FROM UNIVERSE...

TO PLANETS LECTURE 1.1: COSMOLOGICAL HISTORY

OUTLINE

- Cosmological history:
 - From Big Bang to present day
- Stellar Nurseries:
 - The environment in which stars are born
- Star formation and evolution:
 - The life and death of stars



PREVIOUSLY.

Dark Energy Accelerated Expansion



VERY EARLY UNIVERSE



Afterglow Light Pattern 380,000 yrs.

Inflation

^{ars} Quantum Fluctuations

SEPARATION OF FORCES



WHAT IS INFLATION?

Had inflation not taken place, the present-day observable universe would have had to have been relatively large just after the Big Bang.

Once the inflationary epoch had ended, the universe continued to expand in a more gradual way down to the present day.

 10^{40} 10^{30} 10^{20} 10^{10} Distance (cm) In the inflationary model, the present-day 10-10 observable universe was very tiny just after 10-20 the Big Bang. This region, as well as the rest 10-30 of the universe, then underwent a tremendous 10-40 expansion during the inflationary epoch. 10-50 10-60 Inflationary epoch 10^{15} 10-45 10-35 10^{-25} 10-15 10^{-5} 10⁵ Present Time after Big Bang $(s) \longrightarrow$

WHAT CAUSED INFLATION?



WHAT CAUSED INFLATION?



WHAT CAUSED INFLATION?





 $-500 \,\mu\mathrm{K}$

check this out: <u>http://thecmb.org/</u>

500 µK

Especially when the oldest light is the cosmic microwave background (CMB) from 380,000 years after the Big Bang...



INFLATION: SOLVES THE FLATNESS PROBLEM...



INFLATION: SOLVES THE ISOTROPY PROBLEM...



Radiation from A takes 13.7 billion years to reach us Radiation from *B* takes 13.7 billion years to reach us

INFLATION: SOLVES THE ISOTROPY PROBLEM...



Radiation from A takes 13.7 billion years to reach us Radiation from *B* takes 13.7 billion years to reach us

INFLATION: EXPLAINS WHERE MATTER COMES FROM...

$$\Delta E \Delta t = \frac{\hbar}{2} \left\{ \longrightarrow \Delta t = \frac{1}{\Delta m} \frac{\hbar}{2c^2} \xrightarrow{e^+} 3.22 \times 10^{-22} \text{ s} \\ E = mc^2 \right\}$$



appear anywhere in space...

...but must disappear after a very short time interval.

INFLATION: EXPLAINS WHERE MATTER COMES FROM...

A virtual particle-antiparticle pair that appears just before the inflationary epoch...

Virtual positron

...is pulled so far apart during inflation that the particles cannot recombine, thus leaving a pair of real particles.



Real positron

Virtual electron



(a) Pair production

(b) Annihilation

- Thermal equilibrium between radiation and matter (Pair production rate = Annihilation rate)
- Expansion redshifted photons, lowering the temperature

EXPANSION AND ANNIHILATION REMOVE ANTIMATER

- Expansion
 - \rightarrow longer λ
 - \rightarrow lower energy
 - \rightarrow lower temperature



- Insufficient energy to create certain particles
 - \blacktriangleright \rightarrow Annihilation events leaving only matter
 - Particle ratio: 10^9 antimatter to $10^9 + 1$ ordinary matter

AND EVENTUALLY SETS ABUNDANCES



Afterglow Light Pattern 380,000 yrs.

Inflation

^{ars} Quantum Fluctuations



Electron-positron annihilation severely reduces the number Time after Bisof electrons to form neutrons

n 10^{-10} Neutrons decay into protons with a half-life of 10.5 minutes 10^{-4} $n \rightarrow p + e^- + \bar{\nu}$ \rightarrow n + ν p + e**Collisions keep number of neutrons** n approximately equal to protons 0.3×10^{9} 3.0×10^{9} 1.0×10^{9} 0.1×10^{5}

Temperature (K)



Temperature (K)



Temperature (K)



NUCLEOSYNTHESIS: VIRTUALLY NO HEAVY ELEMENTS



NUCLEOSYNTHESIS: OBSERVATIONS

- Observations at later times are consistent with models with low baryonic density
- However, the universe requires high densities to explain flatness
 - Most matter in the universe is non-baryonic (dark)!
 - Important for creating the large-scale structure we see today



