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## Exercises for Quantum Field Theory (TVI/TMP) Problem set 3 Feynman path integral and harmonic oscillator

## 1 Gaussian integrals and Wick theorem

(i) Prove the Gaussian integration formula

$$\int_{\mathbb{R}^n} d^n x \exp\left[-\frac{1}{2}x^T A x + b^T x + c\right] = \sqrt{\frac{(2\pi)^n}{\det A}} \exp\left[\frac{1}{2}b^T A^{-1}b + c\right]$$
(1)

valid if A is a positive definite symmetric  $n \times n$  matrix and b is an arbitrary real vector.

(ii) Calculate the two-point function

$$\int_{\mathbb{R}^n} d^n x x_{j_1} x_{j_2} \exp\left[-\frac{1}{2}x^T A x\right]$$

and the four-point function

$$\int_{\mathbb{R}^n} d^n x x_{j_1} x_{j_2} x_{j_3} x_{j_4} \exp\left[-\frac{1}{2} x^T A x\right].$$

Interpret your result using Wick's theorem.

(iii) (\*) Consider one-dimensional Gaussian integral (1) with n = 1. What are the coefficients of Taylor expansion of the RHS counting combinatorially?

## 2 Euclidean Propagator

The propagator (Green's function) of a massive scalar field in d dimensions satisfies

$$(-\Delta + m^2)G(x) = \delta^{(d)}(x).$$
<sup>(2)</sup>

Show that, by going to momentum space, that G is equal to

$$G(x) = \int \frac{\mathrm{d}^d p}{(2\pi)^d} \frac{e^{ipx}}{p^2 + m^2}.$$
 (3)

Show that in spherical coordinates

$$G(x) = Vol(S^{d-2}) \int \frac{p^{d-1} dp}{(2\pi)^d} \int_0^\pi d\theta \sin^{d-2}\theta \frac{e^{ipr\cos\theta}}{p^2 + m^2}$$
(4)

We now specialize to d = 3. Perform the  $\theta$ -integration and find

$$\frac{1}{(2\pi)^2 ir} \int_{-\infty}^{\infty} \mathrm{d}kk \frac{1}{k^2 + m^2} e^{ikr}.$$
 (5)

Close the contour in the upper-half plane to pick up the pole at +im. You should obtain the final result

$$G(x) = \frac{1}{4\pi r} e^{-mr}.$$
(6)

## 3 Free scalar field in the canonical formalism

(i) The solution of the classical Klein-Gordon equation can be written as

$$\phi(x) = \int \frac{\mathrm{d}^3 p}{(2\pi)^3 \sqrt{2E_{\mathbf{p}}}} (a_{\mathbf{p}} e^{-ipx} + a_{\mathbf{p}}^{\dagger} e^{ipx}), \tag{7}$$

where  $p^{\mu} = (E_{\mathbf{p}}, \mathbf{p})$  and  $p^2 = m^2$ . The canonical momentum is given by  $\pi = \partial_t \phi$ . We promote the fields to operators and impose canonical commutation relations

$$[\phi(t, \mathbf{x}), \pi(t, \mathbf{y})] = i\delta^{(3)}(\mathbf{x} - \mathbf{y}), \quad [\phi(t, \mathbf{x}), \phi(t, \mathbf{y})] = [\pi(t, \mathbf{x}), \pi(t, \mathbf{y})] = 0.$$
(8)

at equal time. Show that this implies that

$$[a_{\mathbf{p}}, a_{\mathbf{q}}^{\dagger}] = (2\pi)^{3} \delta^{(3)}(\mathbf{p} - \mathbf{q}), \quad [a_{\mathbf{p}}, a_{\mathbf{q}}] = [a_{\mathbf{p}}^{\dagger}, a_{\mathbf{q}}^{\dagger}] = 0.$$
(9)

(ii) Derive the Hamiltonian H by a Legendre transformation of the scalar field Lagrangian

$$L = \int \mathrm{d}^3 x (\partial_\mu \phi \partial^\mu \phi - m^2). \tag{10}$$

Express *H* in terms of *a* and  $a^{\dagger}$ . Substract an infinite constant, so that  $H|0\rangle = 0$ , where  $|0\rangle$  is the vacuum state, i.e.  $a_{\mathbf{p}}|0\rangle = 0$ , for all **p**.

