

Neutrino Physics Course

Lecture XXIV

22/7 / 2022

L MU
Summer 2022



e - μ oscillations?

Oscillations \Leftrightarrow

$$\Delta m^2 \ll \sigma_m^2$$

$$\sigma_m^2 \approx E \sigma_E \approx E \Gamma$$

in order to have $e - \mu$

\Rightarrow large E

- $\bar{W}_R \rightarrow Nl$ $l = e, \mu$

$$N = N_1, N_2$$

$$\left(\begin{matrix} N_1 \\ l_1(\theta) \end{matrix} \right)_R, \left(\begin{matrix} N_2 \\ l_2(\theta) \end{matrix} \right)_R \longleftrightarrow W_R$$

N_1, N_2 = mass eigenstates

$$l_1(\theta) = \cos\theta e + \sin\theta \mu$$

$$l_2(\theta) = -\sin\theta e + \cos\theta \mu$$

$$\Delta m_N^2 > \sigma_m^2 = E \Gamma \simeq M_R \Gamma_R$$

$$\Gamma_R \propto M_R$$

(a) $\boxed{\Delta m_N^2 > \alpha M_R^2}$

$$m_N \propto \vartheta_R$$
$$M_{m_R} \equiv M_R \propto \vartheta_R$$

(b) $t_{coh}(N) \equiv \frac{E^2}{\Delta m_N^2} \frac{1}{\Gamma}$

$$\simeq \frac{M_R^2 \alpha^{-1}}{\Delta m_N^2 M_R} \simeq 100 \frac{M_R}{\Delta m_N^2}$$

$$\Delta m_{\alpha}^2 \simeq M_R^{-2} \Rightarrow$$

$$l_{coh}(N) \simeq l^2/M_R \approx l^{-2} l^{-14} \text{ cm}$$

$\underbrace{\hspace{10em}}$

$$M_R \gtrsim 10^4 \text{ GeV}$$

$\Rightarrow \boxed{N_{1,2} \neq \text{coherent}}$

• Q. Can $l_1(0) \rightarrow l_2(0)$?

$$(a) \quad \Delta m_{e\mu}^2 \simeq m_\mu^2 \ll \Delta m^2$$

$\boxed{m_\mu^2 \ll \alpha M_R^2}$

\mathcal{W}

$$l_{\text{coh}} = \frac{E^2}{\Delta m_{\mu}^2} \frac{1}{\Gamma_R} \propto \frac{\mu_R \alpha^{-1}}{m_\mu^2}$$

$\gtrsim \underline{100 \text{ cm}} \Rightarrow \underline{\text{oscillations}}$

$$\frac{\mu_R \alpha^{-1}}{m_\mu^2} \gtrsim 10^4 \text{ cm}$$

$$\mu_R \gtrsim 10^2 10^{-2} \text{ GeV}^2 \text{ cm}$$



$$(m_\mu \simeq 10^{-1} \text{ GeV})$$

$$\text{GeV cm} \simeq 10^{14}$$



$$M_R \gtrsim 10^{14} \text{ GeV}$$

$$\cdot \tau_\mu \simeq 10^{-6} \text{ sec} \simeq 10^4 \text{ au} = 100 \text{ m}$$

$$\tau_\mu(E) \simeq \frac{E}{m_\mu} \tau_\mu(\text{rest})$$

$$E \simeq M_R \simeq 10^{19} \text{ GeV}$$

$$\Rightarrow \frac{E}{m_\mu} \simeq 10^{15}$$

$$\boxed{\tau_\mu(E=M_R) \simeq 10^{19} \text{ au}}$$

Untangling seesaw

$$M_\nu = - M_D^T \frac{1}{M_N} M_D$$



not complete in LRSM

$$\rightarrow M_N = Y_{\Delta_R}^* v_R \quad M_{w_R} = g v_R$$

$$M_D = Y_0 v_{SM} \quad M_{w_L} = \frac{g}{2} v_{SM}$$

$$\rightarrow v_L \approx \frac{v_{SM}^2}{v_R}$$

induced

$$\begin{aligned} \mathcal{L}_Y (\Delta) &= l_L^T C Y_{\Delta_L} \Delta_L l_L + \\ &+ l_R^T C Y_{\Delta_R} \Delta_R l_R + h.c. \end{aligned}$$

$$\langle \Delta_L \rangle = \begin{pmatrix} 0 & 0 \\ 0 & 0 \end{pmatrix} \Rightarrow$$

$$\boxed{\delta M_\nu = \gamma_{\Delta_L} \Delta_L}$$

(LR)