NUCLEOSYNTHESIS: OBSERVATIONS

- Observations consistent with models if baryon density is low
- Universe requires high densities to explain flatness
 - Suggests most matter in the universe is non-baryonic!
 - "Dark" because it does not emit or interact with light
 - Can begin collapsing before regular matter → large-scale structure we see today



DARK MATTER SIMULATION



HUBBLE ULTRA DEEP FIELD



Spitzer "First Light"



Spitzer "First Light"



Spitzer "First Light"





Bark Bages First Light

FIRST STARS

- With no heavy elements, the internal pressure of gas clouds would be higher
- Collapse requires greater mass
 - First stars: $30 1000 \text{ M}_{\odot}$
 - Short lifetimes with large supernova explosions → heavy elements
- Surface temperatures: 10^5 K
 - Flood of short wave radiation
 - Epoch of reionisation

yrs 400 million yrs



BIG BANG SUMMARY



Distance (light years)

BIG BANG SUMMARY: GATHERING THE INGREDIENTS



If you wish to make an apple pie from scratch, you must first invent the Universe.
BIG BANG SUMMARY: GATHERING THE INGREDIENTS

- Although instead of apple pie, we are more interested in the ingredients needed for life
 - Human composition (by number):

н

Li

3

Na

11

K

19

Rb

37

Cs

55

Fr

87

H : 63 % O : 25.5 % C : 9.5 % N : 1.4 % other : 0.6 %





BIG BANG SUMMARY: GATHERING THE INGREDIENTS

- We need a hot, dense beginning to generate the elemental abundance we see in the universe today
- We need the universe to rapidly expand through inflation (not just regular expansion)
 - To create the space between astronomical objects today (crowded environments are anyway not suitable for life)
 - To cool down: energy \longrightarrow particles \longrightarrow structures
 - To inflate tiny quantum fluctuations into large scale structures that can later collapse due to gravity

THE INTERSTELLAR MEDIUM AND GIANT MOLECULAR CLOUDS





Diffuse Cloud (Giant Molecular Cloud)

Diffuse ISM

Planetary System

Dense Cloud (Molecular Core)

Accretion Disc

. .

INTERSTELLAR MEDIUM (ISM)

INTERSTELLAR MEDIUM (ISM)

INTERSTELLAR MEDIUM (ISM)



INTERSTELLAR MEDIUM (ISM): COMPOSITION AND PHASES

- Hydrogen is the dominant element in the ISM followed by helium and other 'metals'
- Gas in the ISM comes in several different phases depending on the temperature and density

Element	Abundance/H
He	0.1
С	3.6×10-4
Ν	1.1×10-4
0	8.5×10-4
Si	3.6×10 ⁻⁵

Phase	n (cm-3)	Т (К)	M (10 ⁹ M⊙)	Filling factor	Pressure (~nT)	
molecular	>300	10	2	0.01	3000	Sites of star formation
cold atomic	50	80	3	0.04	4000	Most mass in HI
warm atomic	0.5	8×10 ³	4	0.3	4000	Sets the ISM pressure
warm ionised	0.3	8×10 ³	1	0.15	2400	
hot ionised	3×10-3	5×10 ⁵	tiny	0.5	1800	
HII regions	1 - 104	105	tiny	tiny	10 ⁵ - 10 ⁹	stars

Despite the huge range of conditions, most phases are in approximate pressure balance

INTERSTELLAR MEDIUM (ISM): COMPOSITION AND PHASES



GIANT MOLECULAR CLOUDS (GMC)



https://www.youtube.com/watch?v=gpZtlAPQ2QU

GIANT MOLECULAR CLOUDS (GMC)



MOLECULAR GAS TRACER: CO

- Because H₂ is a symmetric molecule, it has zero dipole moment in the ground state
- It cannot radiate unless excited, but molecules are only formed in cold, shielded regions so gas is usually in the ground state

Molecular Hydrogen cannot be observed directly in cold clouds

- We can instead observe molecular CO, which has a permanent dipole moment
- Requires less energy to excite rotational and vibrational transitions (sometimes rovibrational) of a non-symmetric molecule

Assume that the ratio of CO to H₂ is constant! (not always true at very high densities)