- (iii) What is the energy of the state  $|\mathbf{p}_1, ..., \mathbf{p}_n\rangle = a_{\mathbf{p}_1}^{\dagger} \cdots a_{\mathbf{p}_n}^{\dagger} |0\rangle$ ?
- (iv) Verify that the (quantum) Heisenberg equations of motion are satisfied.
- (v) Compute the commutator  $[\phi(x^0, \mathbf{x}), \phi(y^0, \mathbf{y})]$  at non-equal times  $x^0$  and  $y^0$ . You should find

$$[\phi(x^0, \mathbf{x}), \phi(y^0, \mathbf{y})] = D(x - y) - D(y - x), \tag{11}$$

where

$$D(x) = \int \frac{\mathrm{d}^3 \mathbf{p}}{(2\pi)^3 2E_{\mathbf{p}}} e^{-ipx}.$$
(12)

(vi) Show that the measure

$$\frac{\mathrm{d}^{3}\mathbf{p}}{2E_{\mathbf{p}}}\tag{13}$$

is Lorentz invariant. Hint: you can try to evaluate the integral

$$\int d^4 p \Big|_{p^0 > 0} \delta(p^2 - m^2) f(p) \tag{14}$$

with f(p) any Lorentz invariant function.

- (vii) Use the invariance of measure to argue that (11) is zero for space-like separations (even at non-equal times). Why doesn't this argument work for time-like separations?
- (viii) Define the time-ordered two-point function

$$D_F(x-y) := \langle 0| T(\phi(x)\phi(y)) | 0 \rangle := \theta(x^0 - y^0) \langle 0| \phi(x)\phi(y) | 0 \rangle + \theta(y^0 - x^0) \langle 0| \phi(y)\phi(x) | 0 \rangle.$$
(15)

Show that it can be written as

$$D_F(x) = \int \frac{\mathrm{d}^4 p}{(2\pi)^4} \frac{i e^{-ipx}}{p^2 - m^2 + i\epsilon},\tag{16}$$

where  $\epsilon$  is a small positive number, which we send to zero in the end. (*Hint:* Start from equation (16) and close the  $p_0$  contour appropriately. Find the poles of the integrand and use Cauchy's integral theorem.)

(ix) Prove that

$$T(\phi(x)\phi(y)) = N(\phi(x)\phi(y)) + \phi(x)\phi(y),$$
(17)

where

$$\vec{\phi}(x)\vec{\phi}(y)) = \begin{cases} [\phi^+(x), \phi^-(y)], \ x^0 > y^0; \\ [\phi^+(y), \phi^-(x)], \ x^0 < y^0. \end{cases}$$
(18)

In the above expression we defined

$$\phi^{+}(x) = \int \frac{\mathrm{d}^{3}\mathbf{p}}{(2\pi)^{3}\sqrt{2E_{\mathbf{p}}}} a_{\mathbf{p}} e^{-ipx}, \quad \phi^{-}(x) = \int \frac{\mathrm{d}^{3}\mathbf{p}}{(2\pi)^{3}\sqrt{2E_{\mathbf{p}}}} a_{\mathbf{p}}^{\dagger} e^{ipx}.$$
 (19)

The operation N is defined through its action on creation and annihilation operators. For any product of such operators, it puts annihilation operators  $a_{\mathbf{p}}$  to the right, and creation operators  $a_{\mathbf{p}}^{\dagger}$  to the left. So for example

$$N(a_{\mathbf{p}}a_{\mathbf{q}}^{\dagger}a_{\mathbf{l}}) = a_{\mathbf{q}}^{\dagger}a_{\mathbf{p}}a_{\mathbf{l}}.$$
(20)

You should also find that

$$\phi(x)\phi(y) = D_F(x-y). \tag{21}$$

What is the vacuum expectation value of normal-ordered product of operators?