- P: $\Delta_L \leftrightarrow \Delta_R$

$$\bar{\Phi} \leftrightarrow \bar{\Phi}^+$$

$$\boxed{f_L \leftrightarrow f_R}$$



$$\gamma_{\Delta_L} = \gamma_{\Delta_R}$$

$$\gamma_0 = \gamma_0^+$$

- C: $f_L \rightarrow C \bar{f}_R^T = i \gamma_2 f_R^*$

$$\begin{array}{c}
 \downarrow \\
 \boxed{\begin{array}{l} Y_{D_L} = Y_{D_R}^* \\ Y_D = Y_D^T \end{array}} \quad \left. \begin{array}{l} \\ \end{array} \right\} \underline{\text{verifying!}}
 \end{array}$$

\downarrow neutrino mass

$$M_\nu = Y_{D_L} v_L - M_D^T \frac{1}{M_N} M_D$$

complete neutrino mass

matrix

$$\downarrow LR = C$$

$$M_\nu = \gamma_{\Delta R}^* v_L - M_D \frac{1}{M_N} M_D$$

$$M_N = M_{v_R}^* = \gamma_{\Delta R}^* v_R$$

$$\Downarrow \quad L_R = C$$

$$M_\nu = \epsilon M_N - M_D \frac{1}{M_N} M_D$$

$$\Sigma \equiv \frac{u_L}{v_R}$$

Nemec řek, Tello, G. S.

2012



$$\frac{1}{M_N} M_\nu = \varepsilon - \left(\frac{1}{M_N} M_D \right)^2$$

$$\left(\frac{1}{M_N} \mu_0\right)^2 = \Sigma - \frac{1}{2M_N} \mu_\nu$$

$$M_0 = M_n \sqrt{\varepsilon - \frac{1}{H_N} M_n}$$

(II) $\frac{1}{H_N}$ (I)

$$\Rightarrow \theta = \frac{1}{\mu_N} \mu_0 = \sqrt{\varepsilon - \frac{1}{\mu_N} \mu_0}$$

A simple black line drawing of a curved branch with two small, rounded leaf-like shapes at its tip.

oicill.

$$\boxed{\begin{aligned} \underline{M}_N &= V_R \underline{m}_N V_R^T \\ \underline{M}_D &= V_L^* \underline{m}_D V_L^+ \end{aligned}}$$



Message (LRSM)

SM + ? (minimized)

(a) $SM + V_R$ (Type I
seesaw)
semantics

Why?

$$\underline{M}_v = - \underline{H}_0^T \frac{1}{\underline{M}_n} \underline{M}_D$$

$$M_D = i \sqrt{\mu_1} \circ \sqrt{\mu_2}$$

$$O O^T = I, \quad O \in C$$

$$(b) \quad SH + \Delta_e \equiv \Delta$$

$$l \downarrow \quad \quad \quad \Delta \rightarrow \psi \phi \psi^+$$

predictive

$$\mathcal{L} = \mathcal{L}_{SM} + \mathcal{L}_\Delta + \mathcal{L}_{\Delta,SM}$$



$$\mathcal{L}_D = T_\nu (\partial_\mu \Delta)^+ (D^\mu \Delta^-) - \bar{V}(\Delta)$$

$$\rightarrow \mathcal{L}_{D, \text{sym}} = \underbrace{\ell_L^T C \gamma_\Delta \Delta \ell_L}_{\text{---} \uparrow \text{---}} + h.c.$$