GALACTIC CO DISTRIBUTION

- Observations show that most molecular gas is found in molecular clouds that are concentrated in spiral arms located in the galactic plane
- Total mass of molecular hydrogen: $\sim 2-3 \times 10^9 M_{\odot}$
- This is ~25 50% of the total gas mass in the galaxy
- Molecular gas is destroyed by photodissociation (UV radiation and cosmic rays) and dissociative recombination (reaction with electrons)
- Need a high column density ($\gtrsim 10 M_{\odot} \text{ pc}^{-2}$) of atomic gas to shield molecular hydrogen





THE FORMATION OF MOLECULES

- Two-body collisions of atomic H is too inefficient to explain the amount of molecular hydrogen in the Milky Way, even in the denser spiral arms
- The surface of dust grains provides a reservoir where atoms/molecules can be stored and brought together on longer timescales than in the gas
 - Enables reactions that are too slow in the gas (e.g. hydrogenation of atomic O, C and N to form H₂O, CH₄ and NH₃)
- Liberated energy helps to release the molecule from the dust surface







INCERSCELLAR GRAIN SURFACE CHEMISCRY



INTERSTELLAR MEDIUM (ISM): DUST

- Dust comprises only ~1% of the ISM by mass, but is an important source of opacity and cooling
- Average interstellar extinction can be satisfactorily reproduced with two components (silicate and graphite) and a power-law grainsize distribution:

$$dn(a) \propto a^{-3.5} da \qquad \begin{cases} a_{\min} = 0.005 \ \mu n \\ a_{\max} = 0.25 \ \mu m \end{cases}$$

Turbulence influences spatial and size distribution





INTERSTELLAR MEDIUM (ISM): DUST





THERMAL PHYSICS OF ISM: HEATING

- Heating in the diffuse ISM is dominated by photoionisation (UV, X-ray, cosmic ray) of dust (photoelectric effect) and gas
 - ~100 times more electrons from dust (despite dust-to-gas ratio being 1:100)
- The electron receives a large fraction of the photon energy
- Subsequent collisions of the electron with other atoms redistributes the energy (i.e. thermalisation)







THERMAL PHYSICS OF ISM: COOLING



 Cooling is usually dominated by collisional de-excitation

radiat

Fine structure lines

smaller velocity

after collision

- Electron recombination lines
- Resonance and metastable lines

thermalisation through

elastic collisions

- Cooling rate, Λ , per n^2
- $T < 10^4$ K, excitation of CII and OI dominate cooling
- The step at 10⁴ K corresponds to cooling by the Lyman-alpha line of H
- T > 10⁴ K, cooling due to various lines and bremsstrahlung

COLLAPSE OF MOLECULAR CLOUDS

- Molecular clouds form when a dense enough region of gas forms in the ISM to become optically thick from UV radiation (which would otherwise destroy the molecules)
- This results in a "UV shield" of atomic (or ionised) gas
- Molecules then form on dust grains and are ejected into the cloud
- The cooling is often dominated by dust emission in the infrared
- The cores of these clouds are extremely cold (~10 K) allowing gravity to overpower internal gas pressure



SIMPLEST CASE: FREEFALL

- Consider the gravitational collapse of a cold (isothermal), pressureless fluid of uniform density ρ_0 and radius R_0 initially at rest
- The gravitational acceleration at any point is:

$$a_{\rm grav} = \frac{d^2r}{dt^2} = -\frac{GM(R_0)}{r^2}$$

• Multiplying by v_r and integrating we find

$$\frac{dr}{dt} = -\sqrt{\frac{8\pi}{3}}G\rho_0 R_0^2 \left(\frac{R_0}{r} - 1\right)$$

• Making substitutions $\theta = r/R_0$ and $\chi = (8\pi G\rho_0/3)^{1/2}$ and integrating again gives

$$\sqrt{\theta(1-\theta)} + \cos^{-1}\sqrt{\theta} = \chi^{2}$$

 $M(R_0)$ mass constant

 $v_r |_{t=0}$

Rewriting the mass in terms of density and multiplying by v_r

$$\frac{dr}{dt}\frac{d^2r}{dt^2} = -\left(\frac{4\pi}{3}G\rho_0 R_0^3\right)\frac{1}{r^2}\frac{dr}{dt}$$

