|K^0\rangle
 \end{aligned}$$

$$\begin{aligned}
 (c) \quad |K_1\rangle &= \frac{1}{\sqrt{2}}(|K^0\rangle - |\bar{K}^0\rangle) & CP|K_1\rangle &= |K_1\rangle \text{ and} \\
 |K_2\rangle &= \frac{1}{\sqrt{2}}(|K^0\rangle + |\bar{K}^0\rangle) & CP|K_2\rangle &= -|K_2\rangle.
 \end{aligned}$$

with

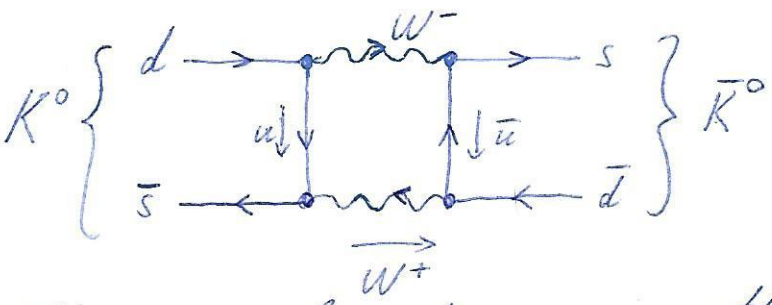
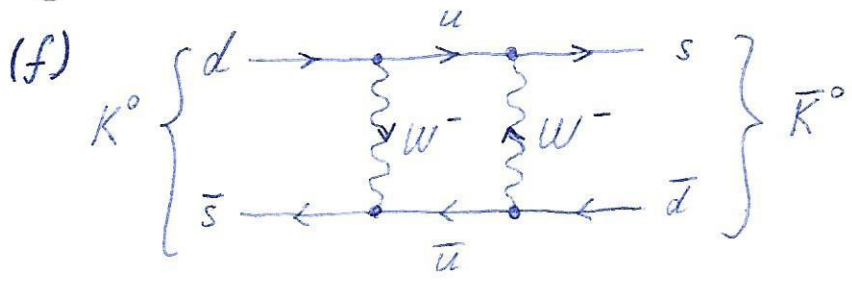
(d) The 2-pion states are parity-even, while the 3-pion states are parity-odd (in the ground state).

$$\Rightarrow CP|\pi\pi\rangle = |\pi\pi\rangle, \quad CP|\pi\pi\pi\rangle = -|\pi\pi\pi\rangle$$

$$\hookrightarrow K_1 \rightarrow 2\pi, \quad K_2 \rightarrow 3\pi$$

So if we observe $K^0 \rightarrow 2\pi$, then actually only the component $|K_1\rangle$ of $|K^0\rangle = \frac{1}{\sqrt{2}}(|K_1\rangle + |K_2\rangle)$ contributes to the decay.

(e) Now the component $|K_2\rangle$ contributes.



\hookrightarrow There are also diagrams with c, t instead of u, which are more suppressed due to smaller CKM-factors.

(g) • The diagonal elements of $H_{weak} = \begin{pmatrix} m - \frac{i}{2}\Gamma & -p^2 \\ -q^2 & m - \frac{i}{2}\Gamma \end{pmatrix}$ describe the quantum mechanical (independent) time-evolution of $|K^0\rangle$ and $|\bar{K}^0\rangle$, with the imaginary piece $\sim i\Gamma$ due to decay.

• The off-diagonal elements $\langle K^0 | H_{weak} | \bar{K}^0 \rangle = -p^2$, $\langle \bar{K}^0 | H_{weak} | K^0 \rangle = -q^2$ are generated by the above loop diagrams.
 ↳ p and q encode the (non-symmetric) dependence on the CKM-elements.

• The eigenvalues of H_{weak} are $\lambda_{\mp} = m - \frac{i}{2}\Gamma \mp pq$, with the corresponding eigenvectors $v_{\mp} = \begin{pmatrix} p \\ \pm q \end{pmatrix}$, so we see that $|K_S^0\rangle$ and $|K_L^0\rangle$ are the eigenstates of H_{weak} .

• We can read off the masses and life-times from λ_{\mp} :
 $K_S^0: m_S = m - \text{Re}[pq], \quad \tau_S = (\Gamma + 2 \text{Im}[pq])^{-1}$
 $K_L^0: m_L = m + \text{Re}[pq], \quad \tau_L = (\Gamma - 2 \text{Im}[pq])^{-1}$

↳ In deriving λ_{\mp} , we chose the signs in $m - \frac{i}{2}\Gamma - \lambda = \pm pq$, such that the imaginary part of $\sqrt{p^2 q^2}$ is positive, so we know that $\tau_S < \tau_L$. Experiment shows that $m_S < m_L$, so $\text{Re}[pq] > 0$.

Problem 2

3/5

① The CKM matrix appears in the w -boson interaction terms, after going from the "gauge basis" to the "mass-basis". The CKM matrix reads

$$V_{CKM} = (L^u)^+ L^d \quad (1)$$

with L^u, L^d the transformation matrices of ν_L and e_L , respectively.

Since $M^{(u)}$ is already diagonal, the CKM-matrix is given by

$$V_{CKM} = L^d \quad (2)$$

To find it, we need to actually diagonalize $M^{(d)}$,

$$M^{(d)} = V_{CKM} M_{\text{diagonal}}^{(d)} V_{CKM}^+ \quad (3)$$

V_{CKM} comprises as columns the normalized eigenvectors of $M^{(d)}$.

A straightforward exercise reveals that

(4/5)

$$V_{CM} = \begin{pmatrix} 0 & \frac{a}{c} & -\frac{b}{c} \\ 0 & \frac{b}{c} & \frac{a}{c} \\ 1 & 0 & 0 \end{pmatrix}, \quad (4)$$

with $c = \sqrt{a^2 + b^2}$.

Comparing (4) to the rotation matrix

$$\begin{pmatrix} 0 & \cos \theta_{mix} & -\sin \theta_{mix} \\ 0 & \sin \theta_{mix} & \cos \theta_{mix} \\ 1 & 0 & 0 \end{pmatrix}, \quad (5)$$

we immediately conclude that

$$\theta_{mix} = \arctan\left(\frac{b}{a}\right). \quad (6)$$

(b) Since the matrix is real, there is no CP violation.

Problem 3

5/5

① once again we have a diagonalized up-quark mass matrix, meaning that V_{CKM} is such that

$$M^{(u)} = V_{CKM} M_{\text{diagonal}}^{(u)} V_{CKM}^{\dagger} \quad (1)$$

We find

$$V_{CKM} = \frac{1}{\sqrt{a^2 + 4b^2}} \begin{pmatrix} 2b & a \\ -a & 2b \end{pmatrix}, \quad (2)$$

meaning that

$$\theta_{\text{mix}} = \arctan\left(\frac{a}{2b}\right) \quad (3)$$

② Since

$$m_s = \lambda_1 \approx 2mb, \quad m_d = \lambda_2 \approx -\frac{a^2 m}{2b} \quad (4)$$

we find

$$\theta_{\text{mix}} \approx \arctan \sqrt{\frac{m_d}{m_s}} \approx 12,5^\circ \quad (5)$$