(1)

$$\gamma(\Delta) = -2 \gamma(\ell_L) = +2$$

$$\gamma(\phi) = +1$$

$$+ \overline{\phi}^T \cancel{\Delta} \cancel{\Delta} \phi$$

$$+ u \phi^T i \sigma_2 \Delta^* \phi + h.c.$$

$$\gamma: \begin{matrix} 1 & -2 & 1 \end{matrix} = 0$$

π

$$T_\nu \Delta_L^+ \cancel{\Phi} \langle D_R \rangle \cancel{\Phi}^+$$

$$V(\Delta) = \bigoplus M_\Delta^2 T_r \Delta^+ \Delta^- + \dots$$

$$+ \mu \phi^\top i \sigma_2 \Delta^* \phi + h.c.$$

$$\Delta = \begin{pmatrix} \delta^+ & \boxed{\delta^{++}} \\ \delta^0 & -\delta^+ \end{pmatrix}$$

$$\Downarrow \quad \phi = \begin{pmatrix} v^0 \\ v^+ \end{pmatrix}$$

$$V(\delta^0) = + M_\Delta^2 \delta^0 \delta^0 * +$$

$$+ \mu v^2 \delta^0 * + h.c.$$

$$\frac{\partial V}{\partial \delta^0 *} = \mu v^2 + M_\Delta^2 \delta^0 + \dots = 0$$



$$\langle \delta_0 \rangle = \mu \frac{v^2}{m_\Delta^2}$$

$$m_D \rightarrow \infty \Rightarrow \langle \delta_0 \rangle \rightarrow 0$$

From (1) $\Rightarrow M_J = \gamma_\Delta \langle \delta^0 \rangle$

Type II seesaw

Maeg, Wittenberg

Lazarides, Shatzi,

Helbigstr, G.S.



$$\delta^{++} (e_i)^T C (\gamma_0)_{ij} e_j =$$

$$= \delta^{++} e_i^T C \frac{(-M_\nu)_{ij}}{\langle \Delta \rangle} e_j$$

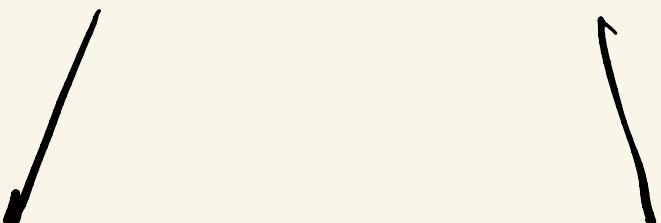


$$\delta^{--} \rightarrow e_i e_j$$

$$\Gamma \propto (M_\nu)_{ij}$$

Seesaw

$$M_0 = \gamma_0 v$$



↓ Type I

$$M_\nu \sim \gamma_D^2 \frac{v^2}{m_N}$$

↓ Type II

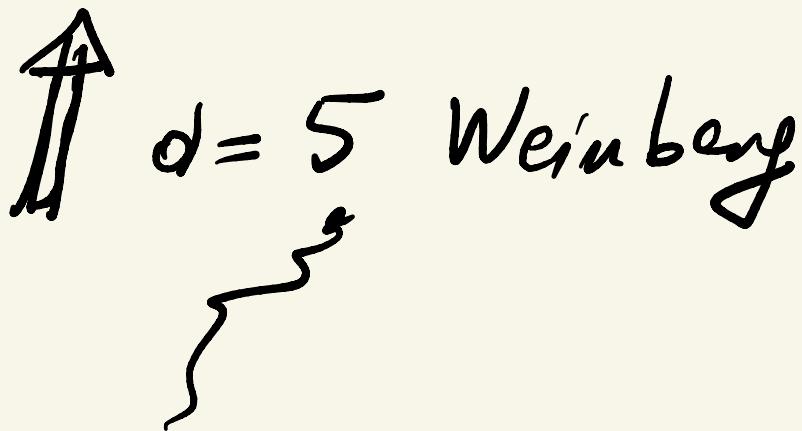
$$M_\nu \propto \gamma_D \mu \frac{v^2}{m_\Delta^2}$$

$$(\mu \sim m_\Delta)$$

$$\gamma_D \frac{v^2}{m_\Delta}$$

Common factor

$$\sim \frac{v^2}{M_{\text{new}}} \sim \frac{M_W^2}{M_{\text{new}}}$$



$$\frac{(\ell^\top i \sigma_2 \phi) c (\phi^\top i \sigma_2 \ell)}{\underline{M_{\text{new}}}}$$

$$\Rightarrow M_\nu \simeq \frac{v^2}{M_{\text{new}}} \simeq \frac{M_W^2}{M_{\text{new}}}$$