Integrating with respect to time

$$\frac{1}{2}\left(\frac{dr}{dt}\right)^2 = \left(\frac{4\pi}{3}G\rho_0 R_0^3\right)\frac{1}{r} + C$$

The integration constant is fixed by the initial condition $\dot{r} = 0$ at $r = R_0$ $C = -\frac{4\pi}{3}G\rho_0 R_0^2$

$$\frac{dr}{dt} = -\sqrt{\frac{8\pi}{3}}G\rho_0 R_0^2 \left(\frac{\kappa_0}{r} - 1\right)$$

• Making substitutions $\theta = r/R_0$ and $\chi = (8\pi G\rho_0/3)^{1/2}$ and integrating again gives

$$\sqrt{\theta(1-\theta)} + \cos^{-1}\sqrt{\theta} = \chi^{2}$$

– mass constant

SIMPLEST CASE: FREEFALL

This predicts that M(R₀) collapses to infinite density (i.e. r = 0) on the free-fall timescale

$$t_{\rm ff} = \sqrt{\frac{3\pi}{32G\rho_0}}$$

- Independent of initial radius so the density profile increases at the same rate everywhere (homologous or self-similar)
- Neglects pressure, rotation, turbulence, magnetic fields, and radiation
- For centrally condensed clouds, the freefall time is shorter near the centre (insideout collapse)

$$t_{\rm ff} = 2.1 \times 10^3 \text{ s} \left(\frac{\rho}{\text{g cm}^{-3}}\right)^{-1/2}$$
$$= 50 \text{ Myr} \left(\frac{n}{\text{cm}^{-3}}\right)^{-1/2}$$

Estimated	ISM	Freefall	Times

Phase	<i>n</i> (cm ⁻³)	t _{ff} (Myr)
molecular	>300	< 3
cold atomic	50	7
warm atomic	0.5	70
warm ionised	0.3	90
hot ionised	3×10-3	900

FLUID DYNAMICS

- Molecular clouds are complex, non-linear systems due to the large number of different physical processes that happen simultaneously
- If we wish to study the stability of the gas and to know what happens under certain physical conditions, we need to solve them accurately
- If we only consider hydrodynamics and gravity...it's already complicated

$$\frac{\partial \rho}{\partial t} + \nabla(\rho \mathbf{v}) = 0$$

$$\frac{D\mathbf{v}}{Dt} = \frac{\partial \mathbf{v}}{\partial t} + \mathbf{v} \cdot \nabla \mathbf{v} = -\frac{\nabla P}{\rho} - \mathbf{g}$$
$$\frac{Du}{Dt} = \frac{\partial u}{\partial t} + \mathbf{v} \cdot \nabla u = -\frac{P}{\rho} \nabla \cdot \mathbf{v}$$
$$P = (\gamma - 1)\rho u$$

The Equation of Continuity (or Mass Conservation)

The Momentum Equation (or the Euler Equation)

The Energy Equation (or Energy Conservation)

Equation of State (EOS)

SOUND WAVES

- Solution to the fluid equations for the simplest possible scenario
 - 1D hydrodynamics (NO gravity) with isothermal EOS ($P = c_s^2 \rho$)
- A trivial solution can be obtained with constant values of density, velocity and pressure, i.e. ρ_0 , v_0 and P_0
- If we perturb the solution we get sound waves, something we actually saw earlier...

500 µK

THE JEANS INSTABILITY

- Now consider a cloud where the pressure gradient acts against the attractive force of gravity
- For simplicity, assume an isothermal EOS, i.e. $P = c_s^2 \rho \equiv \frac{k_B T_0}{\sqrt{1-\rho}} \rho$

$$\nabla P \sim \frac{P_{\text{centre}} - P_{\text{edge}}}{R_0} \sim \frac{c_{\text{s}}^2 \rho_0}{R_0} \qquad \qquad \begin{array}{c} \text{acceleration} \\ & \swarrow a \sim \frac{c_{\text{s}}^2}{R_0} - \frac{4\pi}{3} G \rho_0 R_0 \end{array}$$

For small values of R_0 (very little mass), the pressure term dominates

$$a \sim \frac{c_{\rm s}^2}{R_0}$$
 $t \sim \left(\frac{R_0}{a}\right)^{1/2} \sim \frac{R_0}{c_{\rm s}}$

Cloud expands and dissolves on sound-crossing timescale

Conversely for large values of R₀ (large mass), the gravitational term dominates

$$a \sim -\frac{4\pi}{3}G\rho_0 R_0$$
 $t \sim \left(\frac{3}{4\pi G\rho_0}\right)^{1/2} \sim (G\rho_0)^{-1/2}$

Cloud collapses on freefall timescale

THE JEANS LENGTH AND JEANS MASS

At what length scales do we get gravitational collapse, i.e. a < 0?

$$\frac{c_{s}^{2}}{R_{0}} - \frac{4\pi}{3}G\rho_{0}R_{0} \lesssim 0 \qquad R_{0} \gtrsim R_{J} \sim \left(\frac{c_{s}^{2}}{G\rho_{0}}\right)^{1/2} \quad \text{Jeans Length}$$
• We can also define the mass: $M_{0} \gtrsim M_{J} \sim \left(\frac{c_{s}^{6}}{G^{3}\rho_{0}}\right)^{1/2} \quad \text{Jeans Mass}$
• The collapse timescale varies as t_{col}

$$t_{col} \sim t_{ff} \left[1 - \left(\frac{M_{0}}{M_{J}}\right)^{-2/3}\right]^{-1/2} \quad t_{ff}$$

HIERARCHICAL FRAGMENTATION

- Imagine an infinitely large cloud with small density fluctuations on various size scales
- Smaller fluctuations will "dissolve" away as sound waves
- Larger fluctuations will collapse without affecting the larger cloud
- For sufficiently small densities, the cloud remains isothermal where $M_{\rm J} \propto \rho^{-1/2}$
- The fragmentation scale gets arbitrarily small as density increases
- Originally thought to be the origin of clusters, but this neglects heating



A MINIMUM MASS FOR STAR FORMATION?

- As the cloud gets denser it eventually becomes optically thick to its own cooling radiation from dust
- The Jeans mass now increases with increasing density (approximately adiabatic contraction)
- The intersection of these curves gives a minimum Jeans mass



STAR FORMING CORES

- At small scales (~0.1 pc), the gas has condensed into star-forming cores
- Each core is the progenitor of either an individual star or a binary/multiple system
- Initially, the core can only be seen in the sub-mm
- As it forms a central protostar, it starts to emit at shorter and shorter wavelengths until eventually becoming a star



GIANT MOLECULAR CLOUD STABILITY

- Collapse triggered by:
 - Cloud-cloud collisions (including galactic collisions)
 - Shocks from nearby supernova explosions
 - Passage through a spiral arm of the galaxy
 - Collapse is facilitated by low temperatures and high densities (e.g. regions full of dust)
- Collapse is hindered by:
 - Turbulence
 - Macroscopic flows
 - Cloud geometry
 - Rotation
 - Magnetic fields



STAR FORMATION: COLLAPSE

30560 yrs

1000 AU

STAR FORMATION: COLLAPSE

STAR FORMATION: COLLAPSE

- Conservation of angular momentum leads to rapid rotation in the central region, suppressing further collapse and promoting the formation of a circumstellar disc
- The magnetic field can effectively transfer the excess angular momentum from the center of the cloud by magnetic braking and jets, thus promoting further collapse
 - If too efficient, magnetic braking can prevent a circumstellar disc from even forming (labeled the magnetic braking catastrophe)
- Observations show evidence for both magnetic fields and circumstellar discs, so the "catastrophe" is only numerical

STAR FORMATION: JETS

- Outflows of ionised matter emitted along the axis of rotation.
- Dynamic interactions between compact central objects (e.g. stars and black holes) and a surrounding accretion disc.
- Almost always associated with magnetic fields that twist up and collimate the beam.
 HH-30
- Ionised particles are attached to the field lines like "beads on a wire" and are centrifugally accelerated.