Heed 2022



Type II reason

→ CDF W-mass derivation

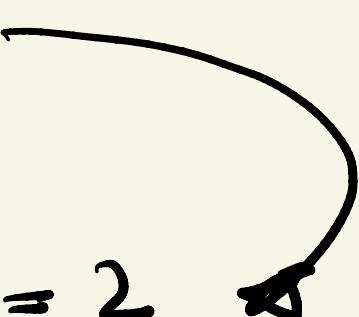
CDF = W-mass deviation

$$\rho_{SM} = \left(\frac{M_Z \cos \theta_W}{M_W} \right)^2 = 1 \quad (\text{tree})$$

$$+ \frac{\alpha}{\pi} \frac{M_Z^2 - M_W^2}{M_W^2} + (\text{O}(2))$$

$$\text{CDF} \Rightarrow \rho \neq \rho_{SM}$$

Type II : $\langle \Delta \rangle$

$$\left(\frac{M_Z \cos \theta_W}{M_W} \right)^2 = 2$$


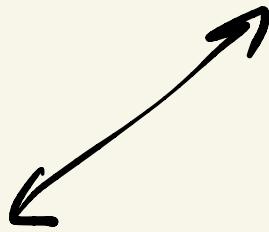
$$P = P_{SM} + \mathcal{O}\left(\frac{\langle\Delta\rangle}{\alpha}\right)^2 + \dots$$

loops

if this was all

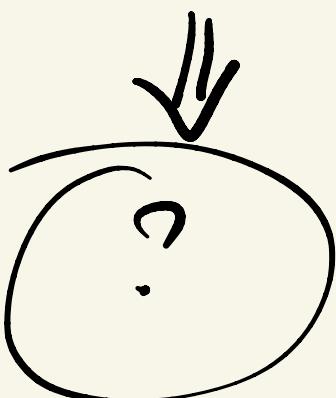
$$\Rightarrow \langle\Delta\rangle \simeq 5 \text{ GeV}$$

but



$$\delta P_{loop} = \frac{\alpha}{\pi} \frac{M_{f++}^2 - M_{f+-}^2}{M_W^2} + \dots$$

both loops and $\langle\Delta\rangle$



\downarrow ? Heesch

$$\langle \phi \rangle \rightarrow e^- \bar{\nu}^- (\text{GeV})$$

$$M_\phi = y_\phi \langle \Delta \rangle$$

- if $\langle \Delta \rangle \approx 6\text{eV} \Rightarrow y_\phi \approx 10^{-9}$
- if $\langle \Delta \rangle \ll \text{MeV} \Rightarrow y_\phi = \dots ?$



observable

$\langle \Delta \rangle \neq \text{Higgs}$

• Higgs : $m_h \simeq \langle \phi \rangle$
 $- \mu^2 \phi^2 + \dots \xrightarrow{\text{!}} \langle \phi \rangle \neq 0$

• $\langle \Delta \rangle = \text{anti-Higgs}$

$$m_\Delta \gg m_W \Rightarrow \langle \Delta \rangle \ll v_{SM}$$

$$\begin{aligned}
 D_\mu \langle \Phi \rangle &= \partial_\mu - ig T_L W_L \langle \phi \rangle \\
 &\quad + ig \langle \phi \rangle T_R W_R
 \end{aligned}$$

$$M_{W_R}^2 = \langle \Delta_R \rangle^2 + \langle \bar{\Phi} \rangle^2$$

$W_L - W_R$ mix

$$\beta_{LR} \propto \frac{\langle \phi \rangle^2}{\langle \Delta_R \rangle^2} \propto \frac{M_{W_L}^2}{M_{W_R}^2}$$

$$M_{W_R} \gtrsim 5 \text{ TeV}$$

$$M_{W_L} \simeq 80 \text{ GeV}$$

$$\beta_{LR} \leq 3 \times 10^{-4}$$

↑

$$\begin{array}{c} M_L \left(M_L^2 \simeq \langle \phi \rangle^2 \right) \\ M_R \left(\langle \phi \rangle^2 \quad M_R^2 \right) \end{array}$$
$$\Rightarrow Z_{LR} \simeq \frac{\langle \phi \rangle}{M_R^2}$$