STAR FORMATION: JETS



Extended jet component

ω




Extended jet component

ω







HH-47

HH-111





STAR FORMATION: CIRCUMSTELLAR DISCS



STAR FORMATION: CIRCUMSTELLAR DISCS



































STELLAR EVOLUTION LECTURE 1.3

BIRTH, LIFE, AND DEATH

STAR FORMATION: PROTOSTARS

- Core that is still gathering mass from its parent molecular cloud (envelope). Material initially falls directly onto the protostar, but this later transitions to the surrounding disc.
- Gravitational contraction releases energy. About half the energy is radiated away by photons, the other half goes into heating the interior.
- The ratio of the energy produced to the energy lost is known as the Kelvin-Helmholz timescale:

$$\tau_{\rm KH} = \frac{\text{gravitational binding energy}}{\text{luminosity}} = \frac{GM^2}{RL}$$

The time spent in the gravitational contraction phase depends on the mass of the protostar.

The gravitational binding energy can be estimated by considering a sphere of radius R and imagine pulling it apart successively by removing thin shells to infinity.

 $m_{\rm shell} = 4\pi r^2 \rho dr$ and $m_{\rm interior} = \frac{4\pi}{2} r^3 \rho$

The energy required for removing one shell is the negative gravitational potential energy $dU = + G \frac{m_{\text{shell}} m_{\text{interior}}}{r}$ Integrating over all shells (assuming ρ is constant): $U = G \frac{16\pi^2}{3} \rho^2 \int_0^R r^4 dr = G \frac{16\pi^2}{15} \left(\frac{M}{\frac{4\pi}{3}R^3}\right)^2 R^5 \sim \frac{GM^2}{R}$

STAR FORMATION: PROTOSTARS

- Cores start out (relatively) low density and transparent. They shine because they are hotter than their surroundings, but are too cool for nuclear fusion.
 - Photons can leak out, keeping the interior cool, allowing collapse to continue. The surface area and luminosity decrease, but the surface temperature remains almost constant.
- Eventually the protostar builds up enough density and it becomes opaque (i.e. photons become trapped). The pressure builds, and hydrostatic equilibrium is reached (but not thermal equilibrium).
 - Changes in the star now keep the luminosity roughly constant, but the surface temperature starts to build up.

PRE-MAIN-SEQUENCE PHASE



HAYASHI TRACKS



- Near vertical track in the pre-main-sequence phase for $M \lesssim 3 M_{\odot}$
- Convective interior:
 energy released by
 contraction on the
 interior is convected to
 the the surface (nearly
 constant in time)
- The loss of energy at the surface allows the star to further collapse
- Smaller radii with a constant temperature produces smaller luminosities

HENYEY TRACKS



- Near horizontal track in the pre-main-sequence phase for $0.5~M_{\odot} \lesssim M \lesssim 10~M_{\odot}$
- Radiative interiors:
 energy loss limited by
 radiative diffusion
 (takes time for energy
 to diffuse out from the
 interior)
 - Steady increase in surface temperature
- Slow contraction with
 increasing surface
 temperature leads to a
 near constant
 luminosity with time

THREE POTENTIAL OUTCOMES

- 1. Brown dwarfs: below ~ 0.08 $M_{\rm sun}$ the core never gets hot enough to ignite H fusion
 - May fuse deuterium (²H) and lithium (⁷Li) if mass is > 65 $M_{\rm J}$.
- 2. Zero-age main sequence (ZAMS): gravitational collapse continues until core reaches ~10 million K
 - Core temperature and pressure rise
 - Pressure = gravity and core collapse halts
 - Energy created by P-P chain fusion = luminosity
- 3. Eddington limit: luminosity so large that radiation pressure pushes away the gas (via the dust)
 - Difficult to make stars $> 20 M_{\odot}$
 - Most massive stars ever observed $\sim 150 300 M_{\odot}$



Surface Temp:

1500K - 3000 K

T Dwarfs > 500 Known Surface Temp: 500 K - 1500 K

Y Dwarfs ~17 Known Surface Temp: 250 K - 500 K





STELLAR EVOLUTION: MAIN SEQUENCE

Stars are classified by their spectra (the elements that they absorb) and their temperature. Seven main types of stars.
 In order of decreasing temperature, O, B, A, F, G, K, and M.



spectral Type	Example(s)	Temperature Range	Key Absorption Line Features	Brightest Wavelength (color)
0	Stars of Orion's Belt	>30,000	Lines of ionized helium, weak	<97 nm (ultraviolet)*
В	Rigel	30,000 K-10,000 K	Lines of neutral helium, moderate hydrogen lines	97–290 nm (ultraviolet)*
А	Sirius	10,000 K-7,500 K	Very strong hydrogen lines	290–390 nm (violet)*
F	Polaris	7,500 K–6,000 K	Moderate hydrogen lines, moderate lines of ionized calcium	390–480 nm (blue)*
G	Sun, Alpha Centauri A	6,000 K–5,000 K	Weak hydrogen lines, strong lines of ionized calcium	480–580 nm (yellow)
К	Arcturus	5,000 K-3,500 K	Lines of neutral and singly ionized metals, some molecules	580–830 nm (red)
м	Betelgeuse, Proxima Centauri	<3,500 K	Molecular lines strong	>830 nm (infrared)

* All stars above 6,000 K look more or less white to the human eye because they emit plenty of radiation at all visible wavelengths.

HERTZSPRUNG-RUSSELL DIAGRAM



- About 90% of stars are on the main-sequence (most of those are red dwarfs smaller than sun). Energy comes from nuclear fusion, as they convert H→He.
- The mass of the star determines the internal structure and, ultimately, how the star will die.



 Below ~ 0.25 M_{sun} convection dominates, the surface temperatures and luminosities gradually increase, and the star quietly transitions into becoming a white dwarf.

- Higher masses develop an inert He core (not hot enough for fusion). The core slowly contracts and heats up, causing intense H burning in a shell surrounding the core.
- Increased radiation pressure causes the surface to expand $10-100 \times$ its original size and cool down, making it appear more red (hence the name Red Giant Branch or RGB).
- The core also contains free electrons. At some point the Pauli exclusion principle (electrons cannot occupy the same energy level with identical quantum numbers) prevents further gravitational collapse due to electron degeneracy pressure.
- When degeneracy pressure is stronger than thermal pressure, the gas is said to be degenerate and no longer behaves like an ideal gas (pressure no longer depends on temperature).

- Nuclear reactions in degenerate cores tend to be explosive because the gas does not expand and cool down when the temperature increases.
- However, nuclear reactions are still very temperature dependent. Runaway burning ensues until the core temperature lifts the degeneracy.
- The ignition of He in low-mass stars is explosive (called a helium core flash). Large amounts of He fuses to C in a matter of seconds. After the flash, the luminosity decreases and the outer layers of the star shrink.
- The process repeats, expanding along the Asymptotic Giant Branch (AGB), but cannot ignite C burning. He shell burns in a series of flashes, shedding its outer layers (planetary nebuae).





DUST FORMATION

- The ISM is not dense enough to form complex interstellar dust molecules efficiently
- The interior of stars are too hot and destroy complex molecules
- Dust condenses out of stellar outflows:
 - Red giant and AGB stars
 - Wolf-Rayet stars
 - Supernovae
 - Neutron star mergers
- Compare the spectrum of red giant atmospheres to PAHs in the lab





STELLAR EVOLUTION: AGB STARS



R SCULPTORIS

STELLAR EVOLUTION: HIGH-MASS STARS

- Massive stars continue nuclear burning beyond
 C, but the process becomes progressively less efficient until Fe, which requires more energy than it releases
- At this point, the internal structure is made up of concentric layers of different elements
- If the mass of the iron core exceeds the Chandrasekhar limit (maximum mass for an electron-degenerate system; ~ 1.4 M_{sun}), the star collapses, rebounds, and explodes in a core-collapse or Type II supernova. The envelope is ejected with typical speeds of about 10⁴ km/s (~0.03c).
- At peak brightness (lasting weeks), their luminosities can compete with the luminosity of their entire host galaxy
- Supernovae are responsible for a large fraction of the ionised gas and heavy elements in the ISM and regulate star formation rates in galaxies



SUPERNOVAE AND PLANETARY NEBULAE



IN SUMMARY...



MAIN POINTS

- H and He are by far the most prevalent elements in the Universe (provides the fuel for sustaining life, i.e. stars), but there is still a rich amount of chemistry to consider
- The life cycle of stars is crucial to planet formation
 - Formation: concentrate gas/solids in accretion discs where planets are formed
 - Destruction: produce the building blocks for future planets (eject solids/heavy elements into the ISM)
- Environment is key! Location in the galaxy, what material is available, neighbouring stars, etc. all play an important role in planet formation